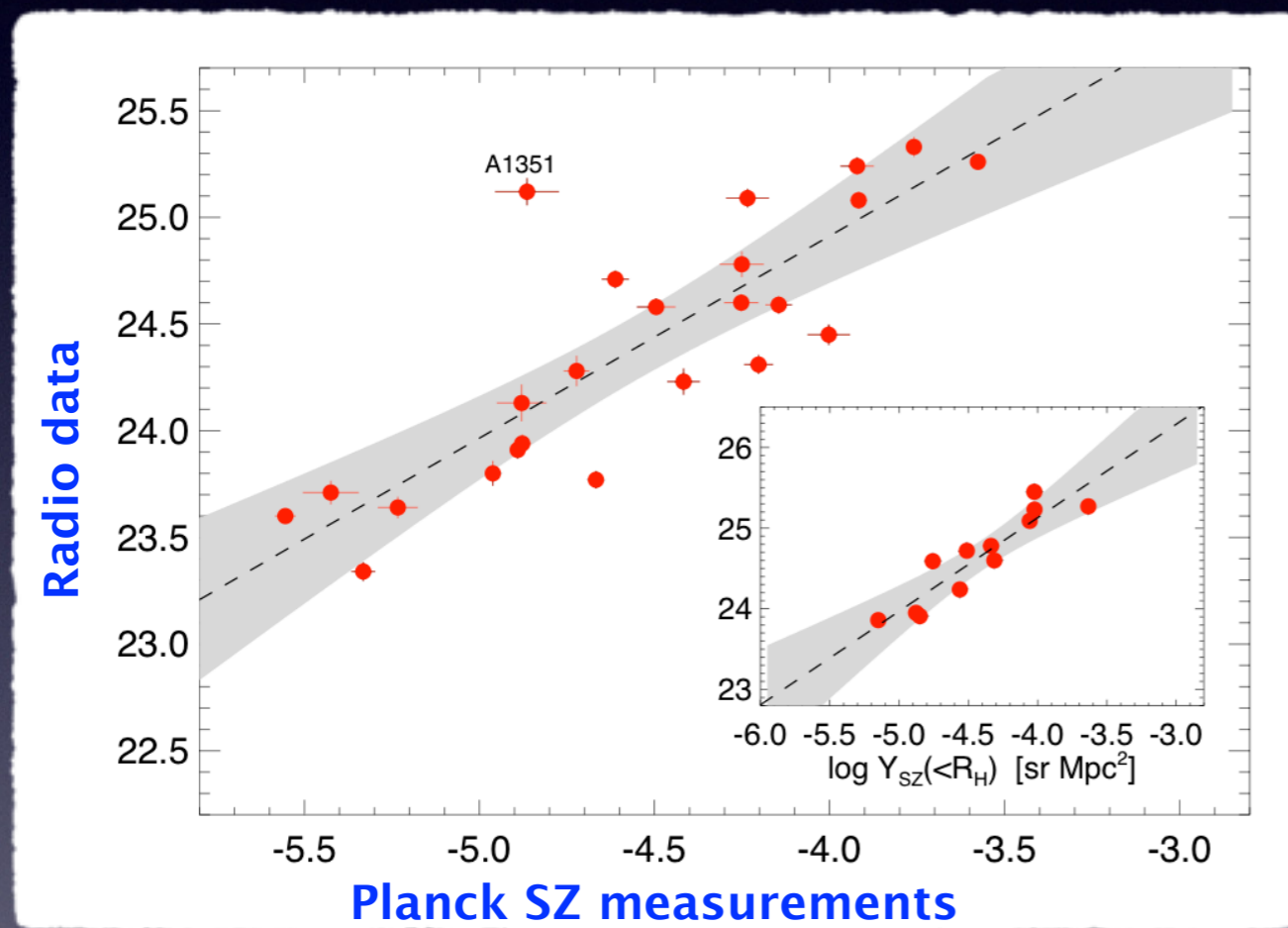


AN SZ TAKE ON

Cluster Diffuse Radio Emissions



Kaustuv Basu

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**Argelander Institute for Astronomy
University of Bonn**

With inputs from:

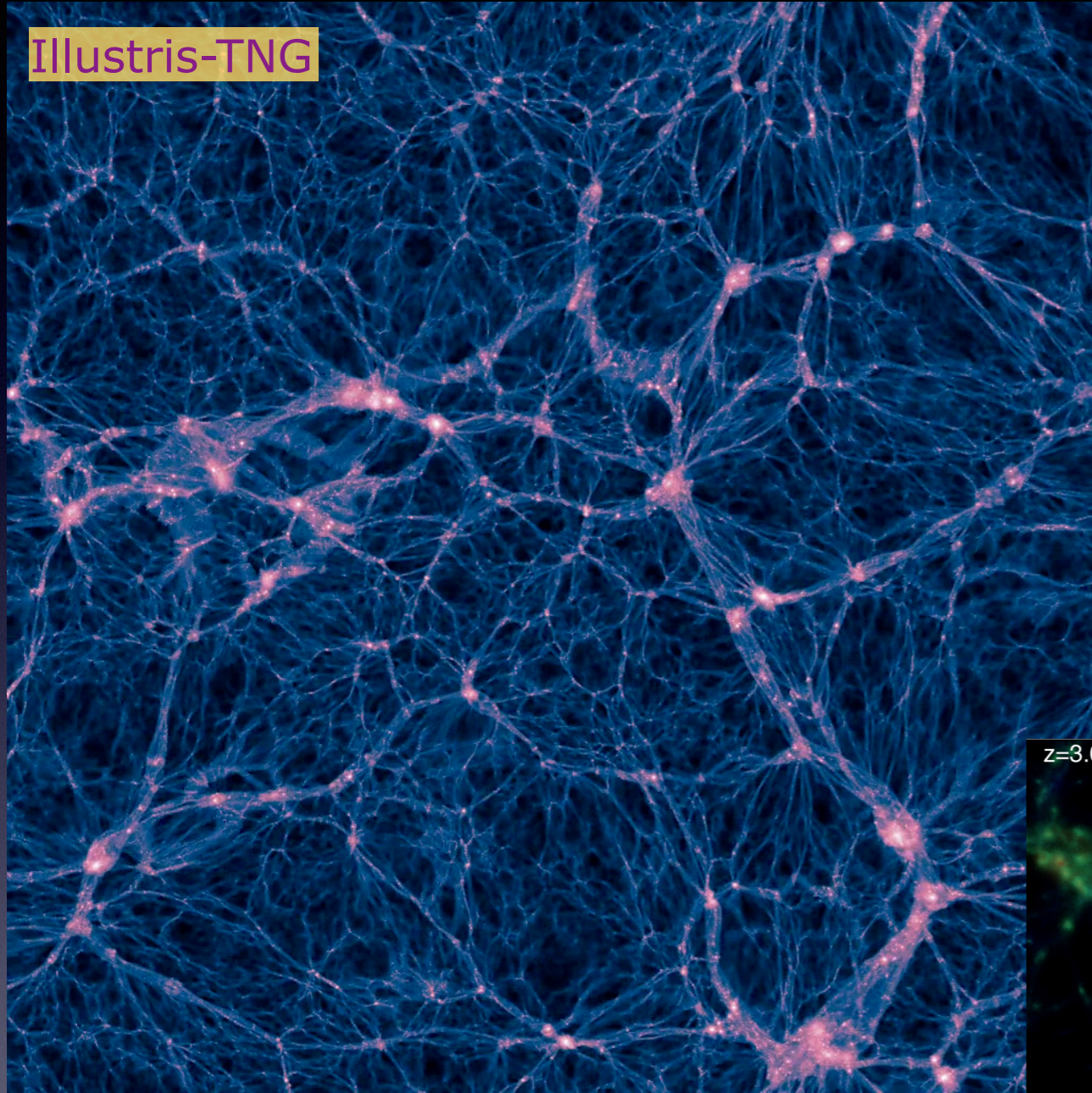
Martin Sommer, Vyoma Muralidhara,
Jens Erler, Franco Vazza,
and many others

In this talk..

- We will learn about the two main types of diffuse synchrotron emissions in galaxy clusters, namely, radio halos and relics, and some open questions regarding them
- We will discuss the thermal SZ effect and how its measurements helped us to gain a better understanding of the nonthermal radio emission
- We will introduce the nonthermal SZ effect, which is caused by the relativistic electrons, and the prospects of measuring its signal

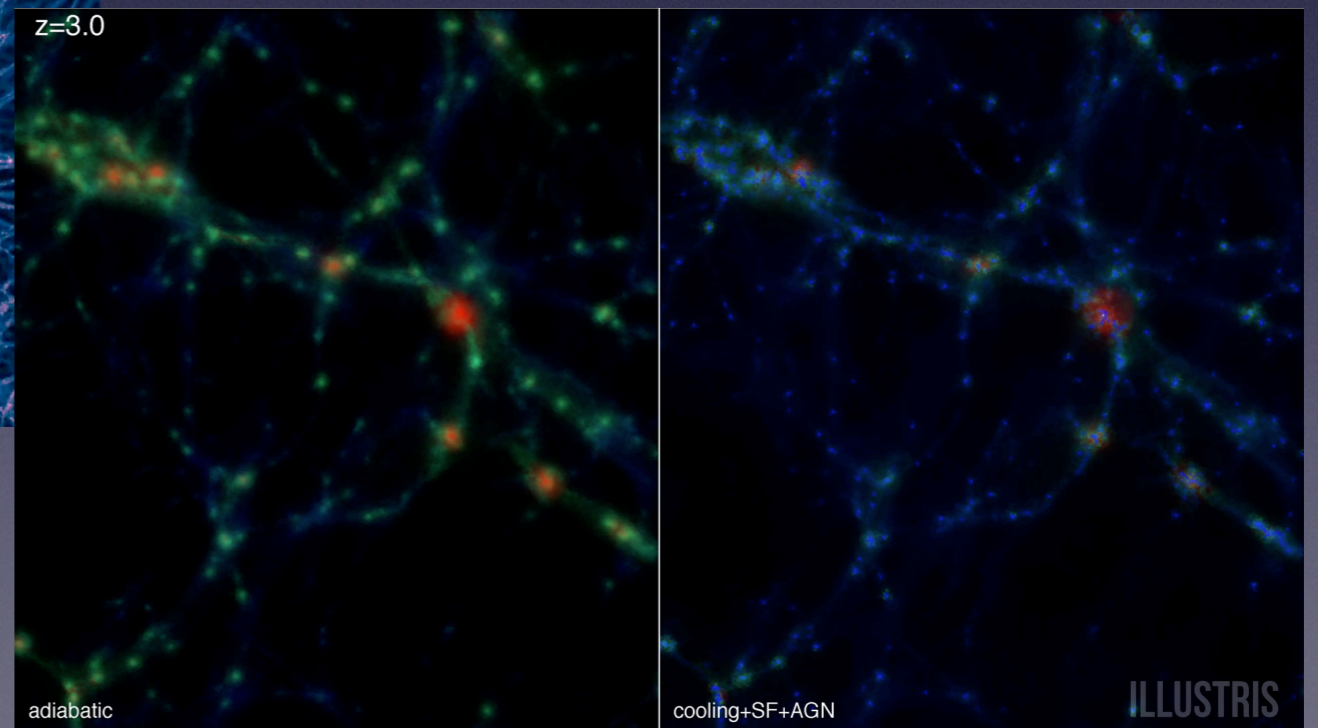
The Big Picture

Illustris-TNG



Slices of the universe at $z=0$ (DM only)

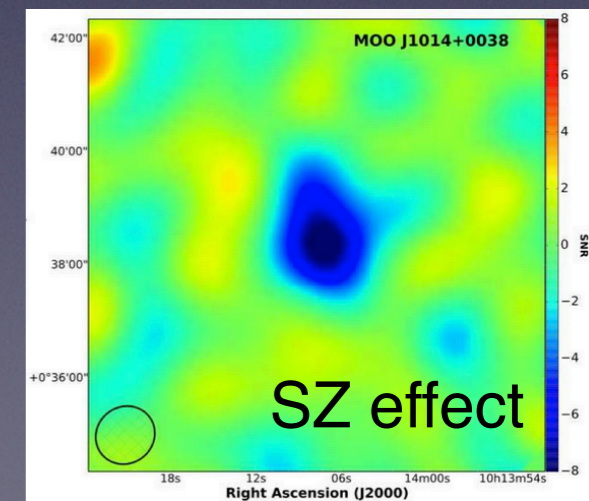
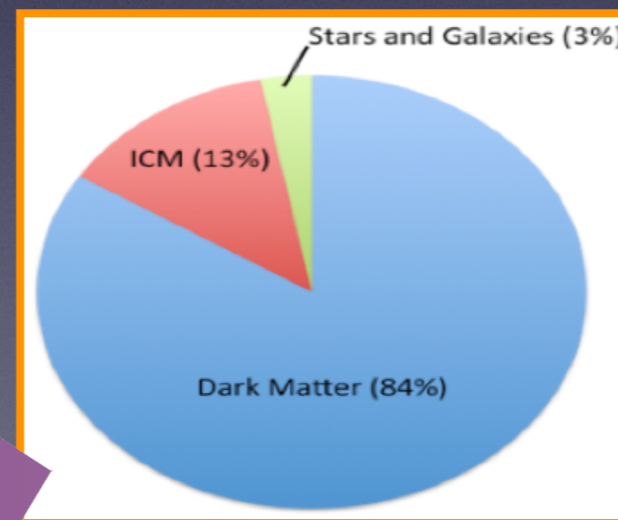
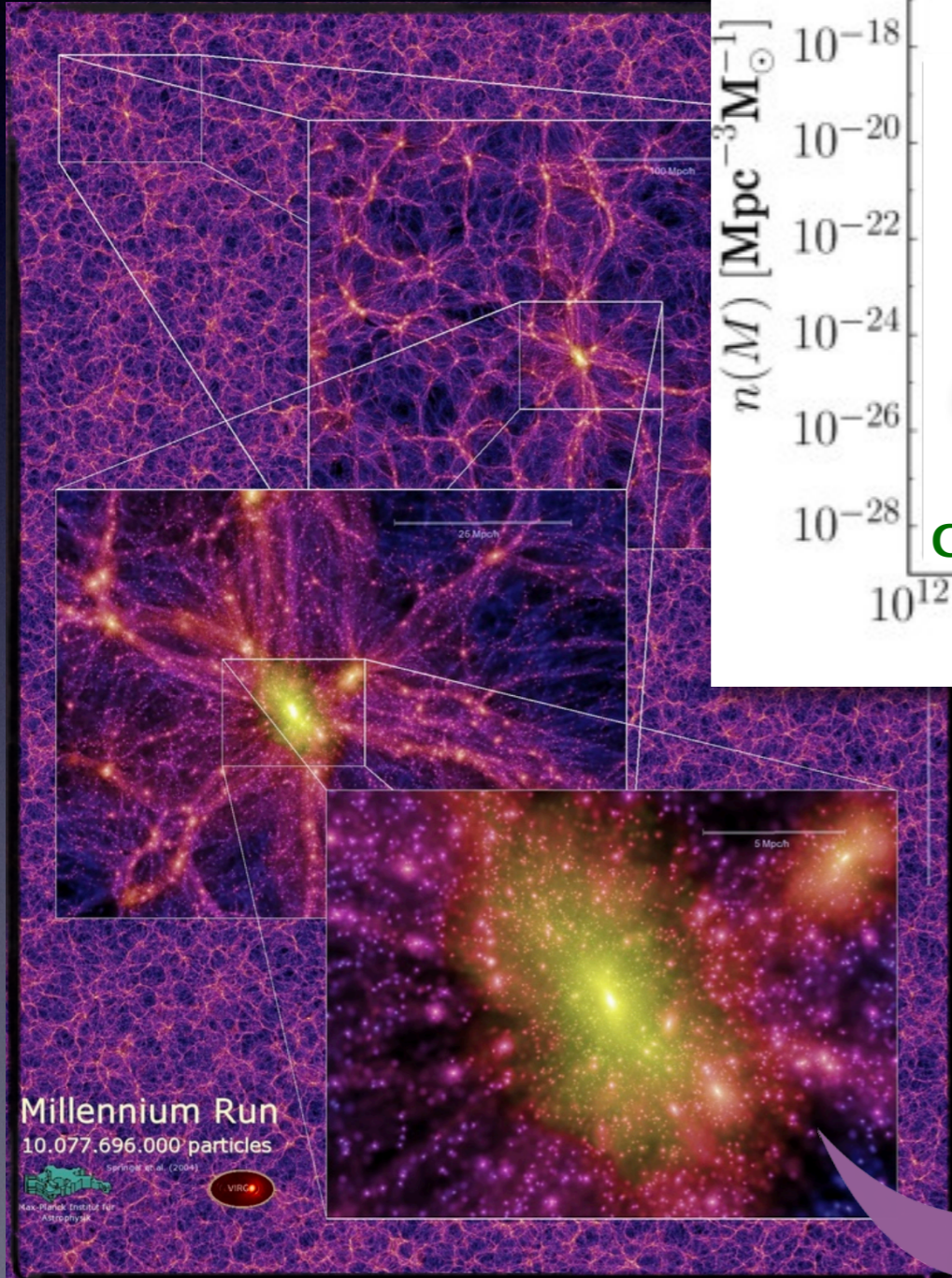
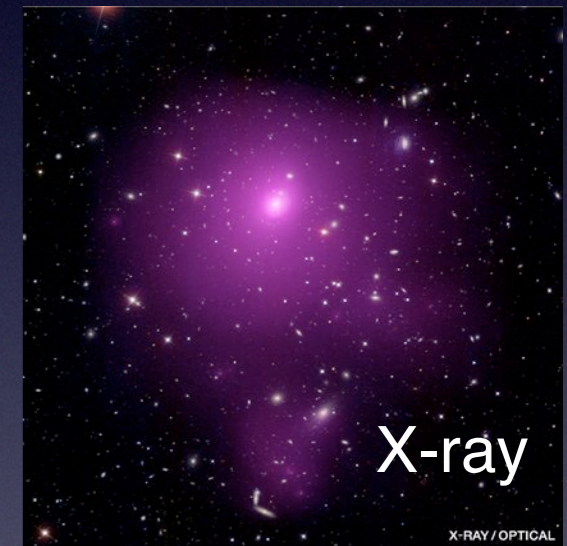
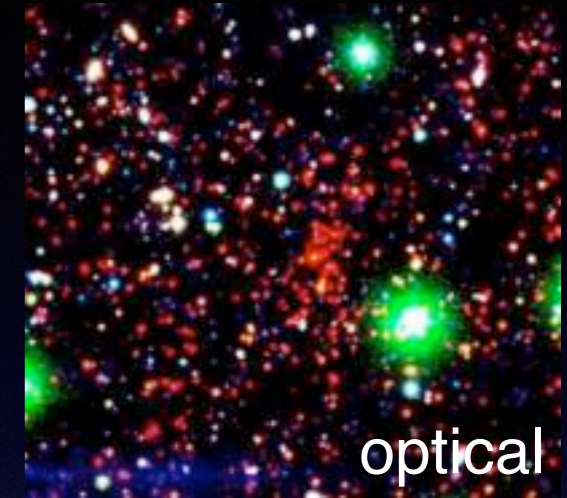
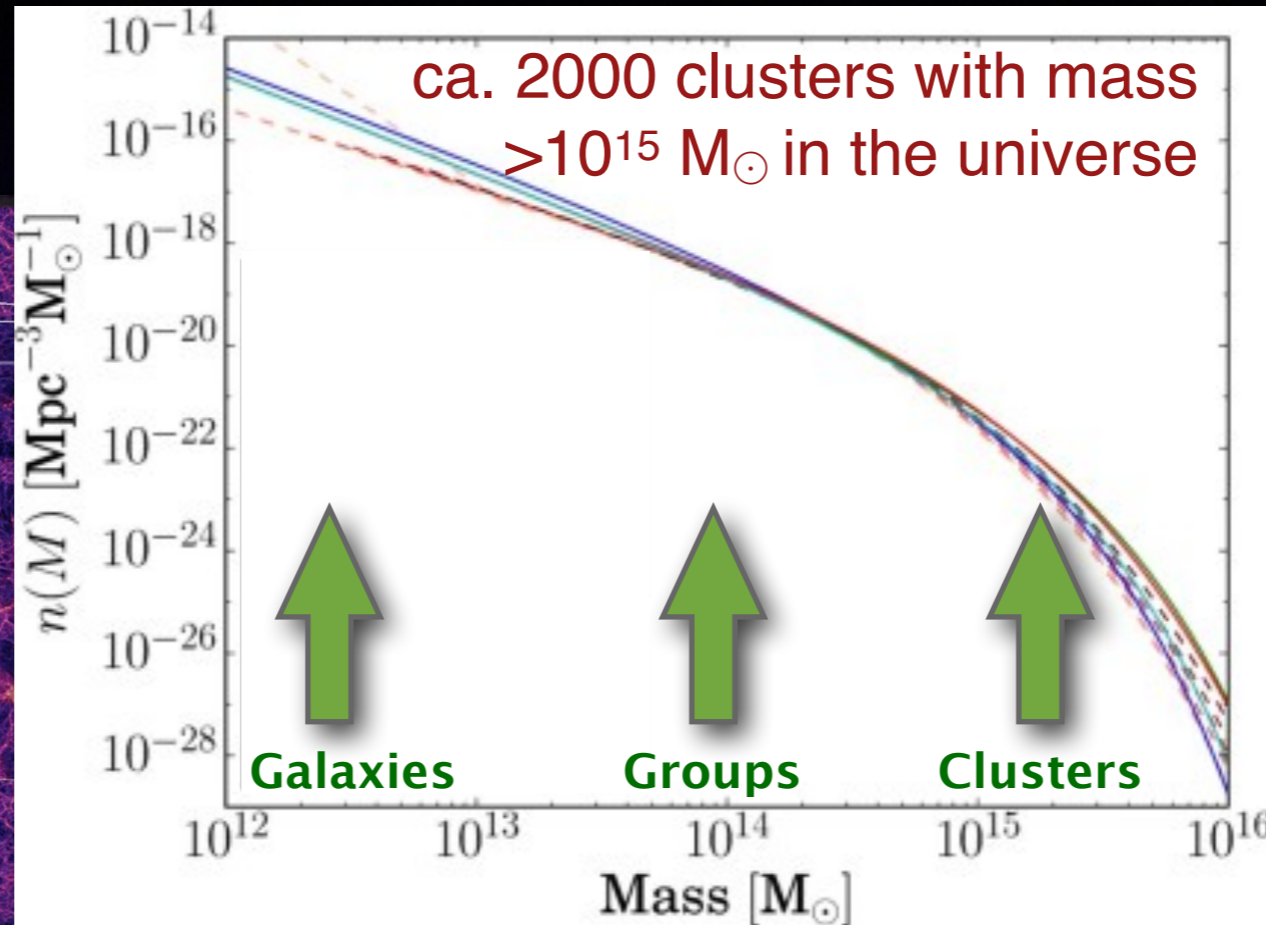
Impact of baryonic physics



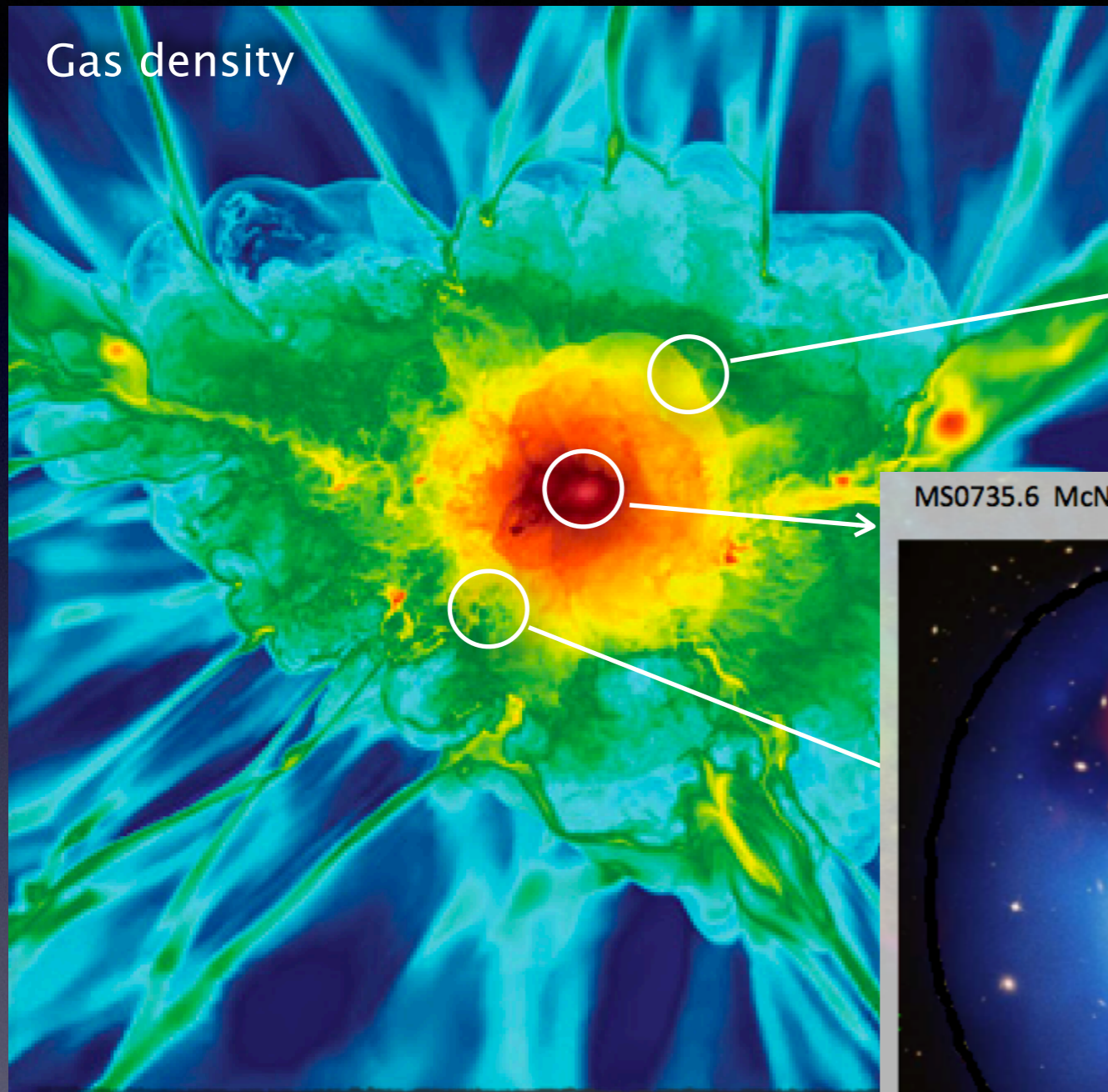
- How do large-scale structures in the Universe evolve? What role do the baryons play? Where are most of these baryons?
- Are nonthermal, ultra-relativistic particles (cosmic rays) also important in this cosmic evolution?
- What can we learn by cross-comparing the thermal and nonthermal emissions from the large-scale structures?

Why Galaxy Clusters?

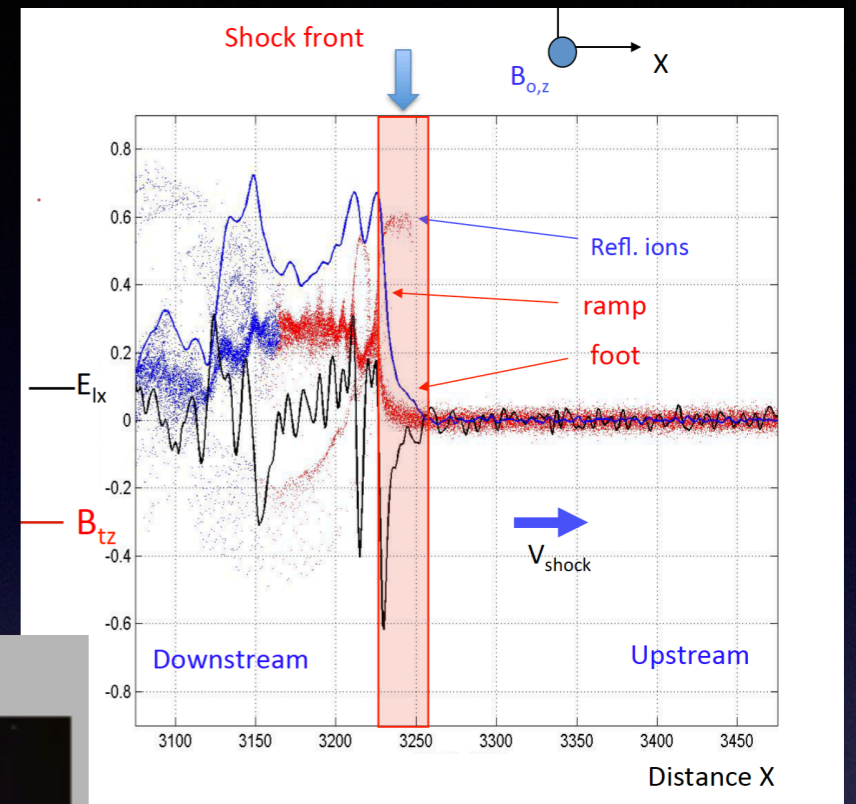
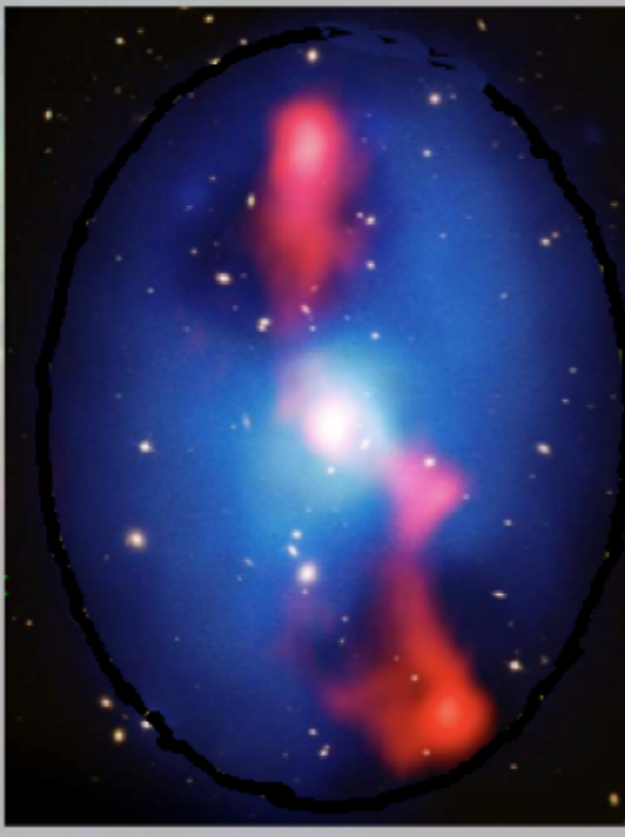
The “massive end” of the halo mass function.



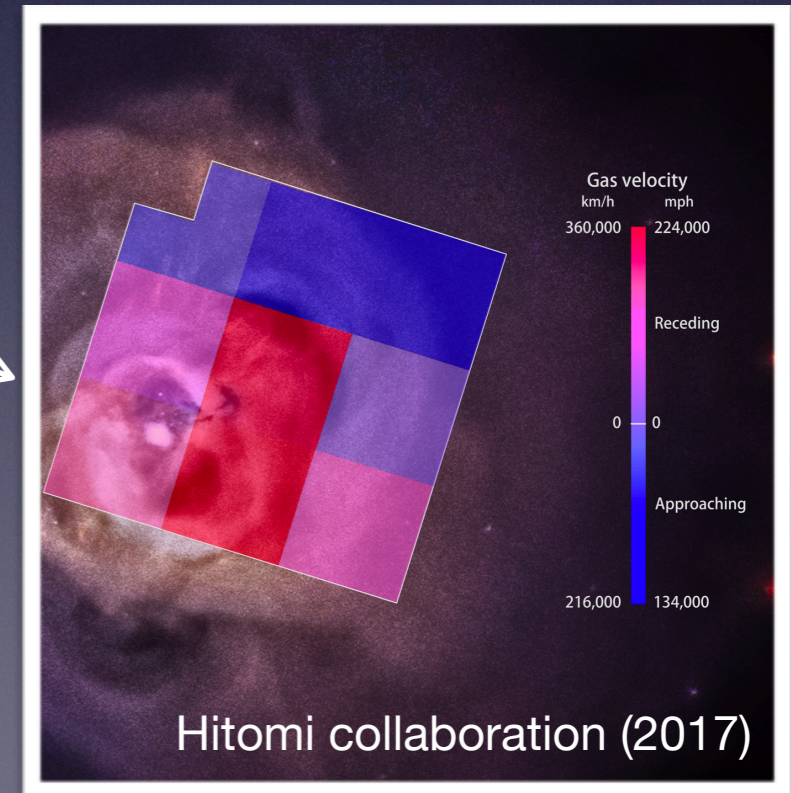
Astrophysics of galaxy clusters



MS0735.6 McNamara+'05

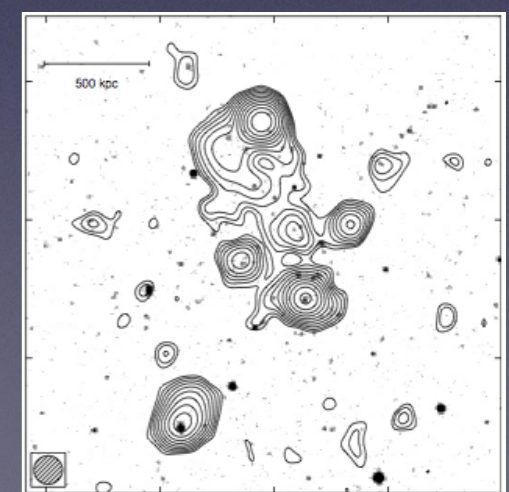
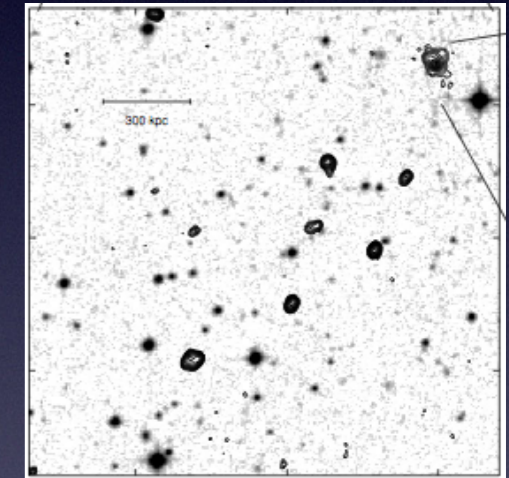
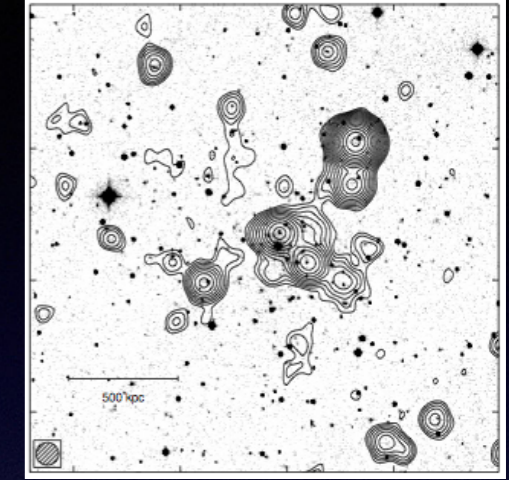
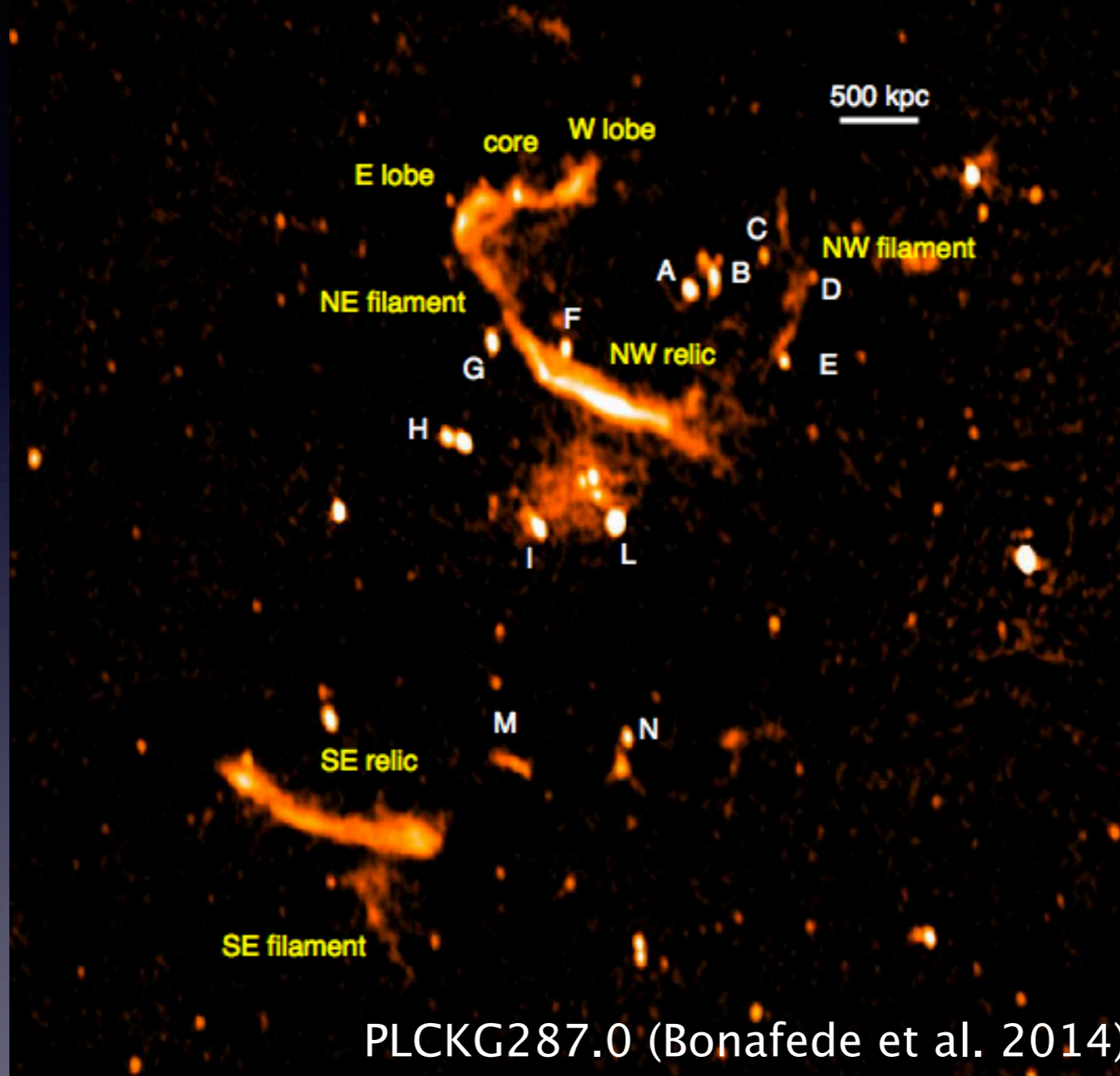
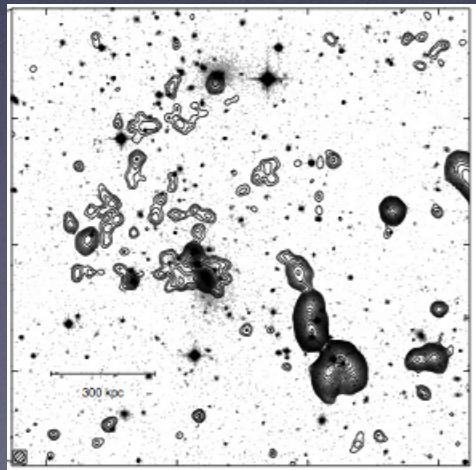
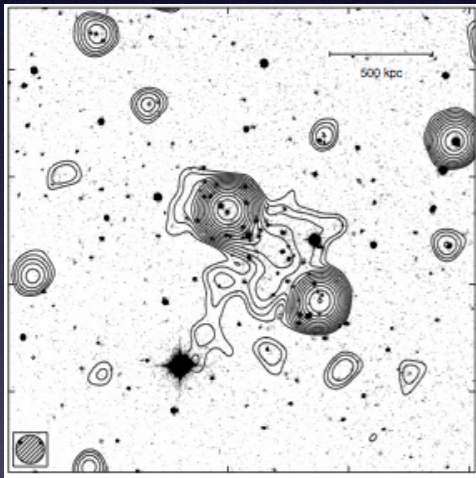
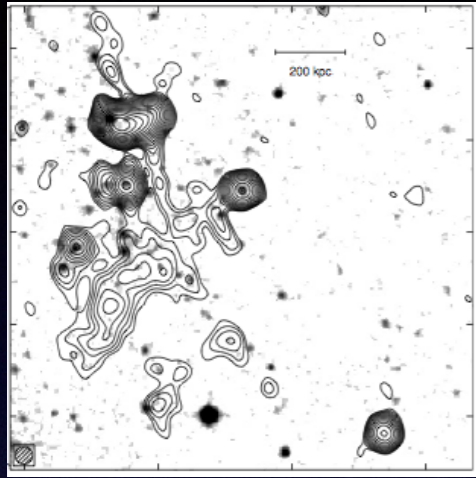


Marcowith et al. (2012)



Part I: Radio halos

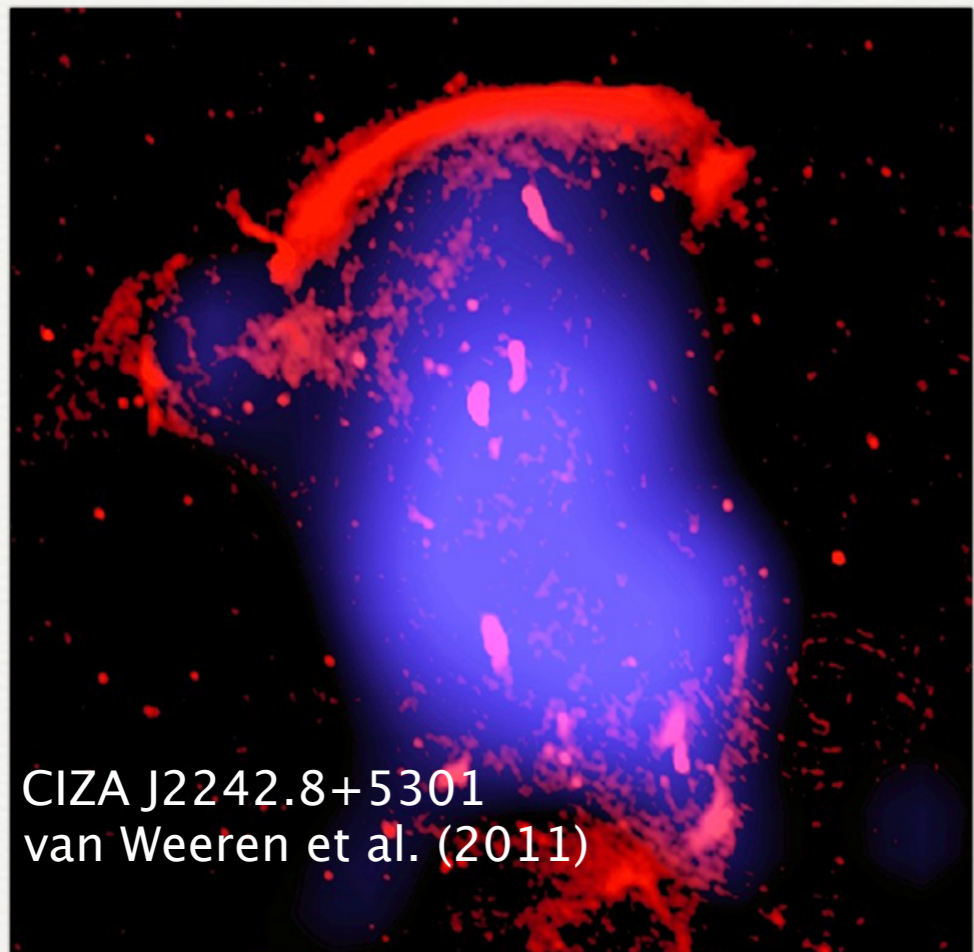
Radio emission from galaxy clusters



All contour plots from Giovannini et al. (2009)

PLCKG287.0 (Bonafede et al. 2014)

The nonthermal ICM



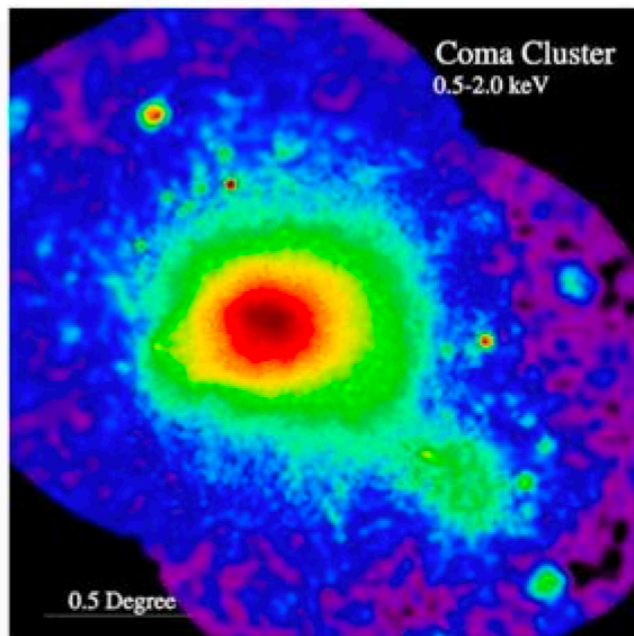
AGN related

Radio jets and lobes,
WAT sources, plasma
bubbles, AGN relics,
radio BCGs, etc.

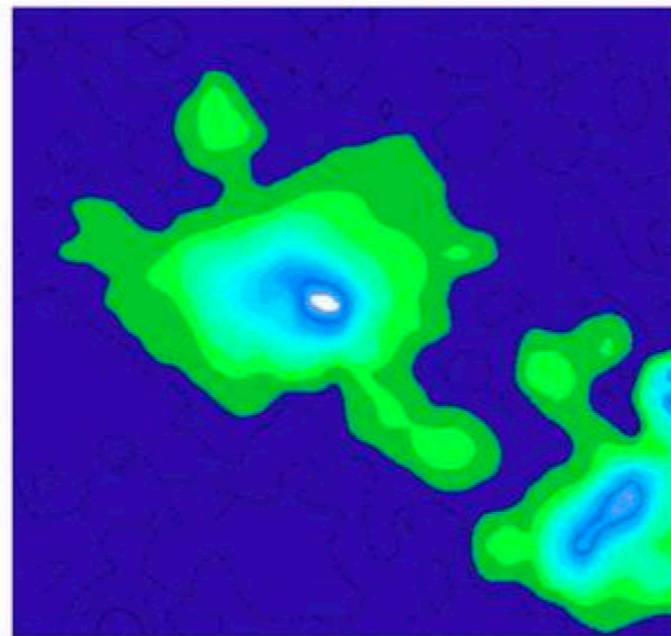
Diffuse radio emissions

Radio halos, radio
relics, mini halos,
radio gischt, radio
phoenix, etc.

MS0735.6 McNamara+'05



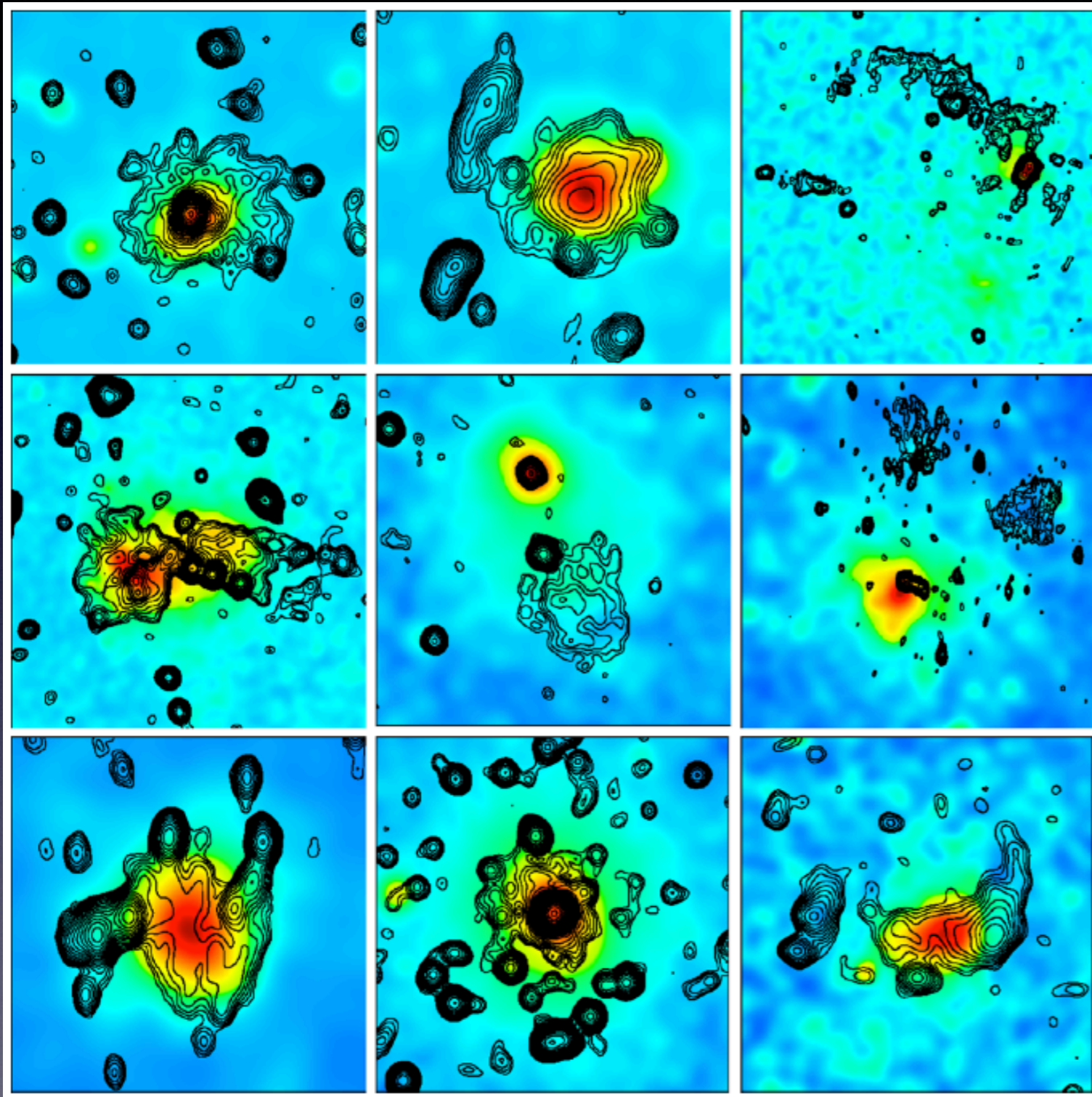
thermal X-ray emission



radio synchrotron emission



Radio halos and relics



Radio halos: $L_{1.4 \text{ GHz}} \sim 10^{24-25} \text{ W/Hz}$

- Mpc scale diffuse sources near cluster centers
- Low surface brightness and generally not polarized
- Mostly steep spectrum ($\alpha \sim 1.2$)
- Morphology roughly similar to X-ray or SZ emission, no severe projection bias

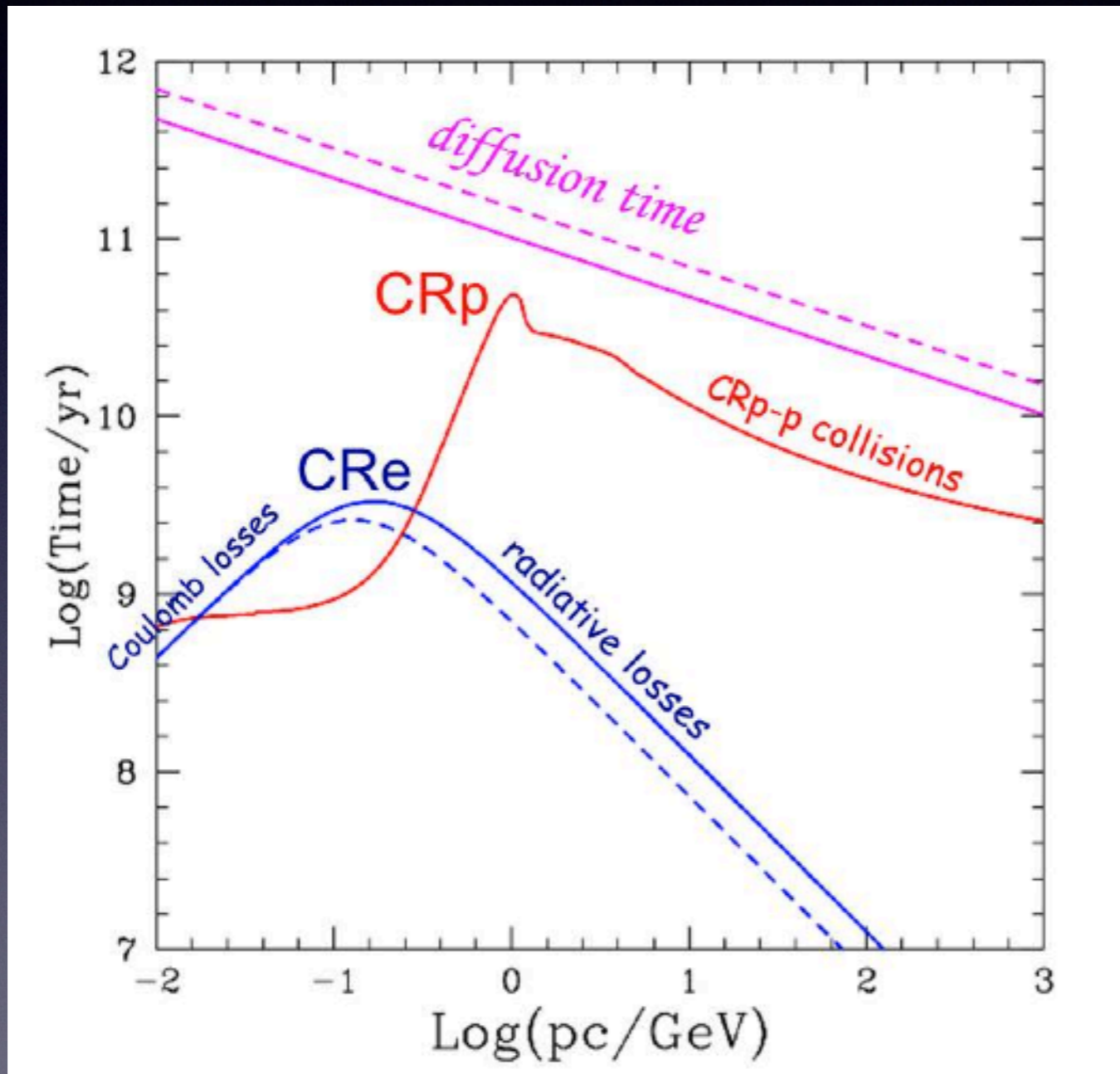
Radio relics: $L_{1.4 \text{ GHz}} \sim 10^{23-25} \text{ W/Hz}$

- Mpc scale elongated sources near cluster periphery
- Higher surface brightness and polarized
- Also steep spectrum ($\alpha \sim 1.2$)
- Morphology resembles shock fronts, subject to projection bias

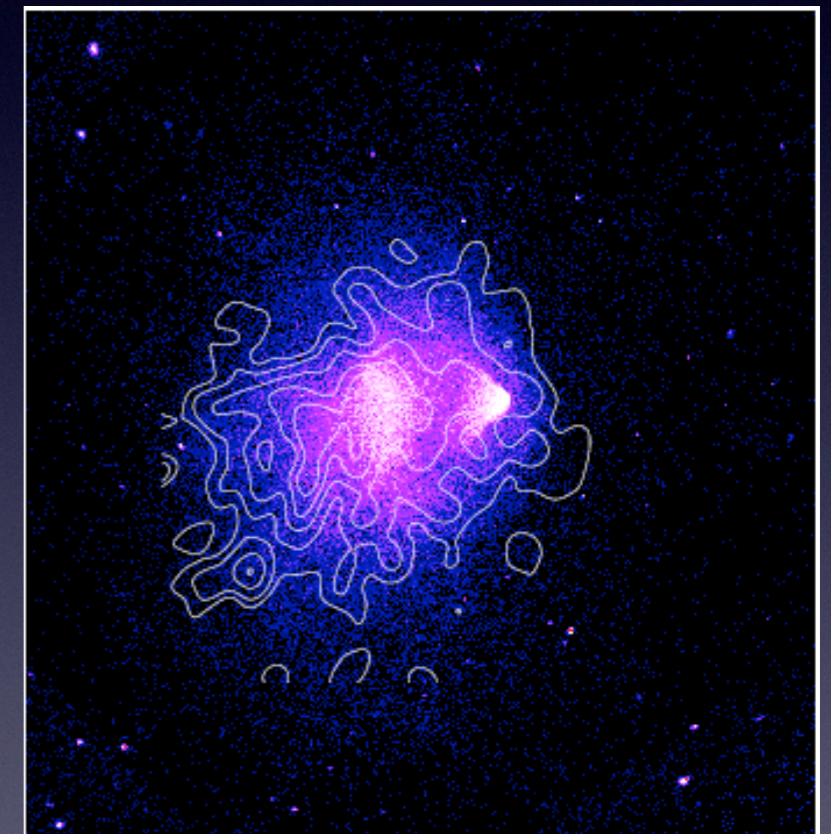
Gallery taken from Feretti et al. (2012) Color \rightarrow X-ray

The “problem” with radio halos

Radio halos imply GeV energy electrons filling up cluster volume ($\sim \text{Mpc}^3$).
But CRe lifetimes are much shorter ($\sim 10^8$ years) than cluster dynamic timescales.



Some in-situ acceleration is necessary for the CRe



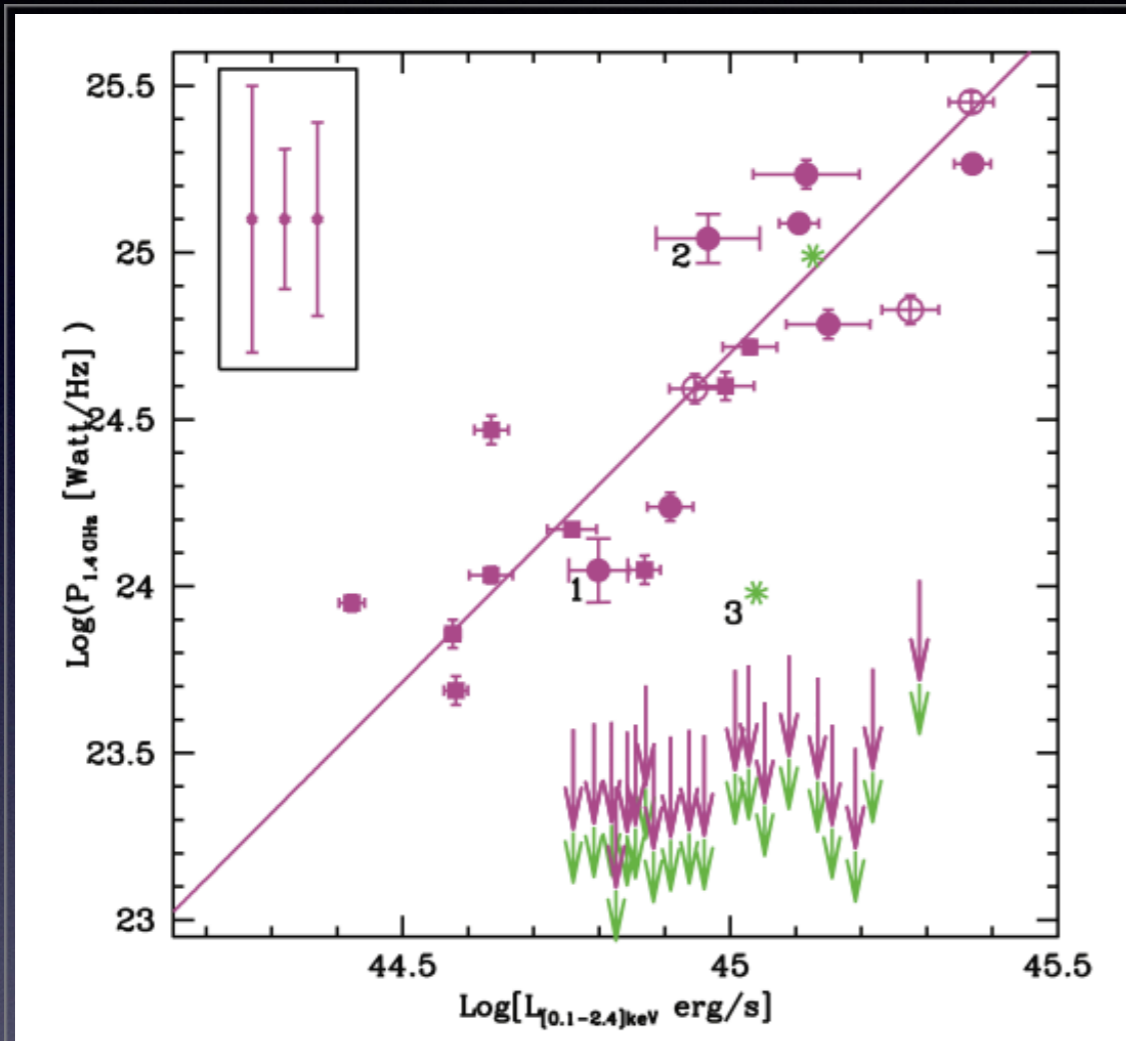
Radio halo in Bullet cluster
(Liang et al. 2000)

(Fig. from Brunetti & Jones 2014)

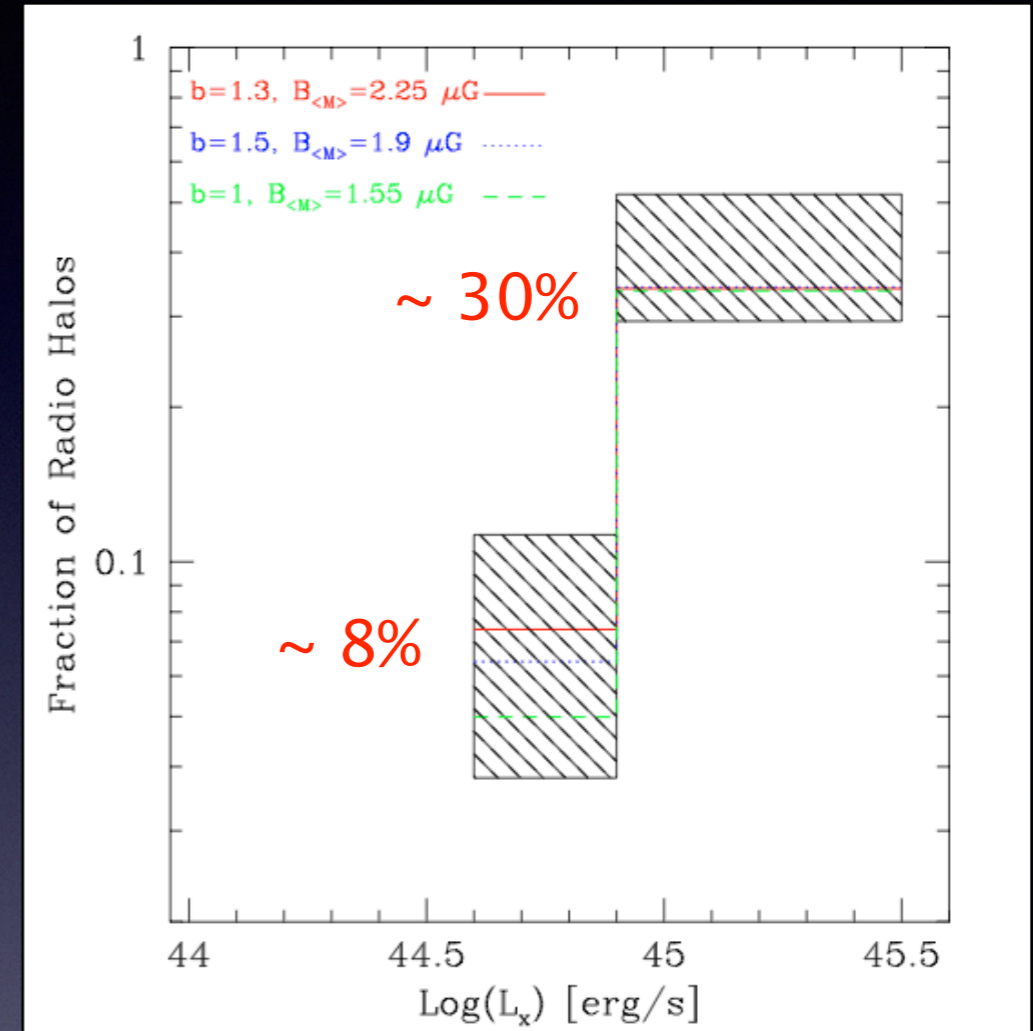
The "consensus" for radio halos

There is a strong bi-modality

They are rare ~60 known halos



Brunetti et al. (2007)



Cassano et al. (2010)

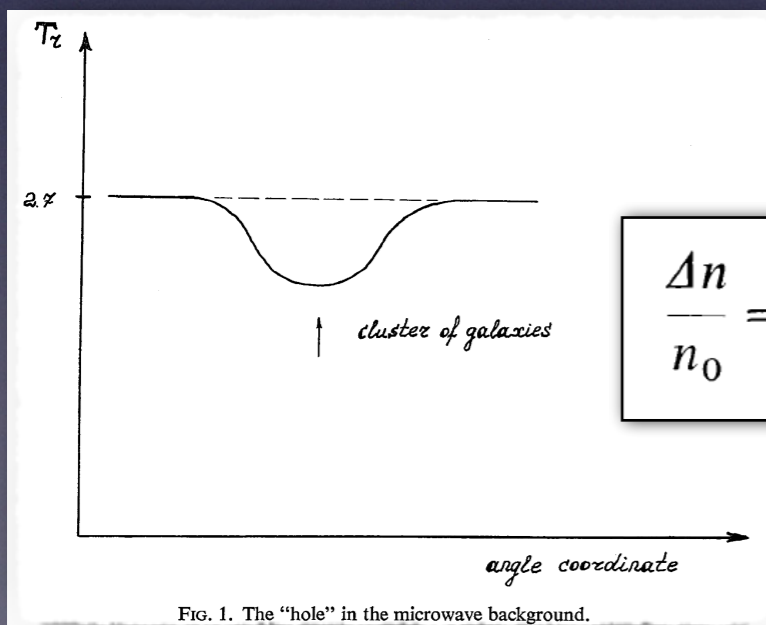
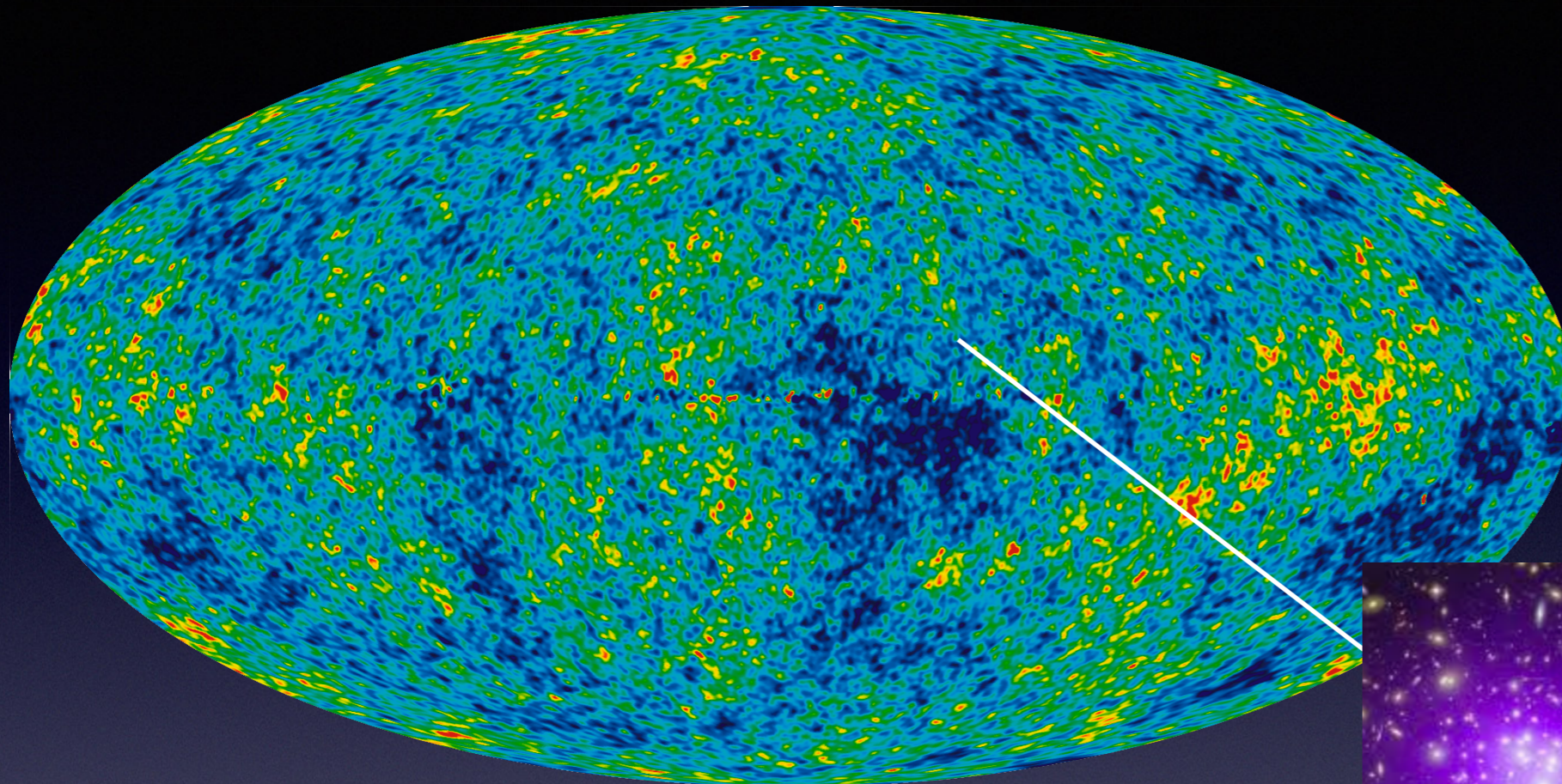
Primary models (or re-acceleration models):
 electrons are accelerated in diffusive shocks via turbulence induced by cluster mergers, through inefficient Fermi-I process

Secondary models (or hadronic models):
 e⁻/e⁺ are produced from collision between thermal ions and cosmic ray protons, the latter having significantly longer lifetimes

*Two competing models for radio halo origin
 (preference towards merger driven turbulence)*

An “SZ take” on radio halos

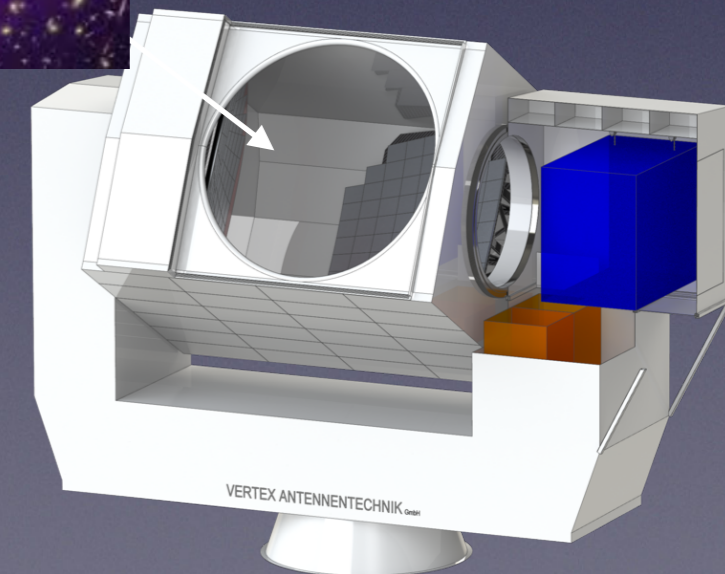
The Sunyaev-Zeldovich Effect



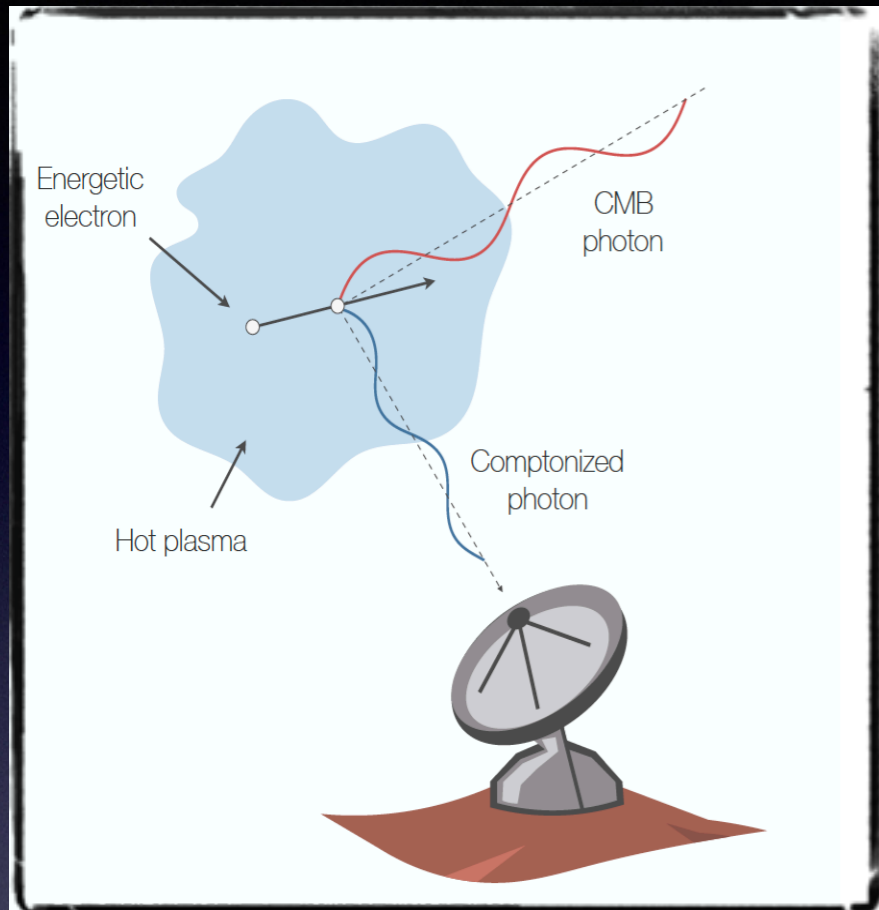
tSZ Effect

$$\frac{\Delta n}{n_0} = \frac{\Delta J}{J_0} = xy \frac{e^x}{e^x - 1} \left\{ \frac{x}{\tanh(x/2)} - 4 \right\}.$$

R. A. Sunyaev & Ya. B. Zeldovich
 Comments on Astrophysics &
 Space Physics, 1972



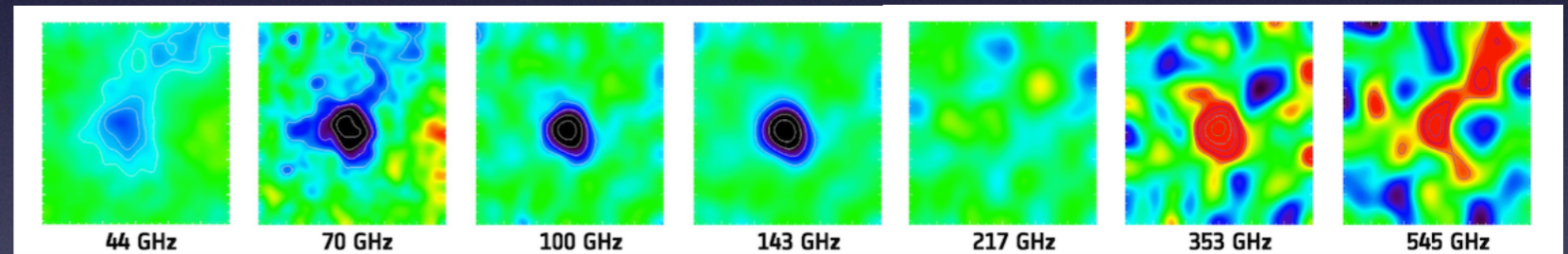
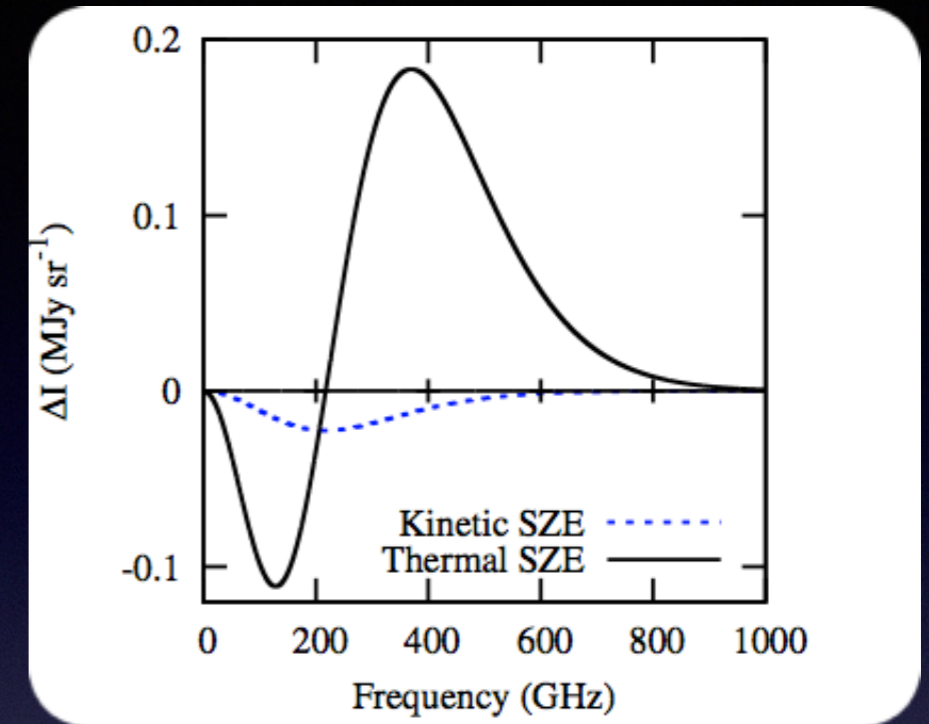
The Sunyaev-Zeldovich effect



Inverse Compton scattering producing unique spectral distortion on the background CMB.

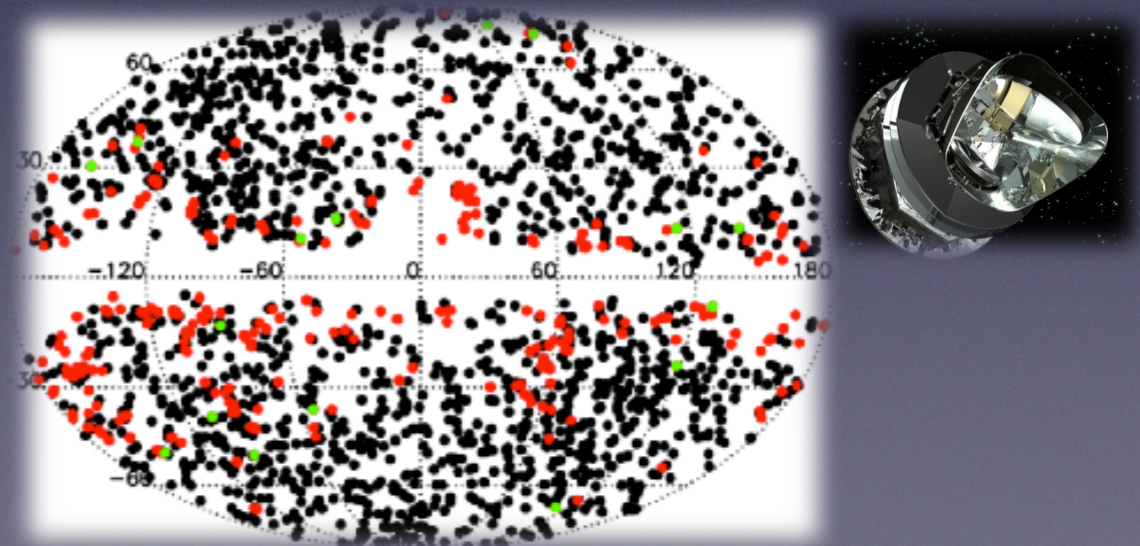
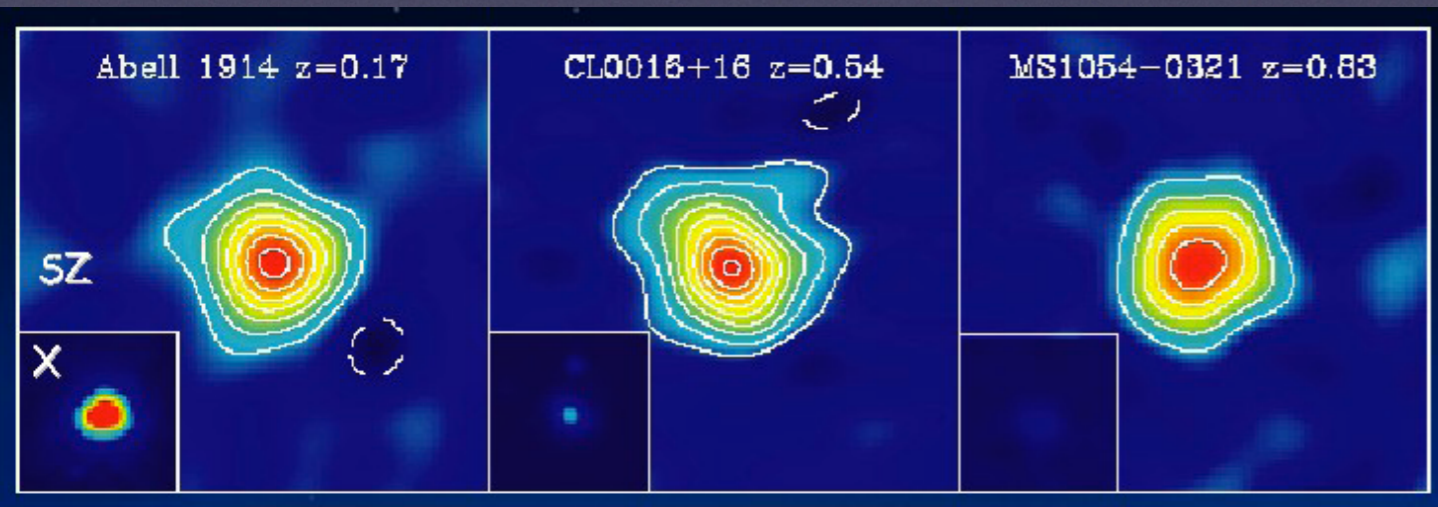
An ideal tool for finding and characterizing galaxy clusters.

Resolved source flux is redshift independent and scales linearly with the gas density!

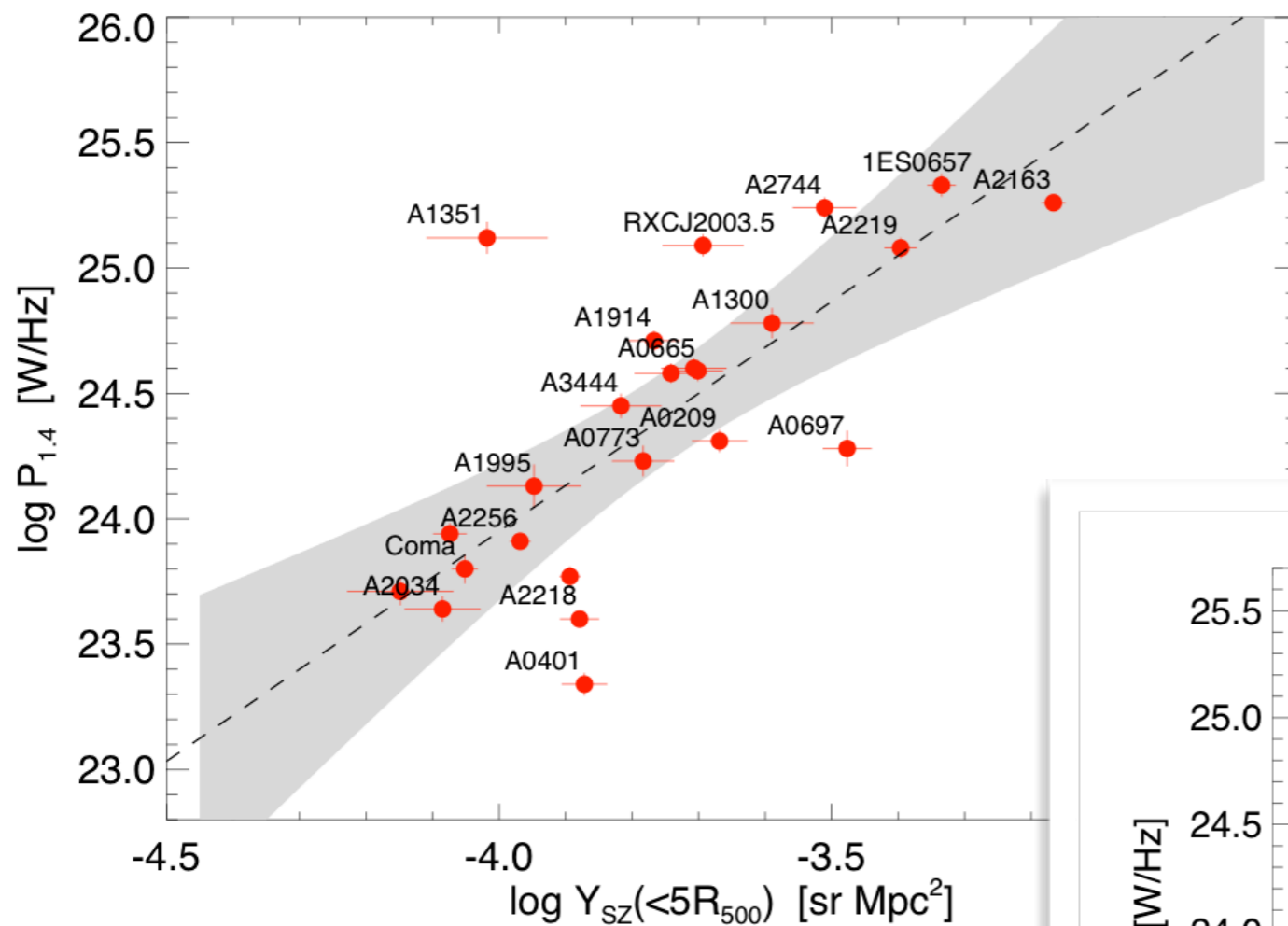


Credit: Planck collaboration

Carlstrom et al.



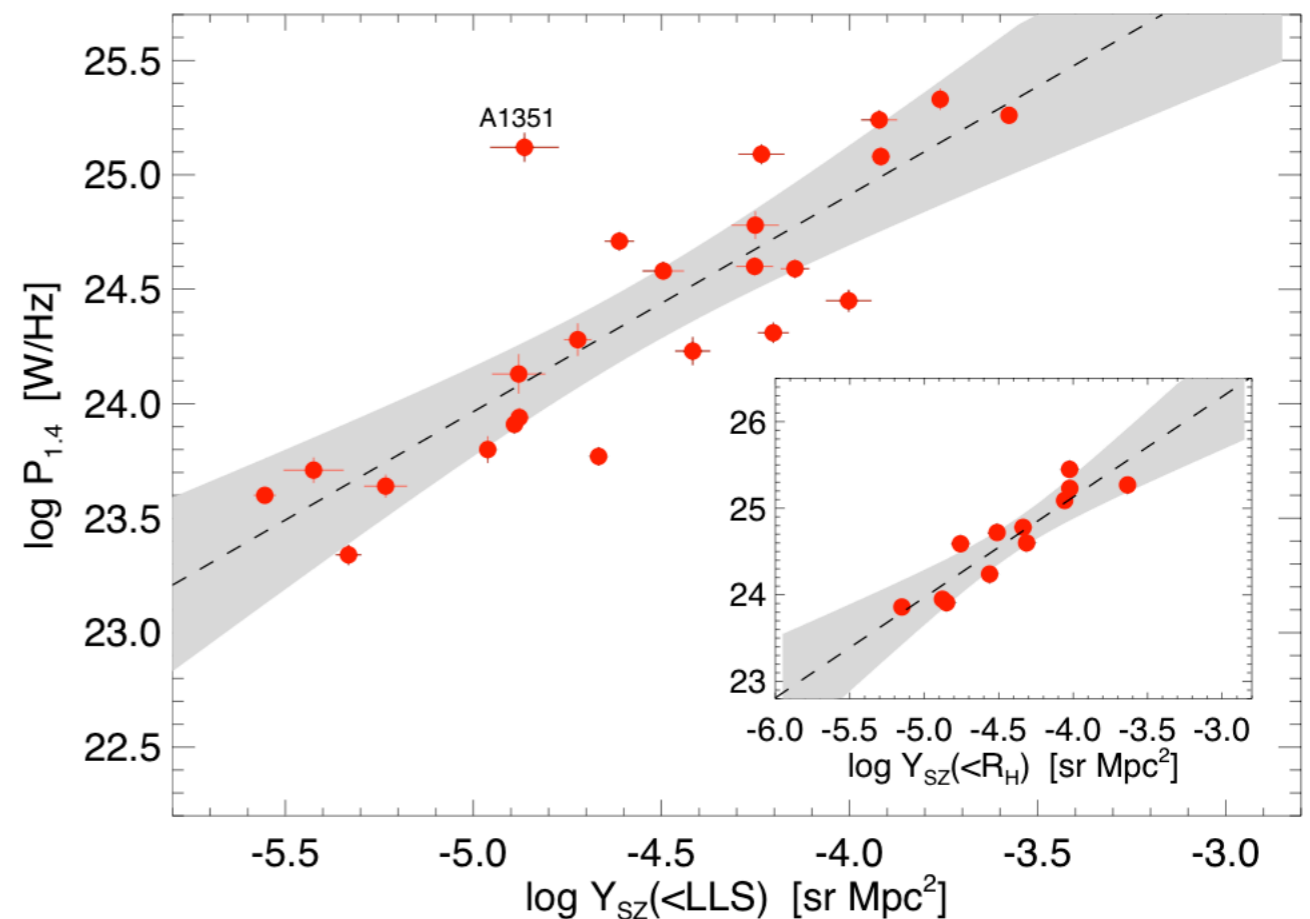
SZ correlation for radio halos



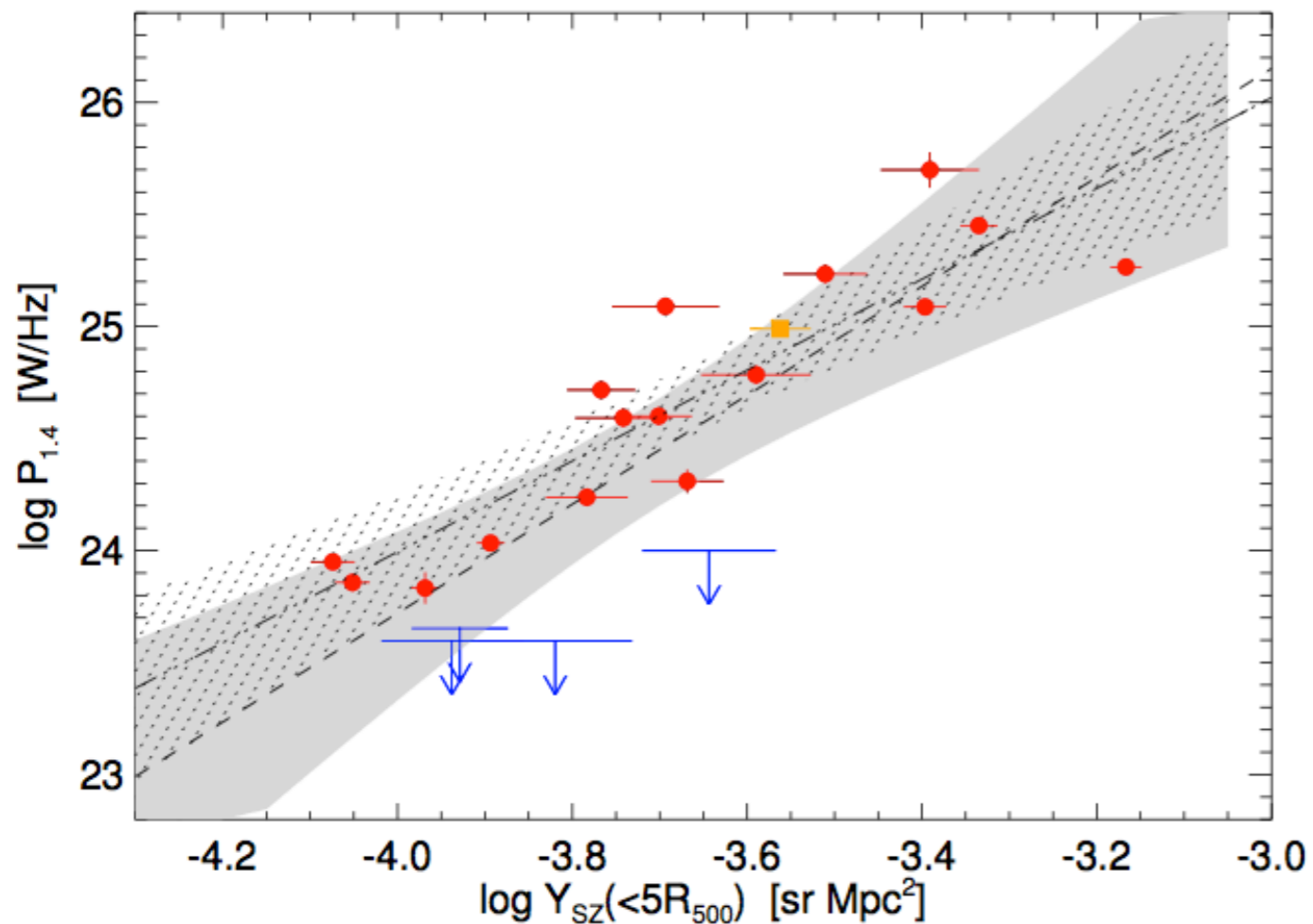
The cluster SZ signal and radio halo power are correlated (as expected from known X-ray correlation)

Basu (2012)

The correlation becomes tighter (and roughly linear) when the SZ signal is scaled to within the radio halo radius



No strong evidence for bimodality

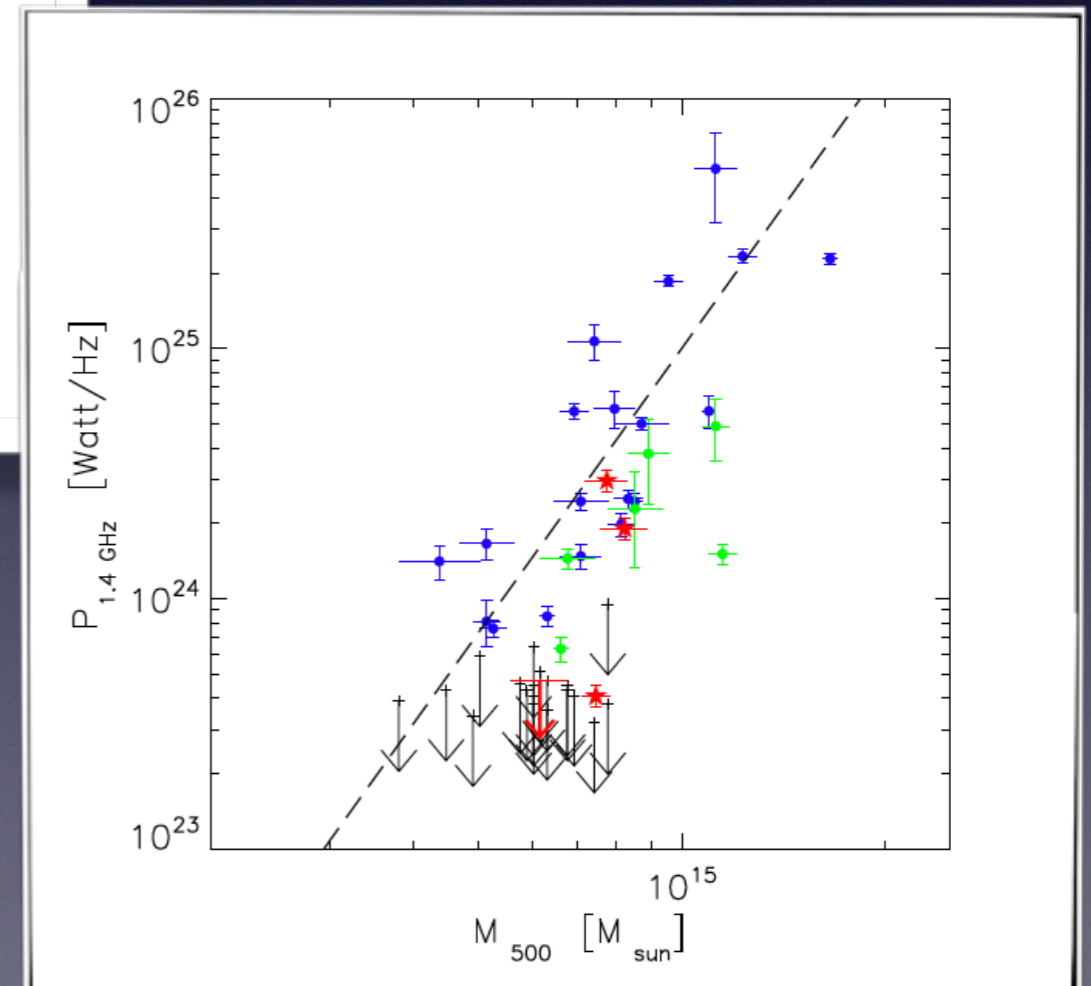


Basu (2012)

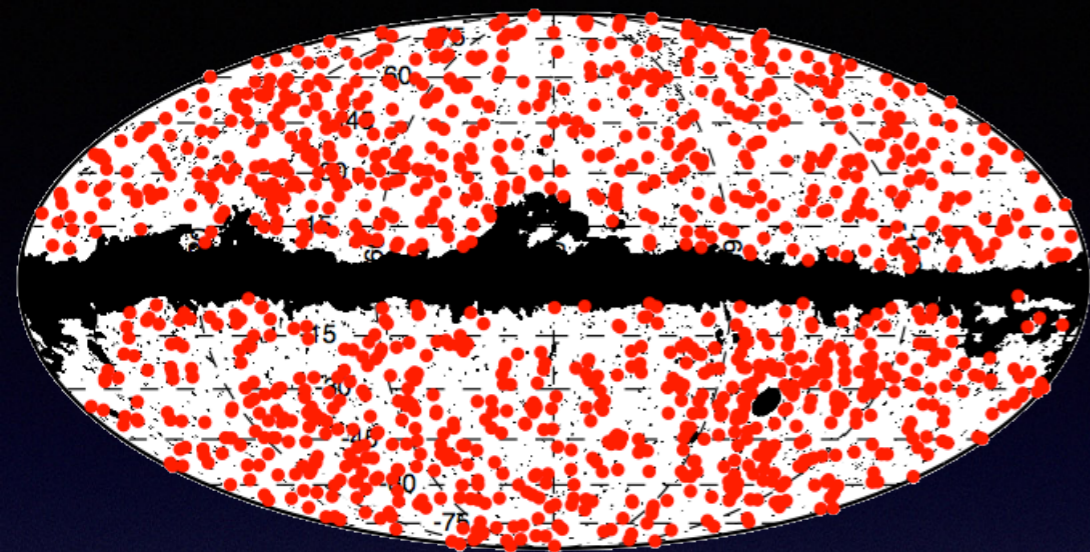
Bonafede et al. (2015)
RHs from *Planck* cluster catalog

We found from a posteriori selection of radio halo clusters, taken from the *Planck* catalog, that the bimodality is weak in the radio-SZ correlation.

This has now been independently verified by many other researchers!



Comparison between SZ/X-ray selection



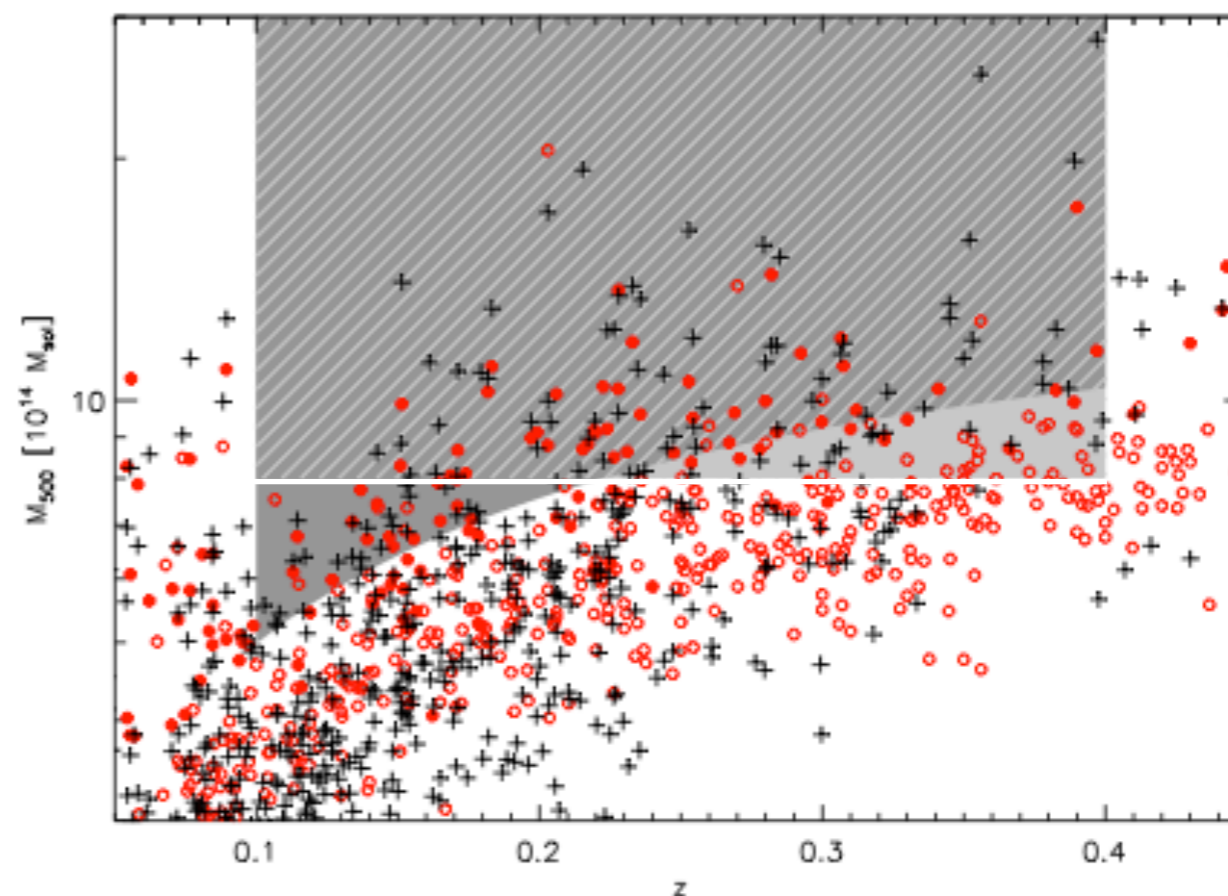
PSZ clusters (Planck coll. 2013)

We used low-sensitivity NVSS survey data to characterize radio halos. Most of these were non-detections, so we developed a regression method to particularly deal with that.

We aimed to find the mass correlation of radio power, as traced by L_x or Y_{sz} , and determine the “radio off” fraction that do not belong to this power-law scaling.

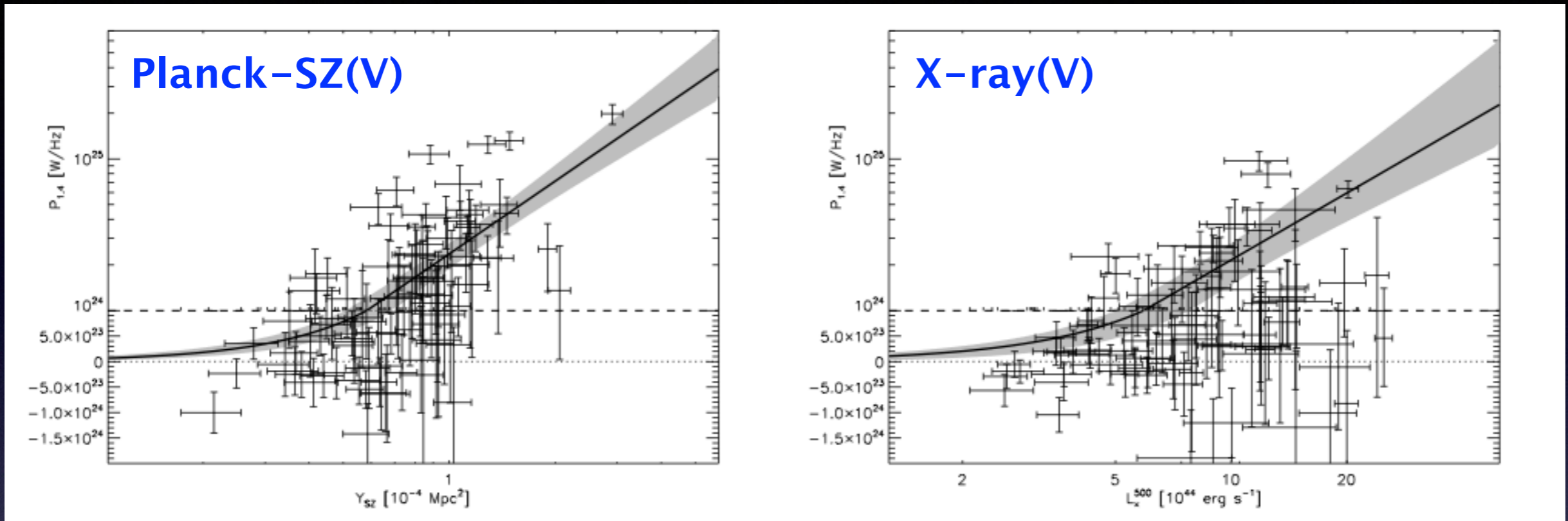
Sub-sample	Mass limit	Primary selection	Flagged due to bad data	Final sample
PSZ(V)	z -dependent	90	1	89
X(V)	z -dependent	86	1	85
PSZ(C)	$8 \times 10^{14} M_{\odot}$	79	0	79
X(C)	$8 \times 10^{14} M_{\odot}$	78	1	77

PSZ and REFLEX+eBCS+MACS



Sommer & Basu (2014)

SZ vs X-ray selection

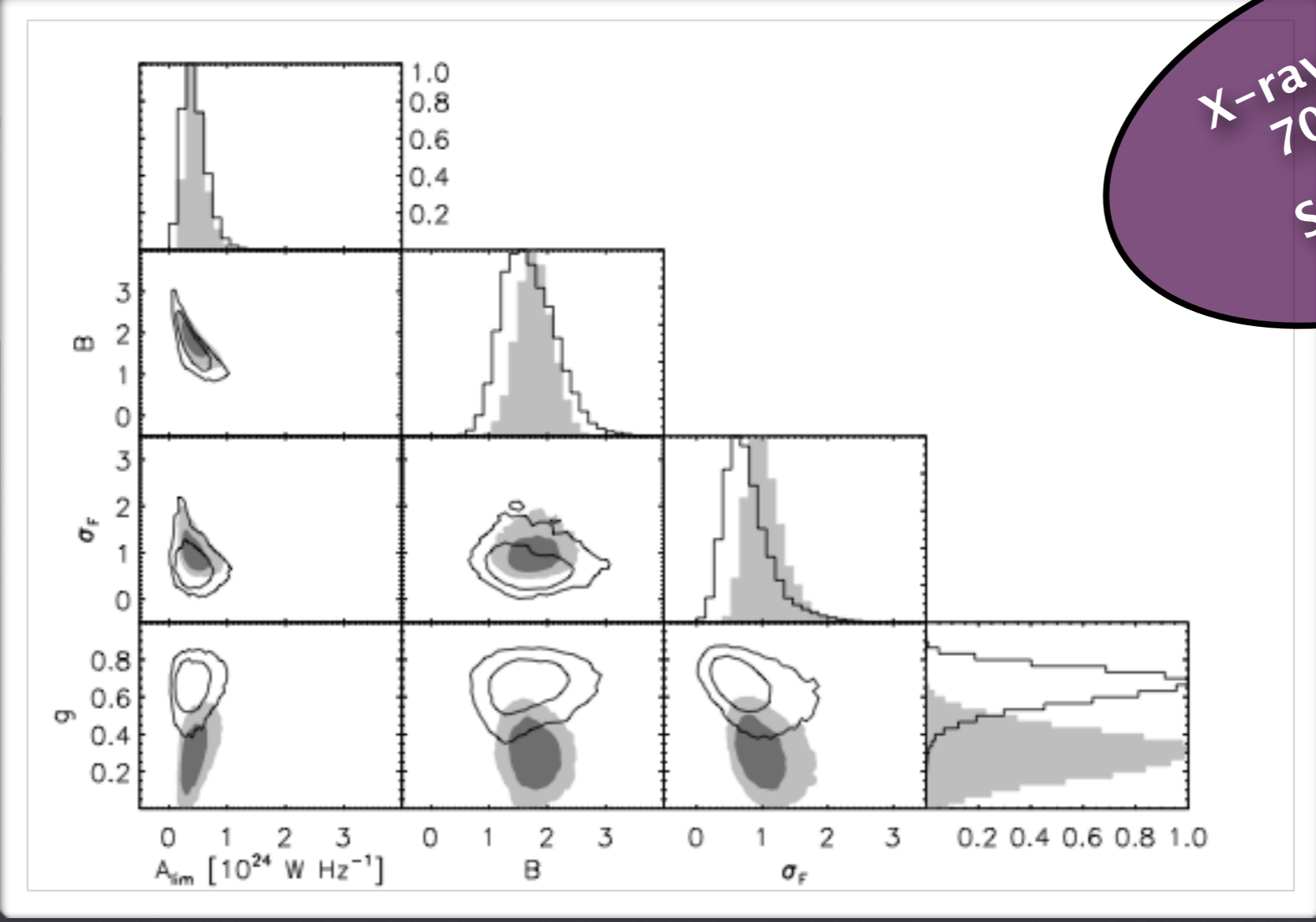
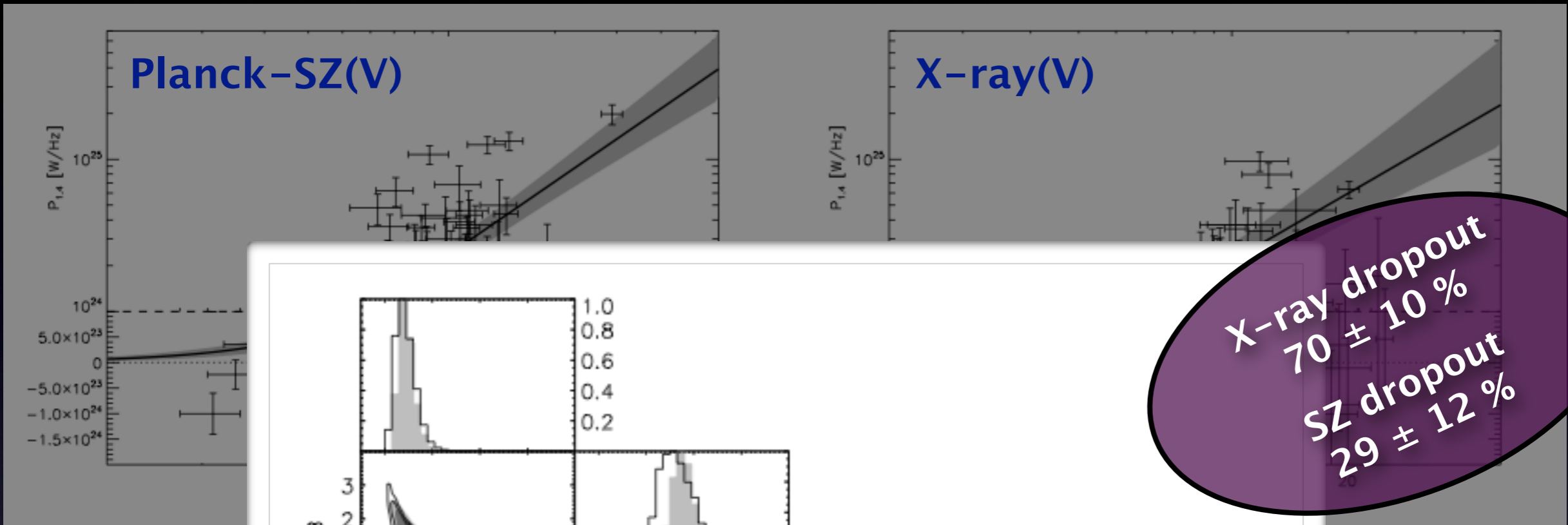


We fit simultaneously for an “on-correlation” population and a “zero” population for both SZ and X-ray sub-samples

The “on-correlation” populations give consistent mass scaling, with large scatter

But the zero-populations are significantly different!

SZ vs X-ray selection

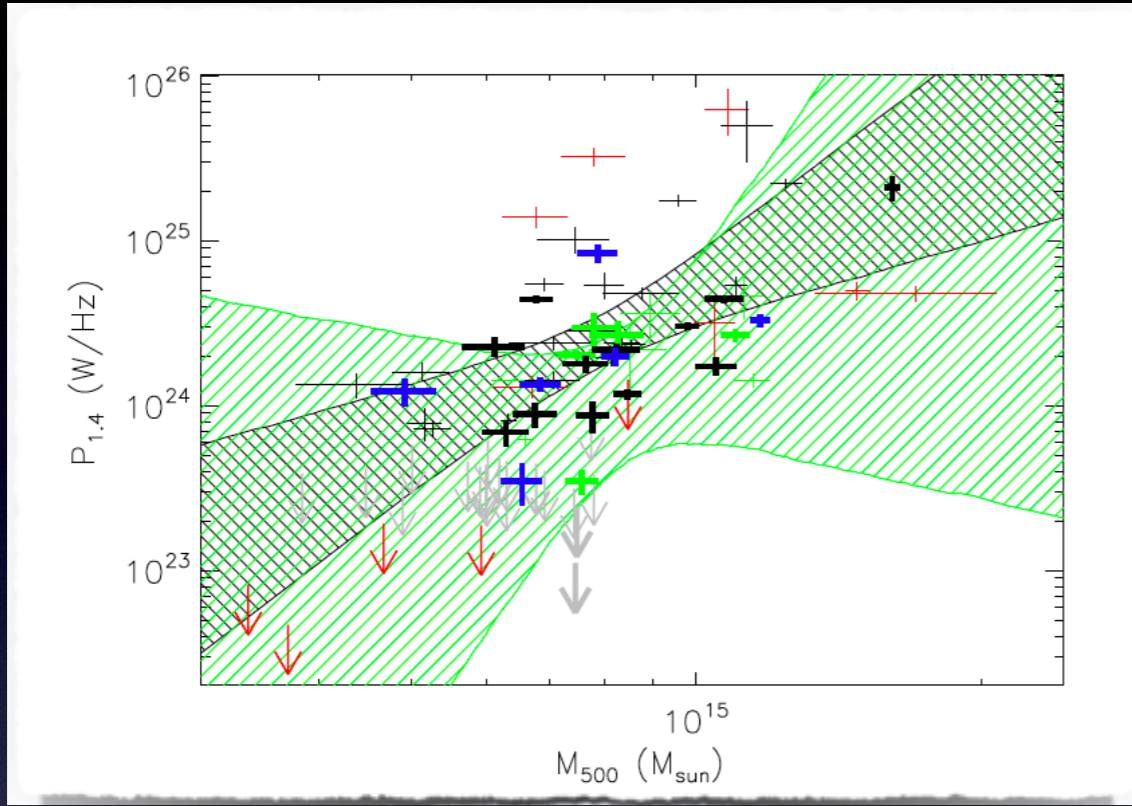


X-ray dropout
 $70 \pm 10 \%$
 SZ dropout
 $29 \pm 12 \%$

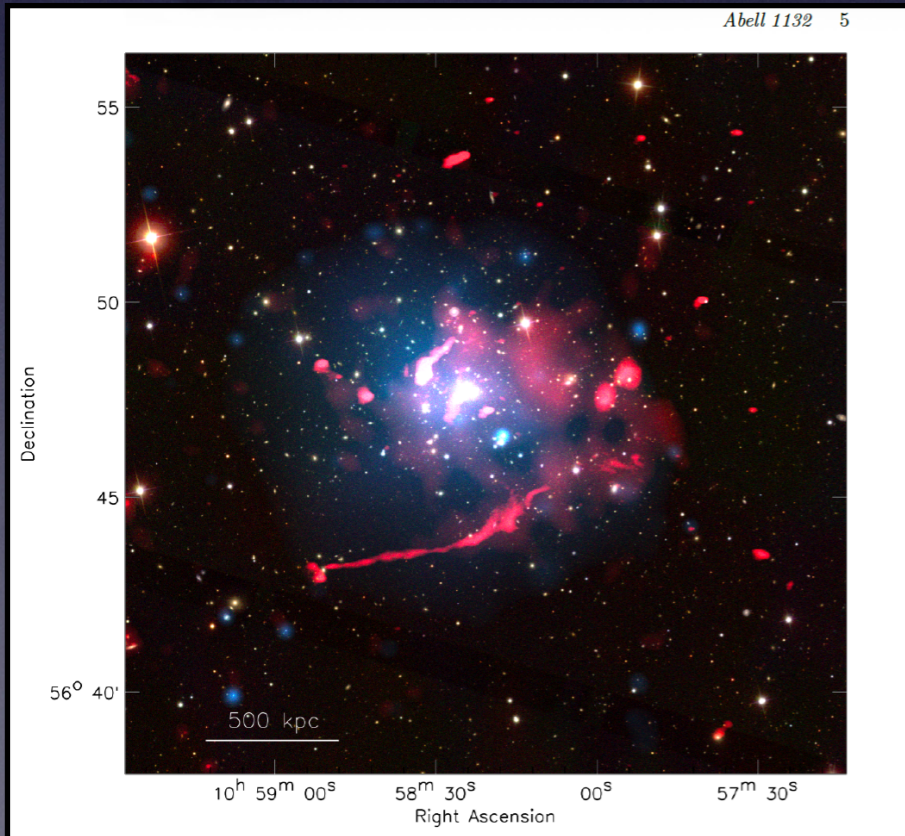
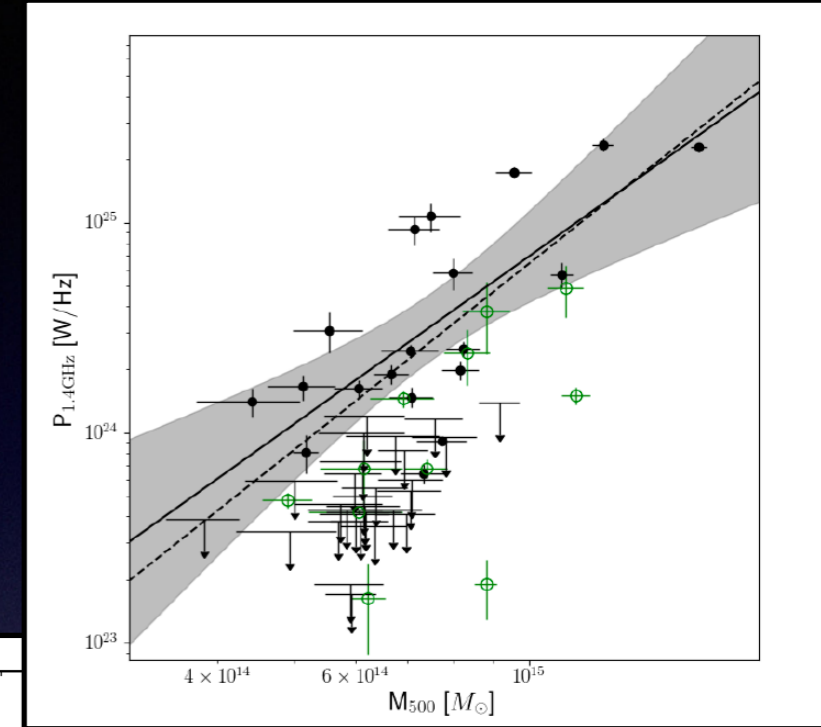
Sommer & Basu 2014

State-of-the-art for radio halos

Sommer, Basu et al. (in prep.)

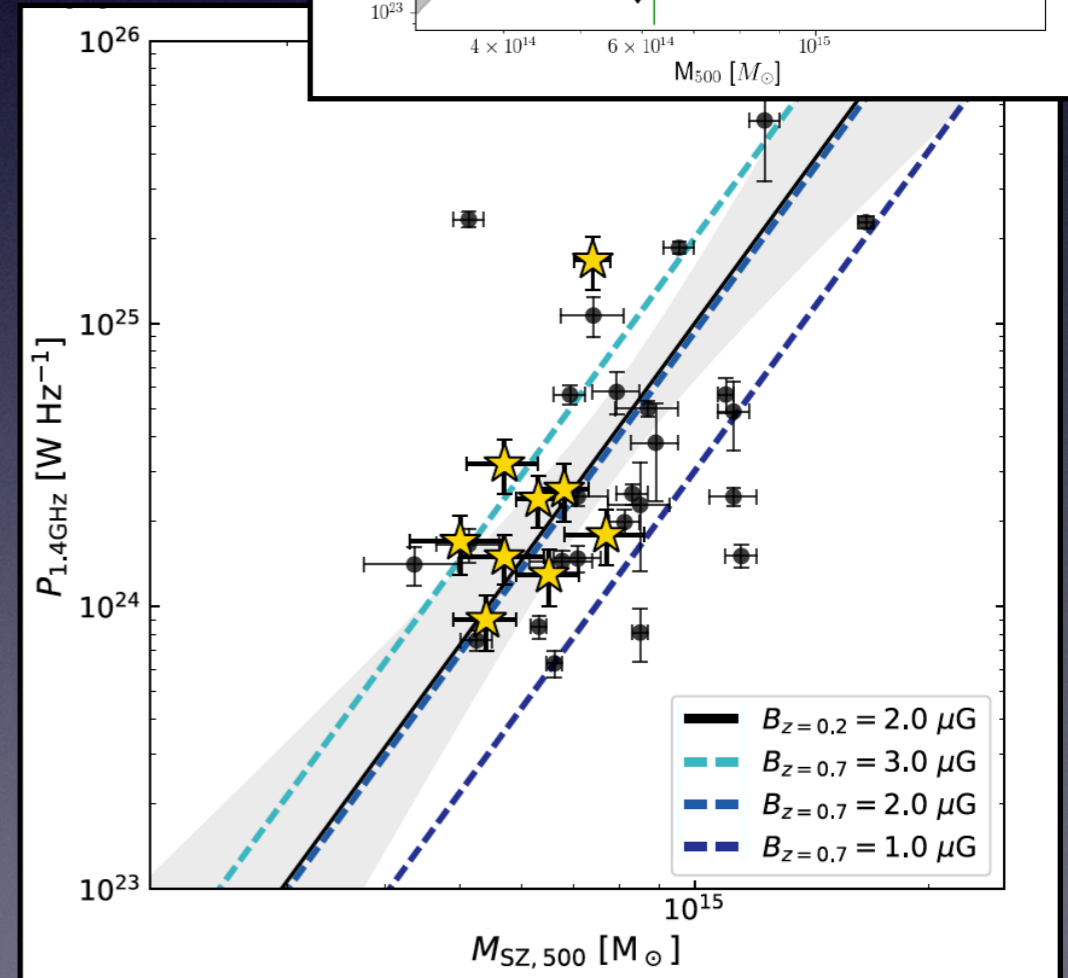


Radio halos in *Planck* selected & other clusters, Cuciti et al. (2021)



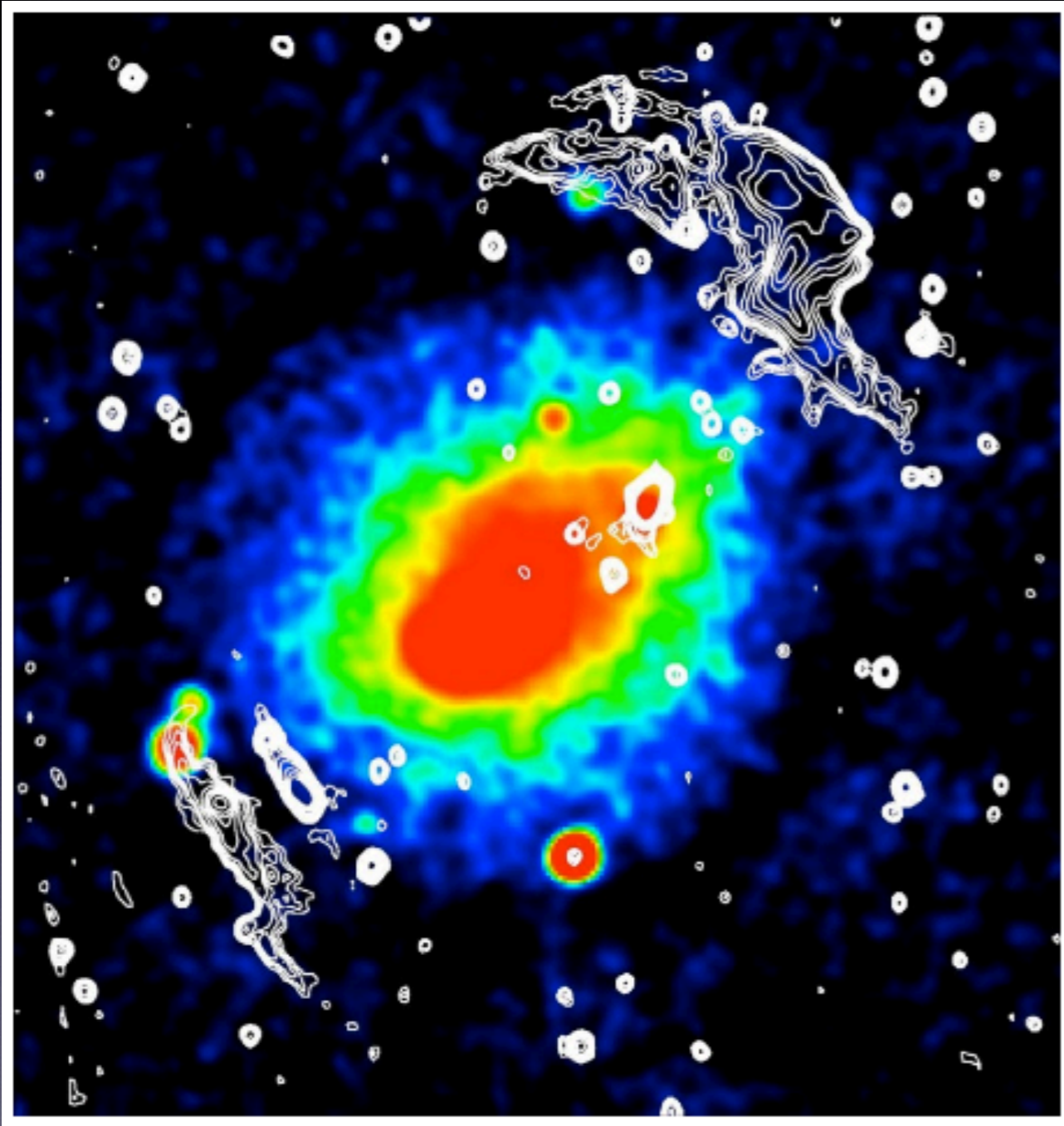
High-redshift ($z \sim 0.6-0.9$) radio halos from Di Gennaro et al. (2020)

LOFAR observation of radio halos, from Wilber et al. (2018)



Part II: Radio relics

Radio relics in clusters



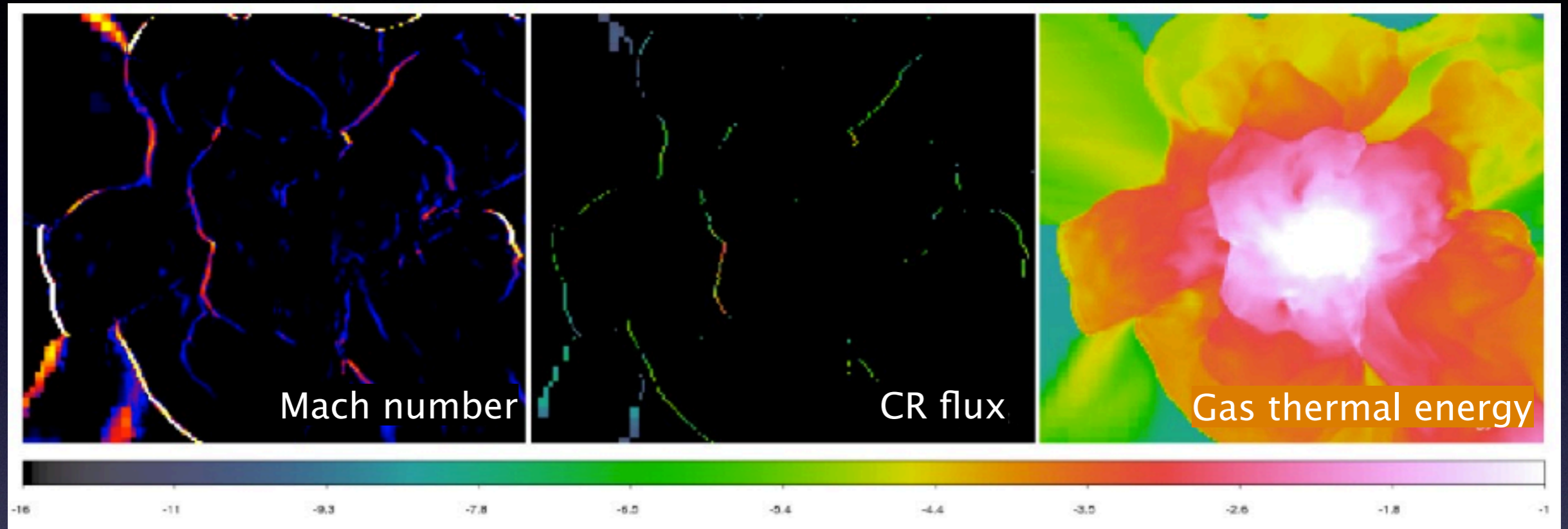
Radio relics: $L_{1.4 \text{ GHz}} \sim 10^{23-25} \text{ W/Hz}$

- Extended (up to $\sim 1 \text{ Mpc}$) diffuse radio sources at the periphery of clusters
- Irregular morphology
- High degree of polarization
- Steep spectrum ($\alpha \sim 1.2$)
- No optical counterpart
- Morphology resembles shock fronts, found only in disturbed clusters

Abell 3667 (Röttgering et al. 1997)
Color: X-ray, contours: radio

Merger shocks & radio relics

Vazza et al. (2012)

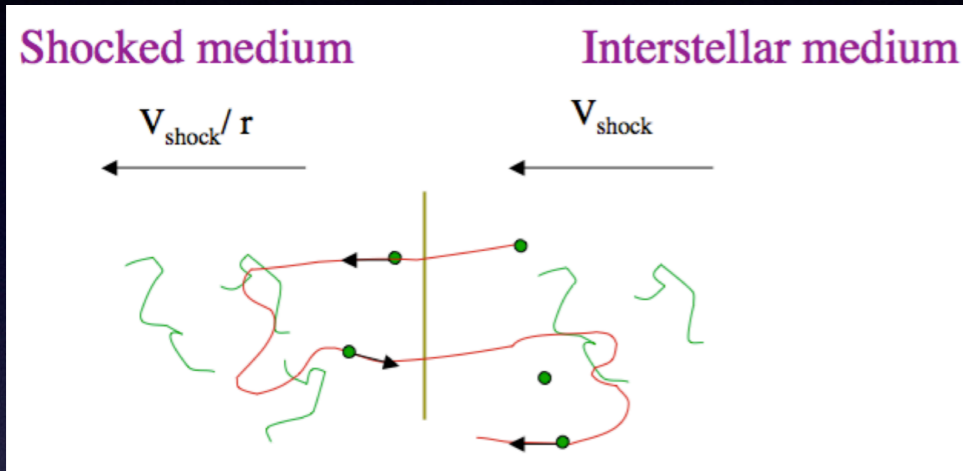


Radio relics are thought to be associated to cluster merger shocks. The shock fronts accelerate electrons (and also protons) with the Fermi-acceleration mechanism, and also compresses the magnetic fields. Those GeV electrons spiraling in the magnetic fields give rise to the synchrotron emission.

- Merger shocks have low Mach numbers ($M \sim 2-4$), so acceleration efficiency will be low
- Simulations predict many shock fronts, but only a few relics are known. Also, most of the relics do not have a detected shock feature.

“Problem” with radio relics

Diffusive Shock Acceleration (DSA) provides the correct CRe energy spectrum, but it is highly inefficient, needs “seed electrons” with relativistic energies



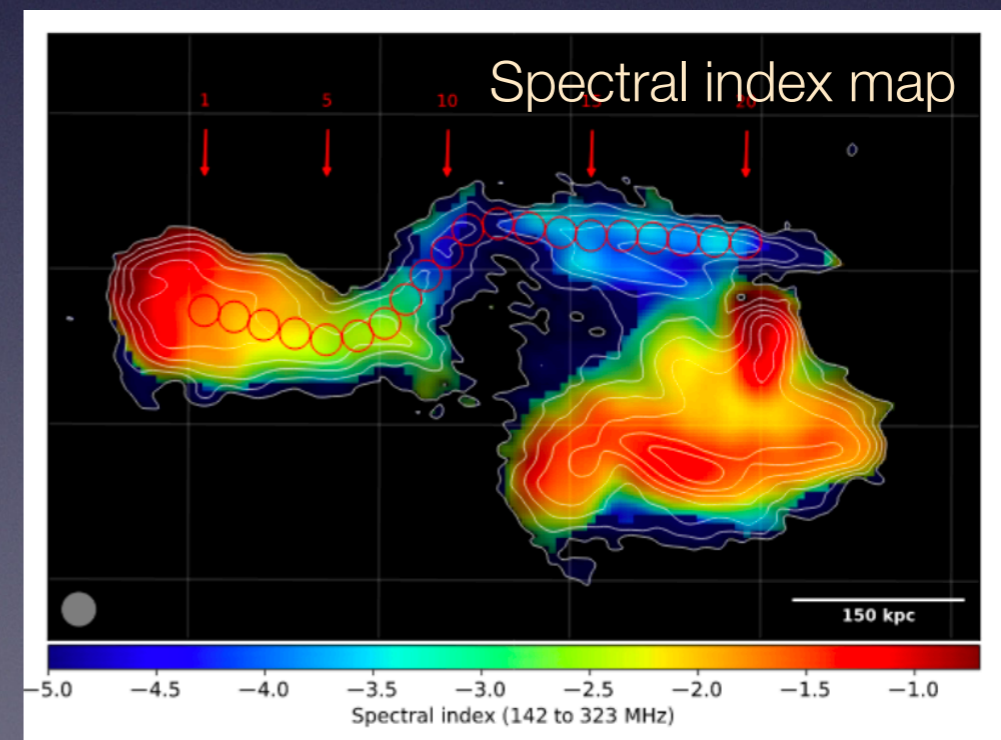
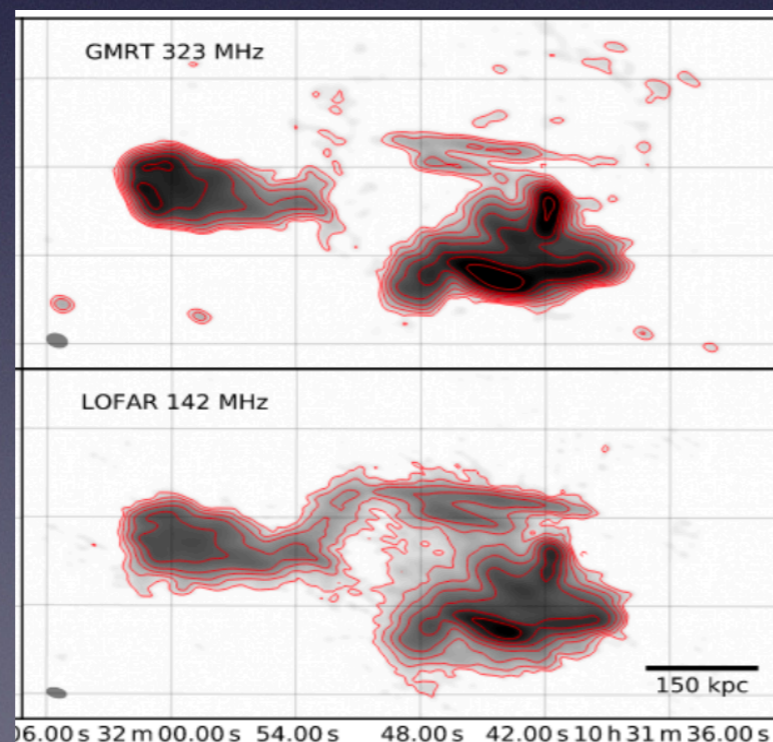
$$\left\langle \frac{\Delta E}{E} \right\rangle = \frac{4}{3}\beta = \frac{4}{3}\beta_{sh} \left(\frac{r-1}{r} \right)$$

- Linear dependence on shock Mach number
- Near-universal power law

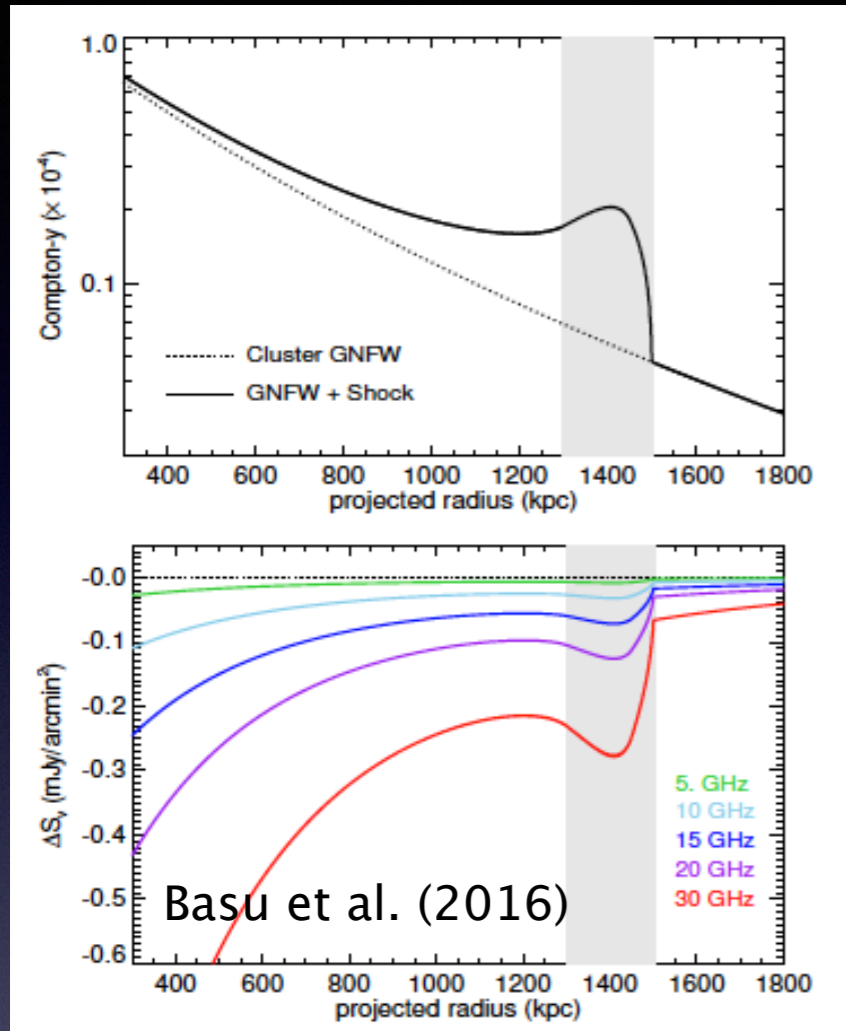
$$n(E) = \left| \frac{dN(\geq E)}{dE} \right| = (x-1) \frac{N_0}{E_0} \left(\frac{E}{E_0} \right)^{-x}$$

$$x = \frac{r+2}{r-1}$$

LOFAR observation of Abell 1033 cluster’s radio relic, from De Gasperin et al. (2017)



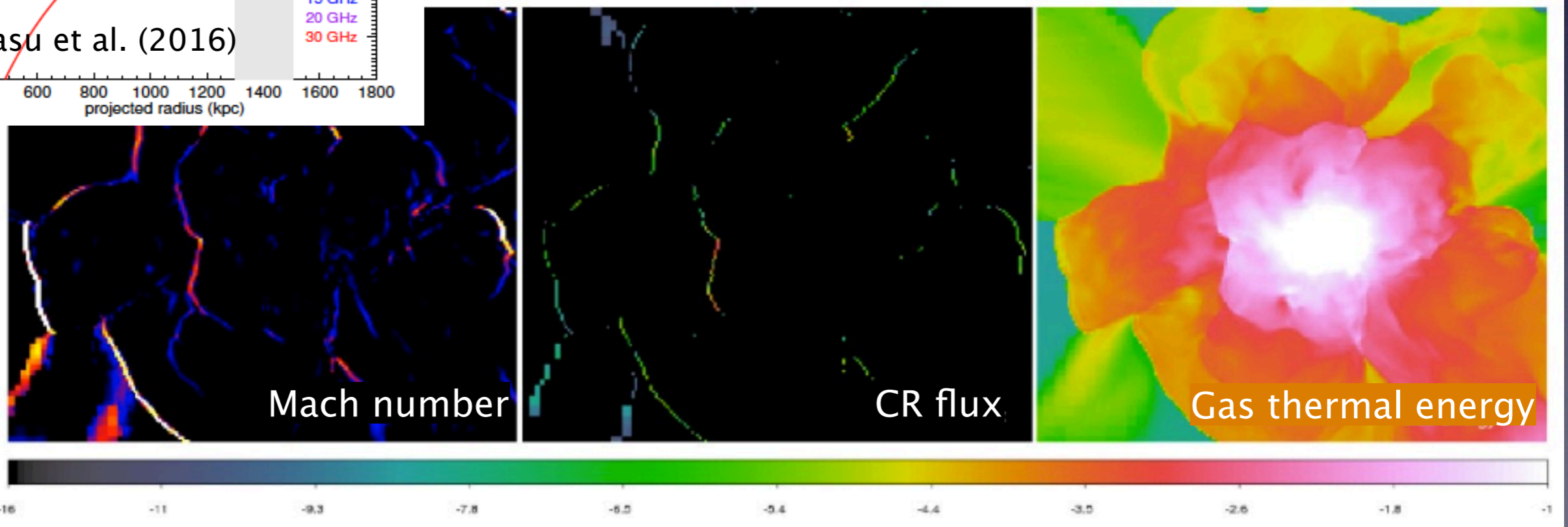
SZ signal expected from cluster shocks



Shocks create a pressure boost, which roughly scales as the **Mach number squared**

On projection ($\int P dl = \text{Compton } y \text{ parameter}$) this looks like a step function. This should be relatively easily detectable also in the cluster outskirts, and also out to high- z

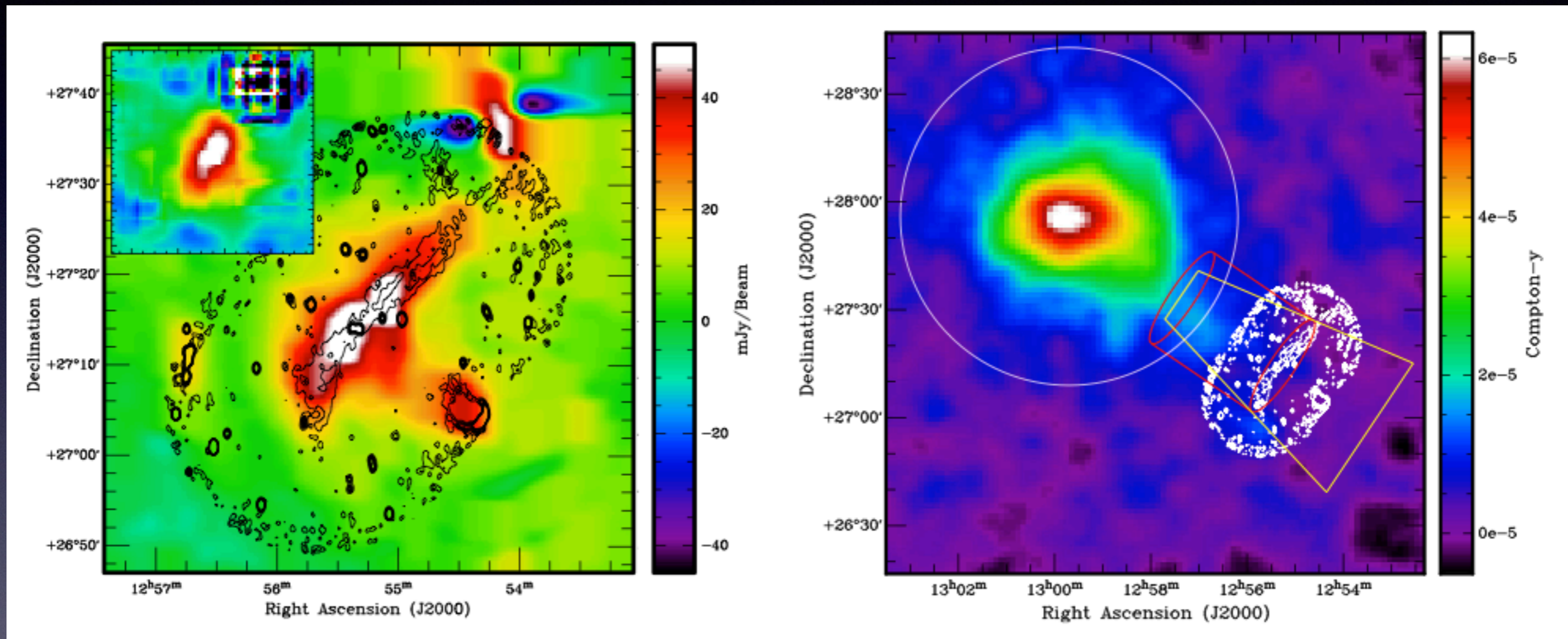
Vazza et al. (2012)



Coma's radio relic with Planck

Coma relic has already been analyzed in X-rays:
Akamatsu et al. (2013), Ogrean & Brüggen (2013)

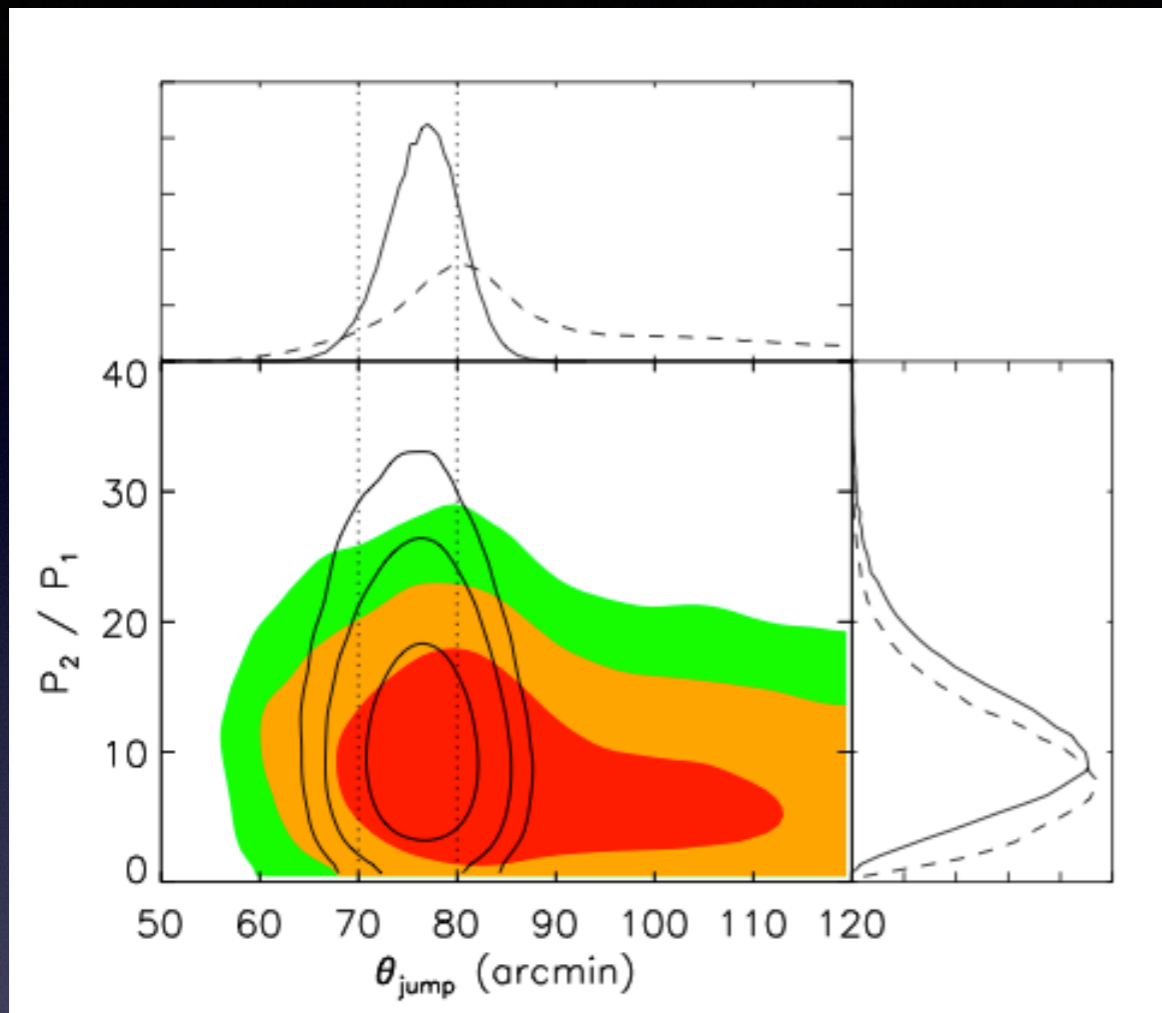
Erler, Basu et al. (2015), MNRAS, 447, 2497



We used new 2.4 GHz radio data for the coma relic, and extracted our own y-map from the Planck 2013 public data release

Results for Coma's relic shock

Erlar, Basu et al. (2015)



- SZ data favors a jump close to the relic without any radio prior, at 79^{+10}_{-9} arcmin (radio relic at 75 arcmin)
- Corresponding pressure ratio at the relic is $8.8^{+6.1}_{-3.4}$
- Pressure ratio and jump location are uncorrelated

$$\frac{P_2}{P_1} = \frac{2\gamma\mathcal{M}^2 - \gamma + 1}{\gamma + 1}$$

$$\mathcal{M} = 2.8^{+0.8}_{-0.6}$$

This is the first “detection” of a pressure discontinuity at a radio relic with the SZ effect. This also happens to be the first SZ shock feature detected near a cluster’s virial radius.

With the latest 2015 Planck data release, we got Mach number $M = 3.4 \pm 0.5$ (pressure ratio $P_2/P_1 = 14.3 \pm 4.5$)

The finest imager for SZ shocks

Measuring SZ shocks with Planck is like observing radio sky with a single-dish antenna ... **we can do better**

Projected pressure map
 $M_{\text{vir}} \sim 2 \times 10^{14}$ merger



(Simulations by F. Vazza, 2012)

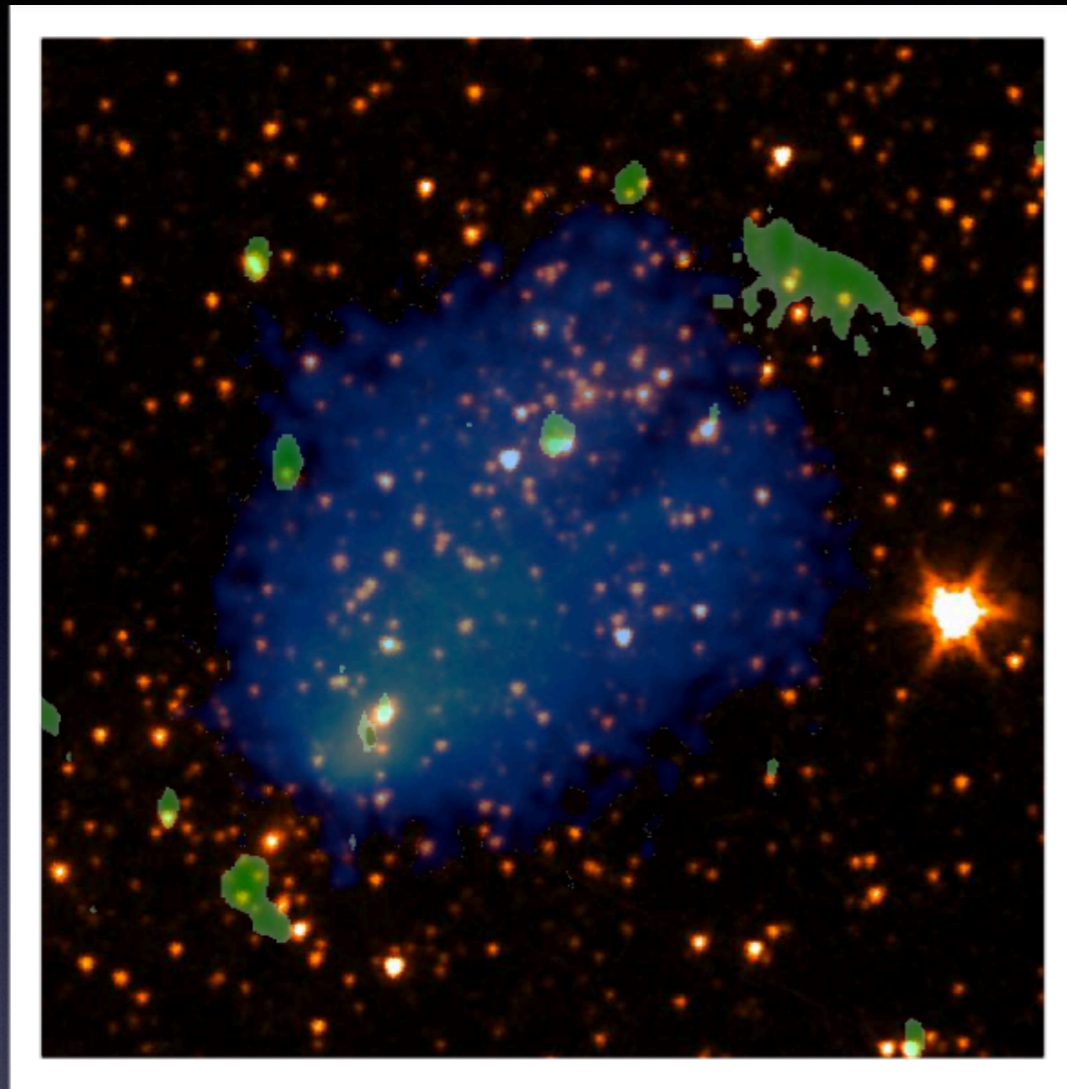


First ALMA-SZ results:

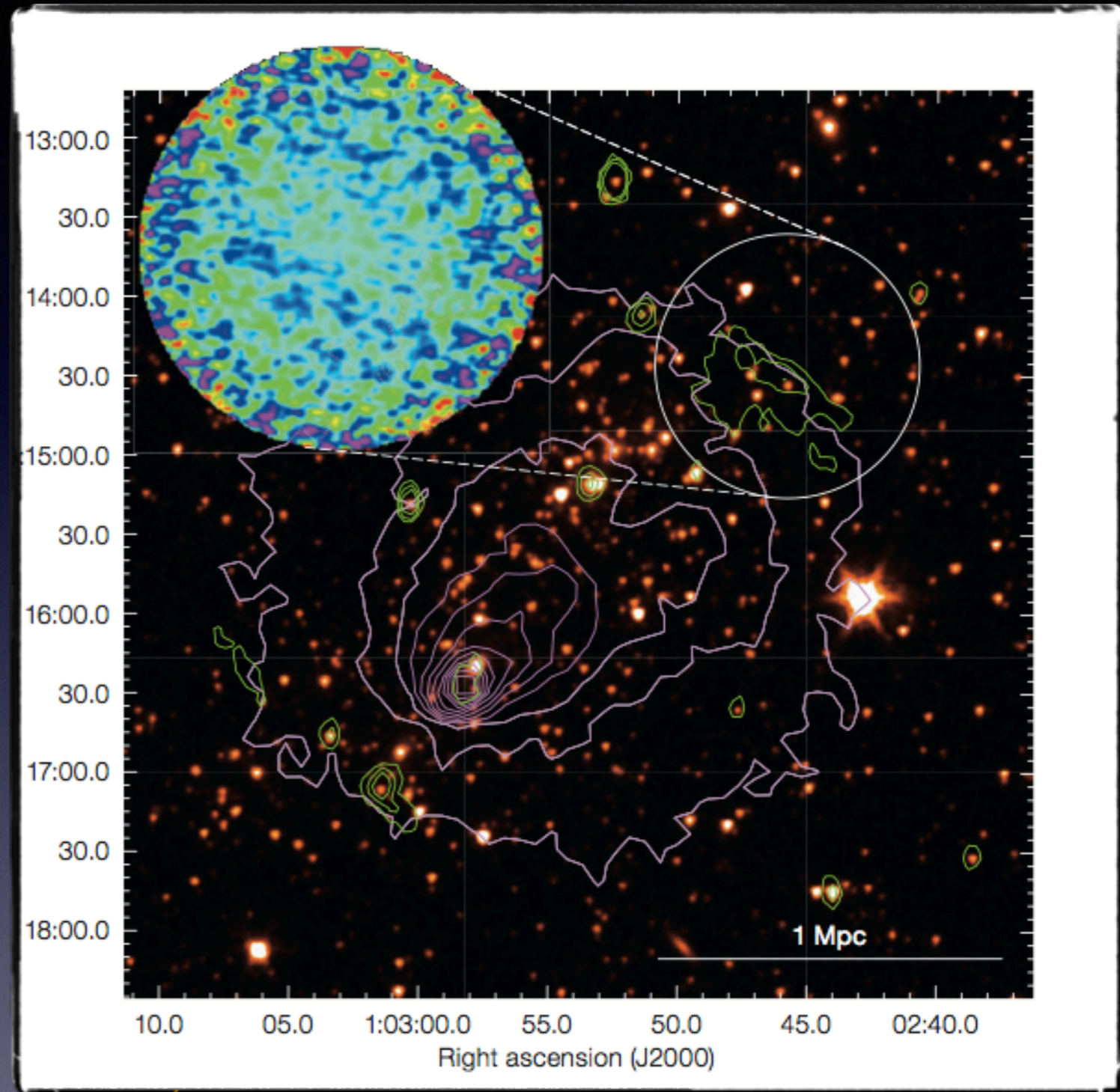
- ★ RXC J1347.5 core (Kitayama et al. 2016)
- ★ El Gordo relic shock (Basu et al. 2016)

High-resolution single dish measurements are also on the way.

A relic shock from at $z=0.9$



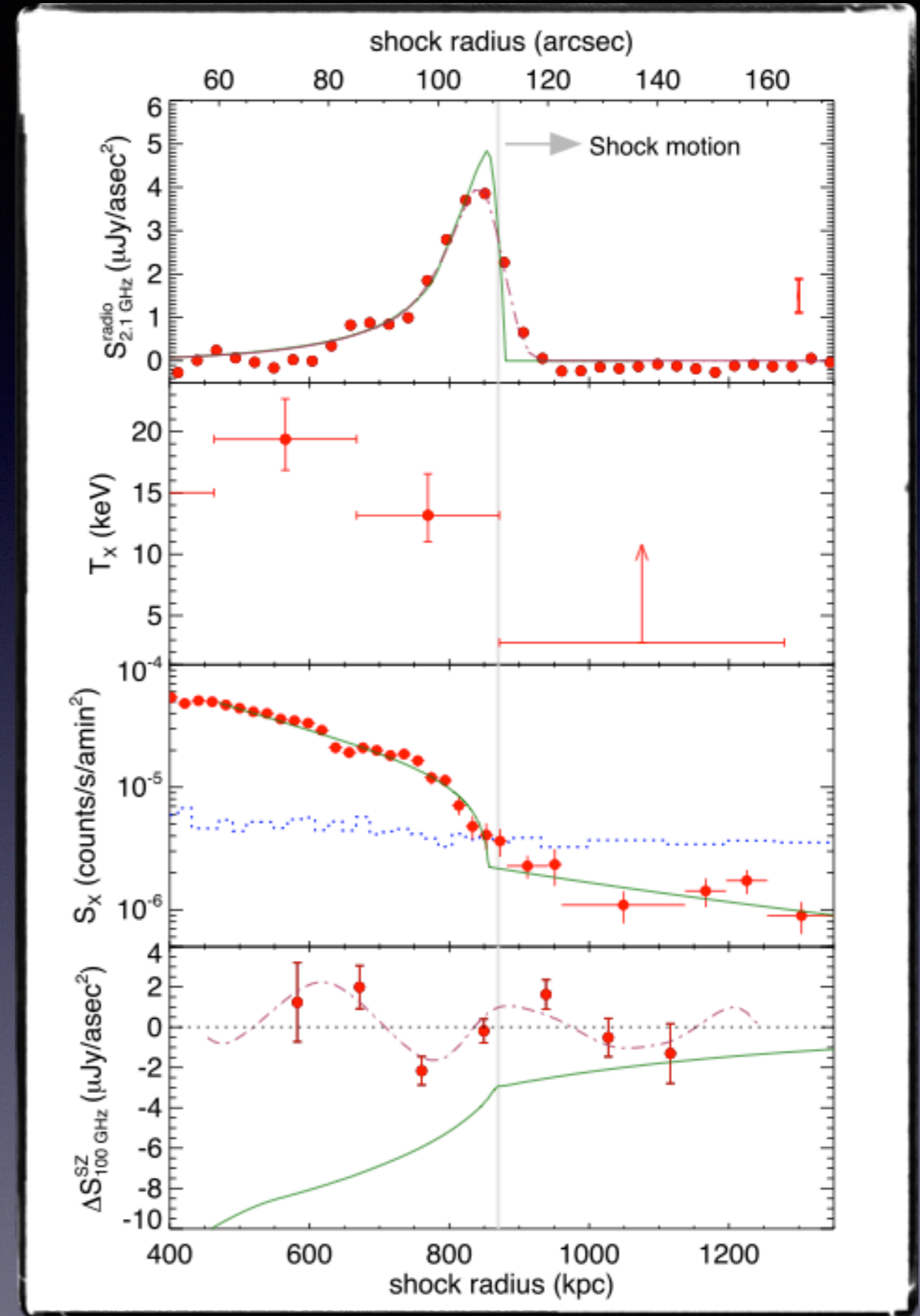
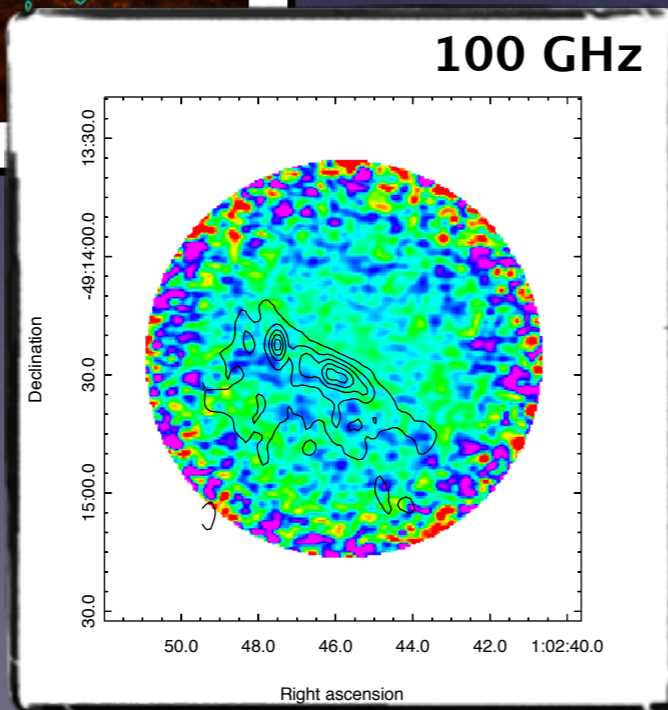
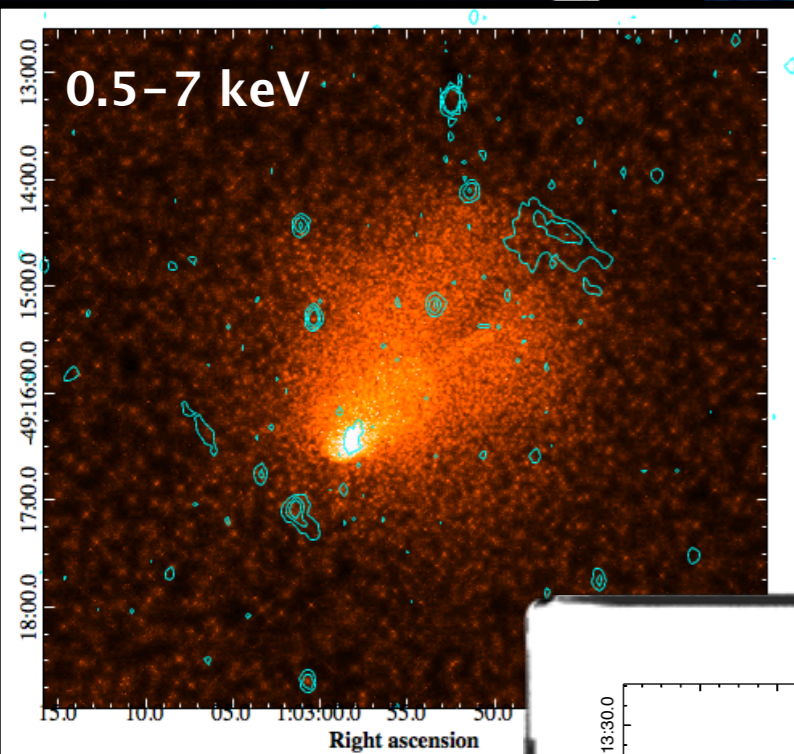
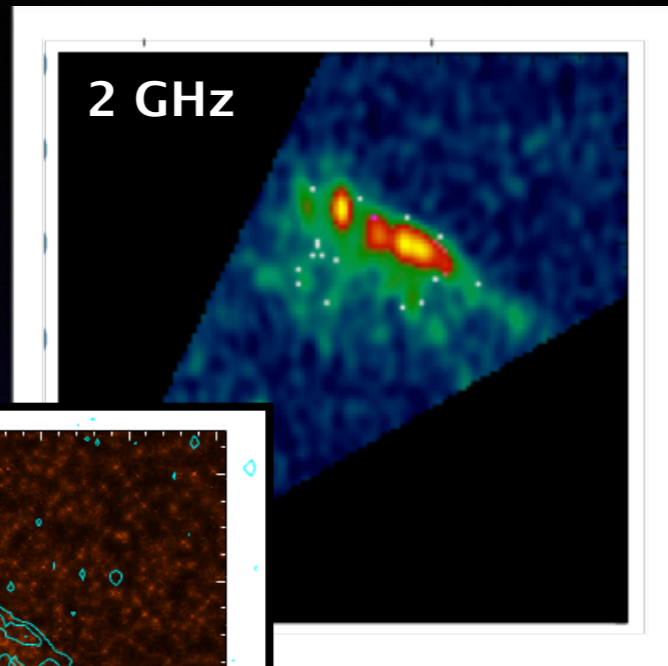
360 ks Chandra + ATCA 2.1 GHz radio
(PI: J. Hughes) (Lindner et al. 2014)



Basu et al. (2016), ApJ, 829

ALMA data ~ 2h on-source
ALMA noise rms ~ $6 \mu\text{Jy}/3''$ beam
(enough to detect $M \sim 2$ shock with $> 5\sigma$)

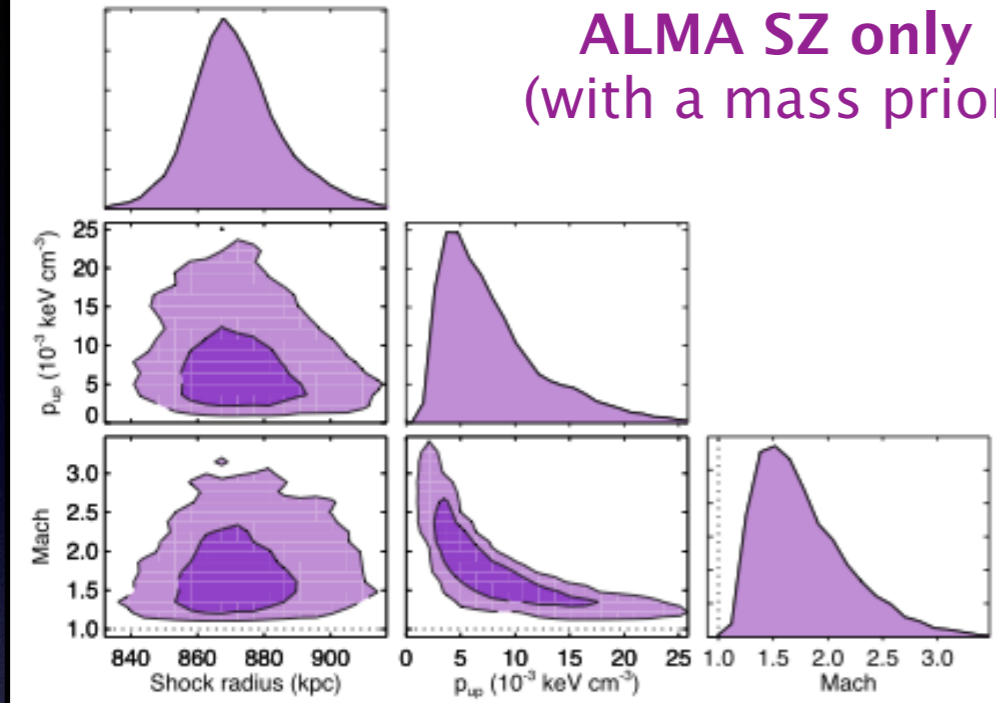
The multi-wavelength view



Basu et al. (2016), ApJ, 829

SZ /X-ray joint modelling of relic shock

ALMA SZ only
(with a mass prior)

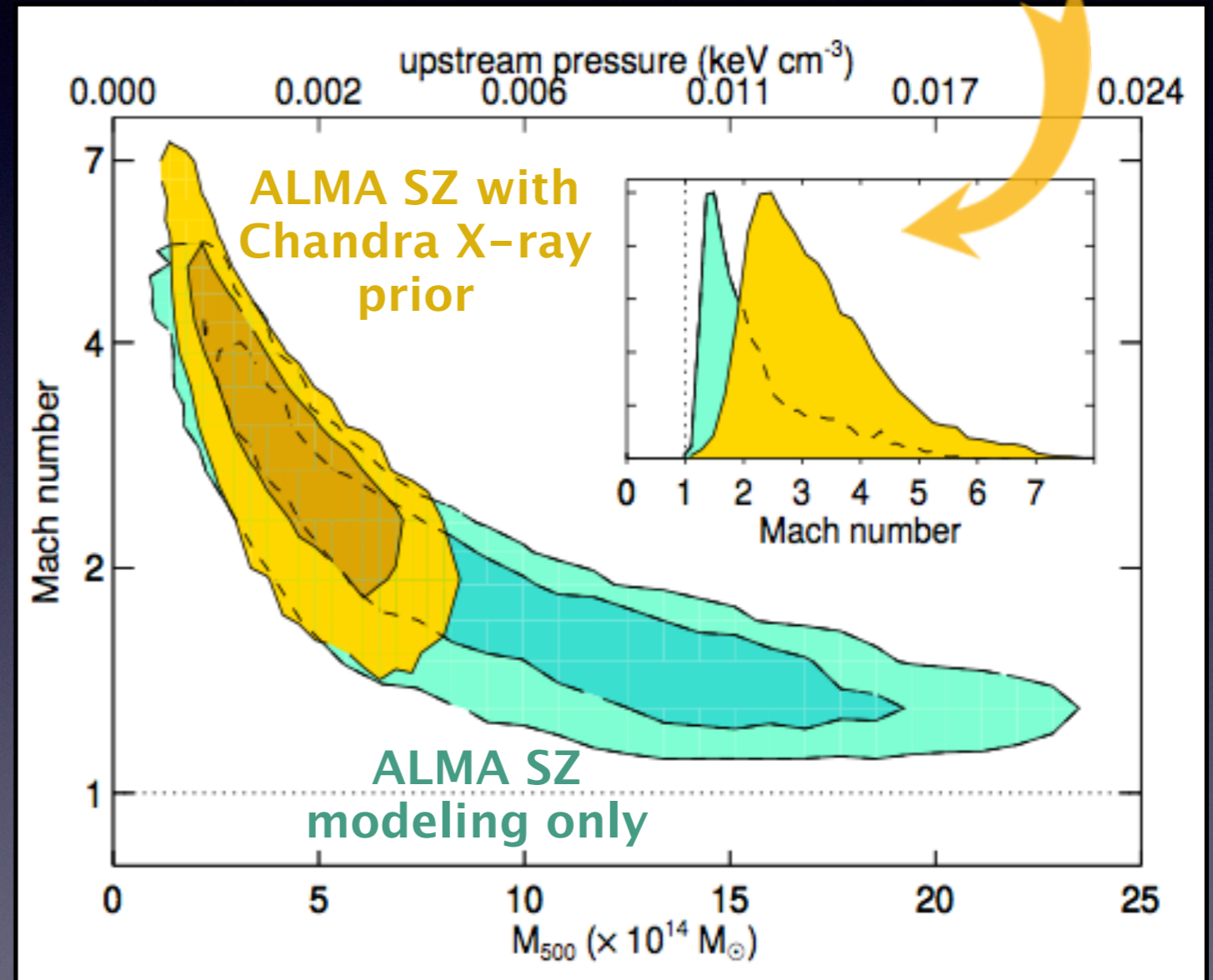
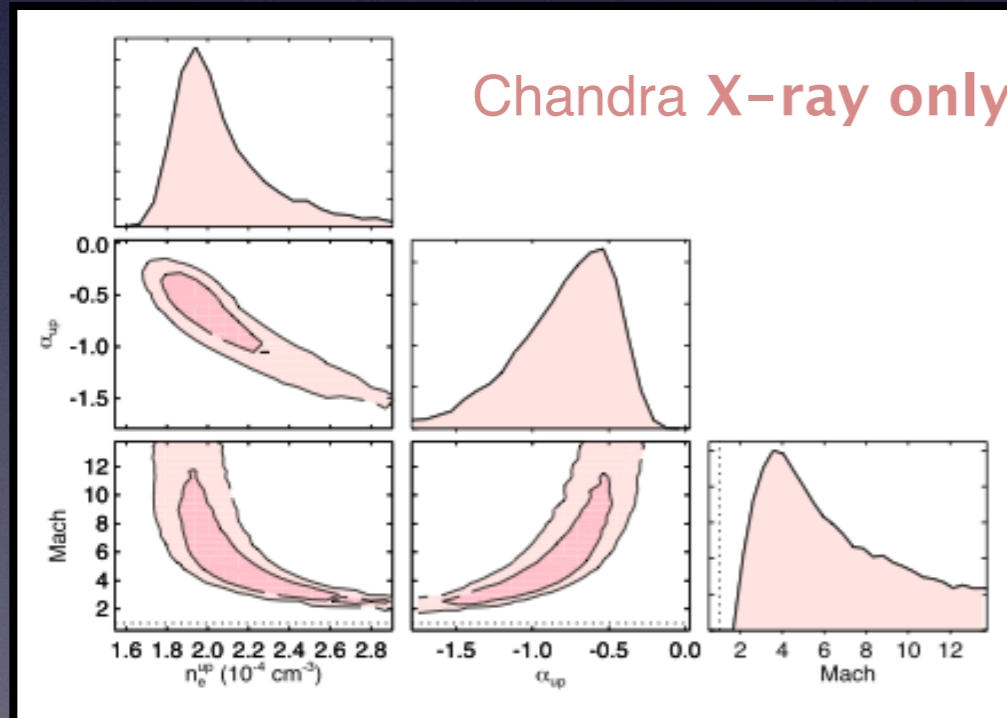


ALMA SZ data alone points to a weak shock: $\mathcal{M} = 1.4^{+1.2}_{-0.2}$

X-ray brightness jump suggests stronger: $\mathcal{M} = 3.5^{+6.4}_{-1.3}$

We use an X-ray pressure prior on the SZ modeling.

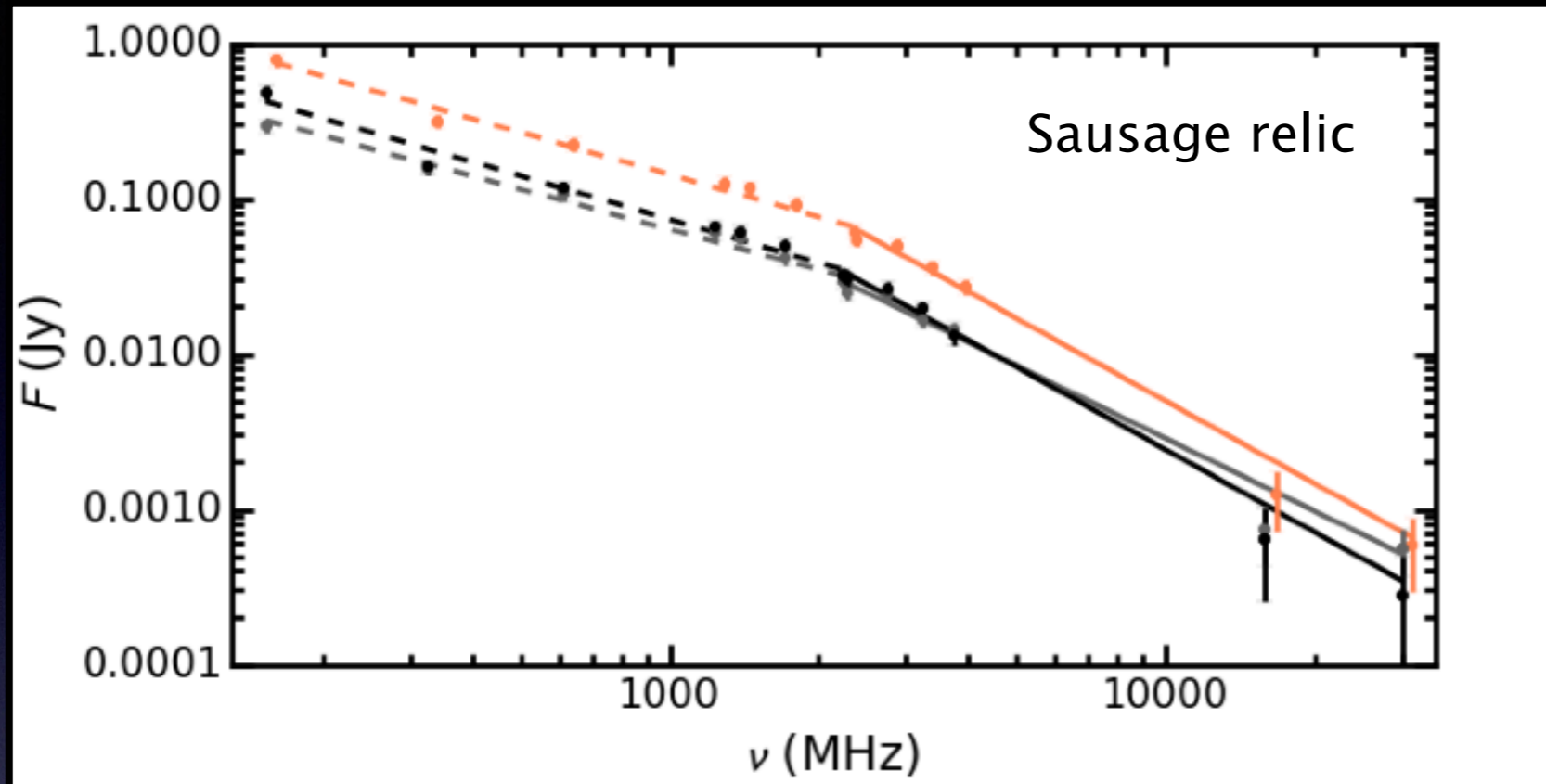
Chandra X-ray only



Basu et al. (2016), ApJ, 829

“SZ contamination” on relic shocks

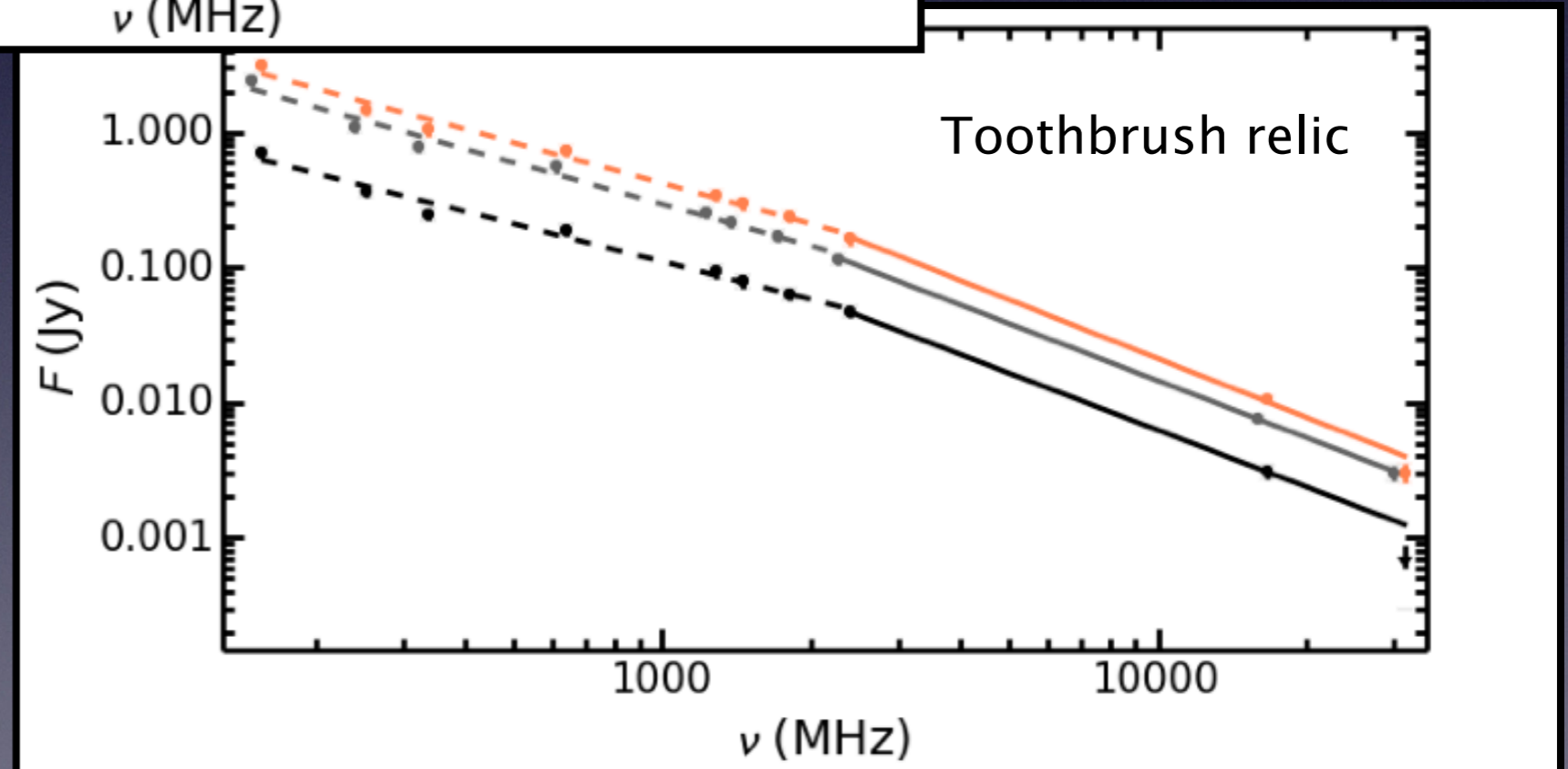
Relic spectral steepening at >10 GHz



Stroe et al. (2015, 2016)

AMI (16 GHz) and
CARMA (30 GHz) data

A gradual spectral steepening is observed above ~ 2 GHz, which cannot be explained from the standard DSA model.



Relic spectral steepening at >10 GHz

RE-ACCELERATION MODEL FOR RADIO RELICS WITH SPECTRAL CURVATURE

HYESUNG KANG¹ AND DONGSU RYU^{2,3,4}

¹ Department of Earth Sciences, Pusan National University, Pusan 46241, Korea; hskang@pusan.ac.kr

² Department of Physics, UNIST, Ulsan 44010, Korea; ryu@unist.ac.kr

Turbulent Cosmic-Ray Reacceleration and the Curved Radio Spectrum of the Radio Relic in the Sausage Cluster

Yutaka FUJITA¹, Hiroki AKAMATSU,² and Shigeo S. KIMURA³

Magnetic Field Evolution in Giant Radio Relics using the example of CIZA J2242.8+5301

J. M. F. Donnert^{1,2,3*}, A. Stroe^{4,1†}, G. Brunetti², D. Hoang¹, H. Roettgering¹

¹ Leiden Observatory, Leiden University, P.O. Box 9513, NL-2300 RA Leiden, The Netherlands

² INAF-Istituto di Radioastronomia, via P. Gobetti 101, I-40129 Bologna, Italy

³ Max-Planck-Gesellschaft, Bonn, Germany

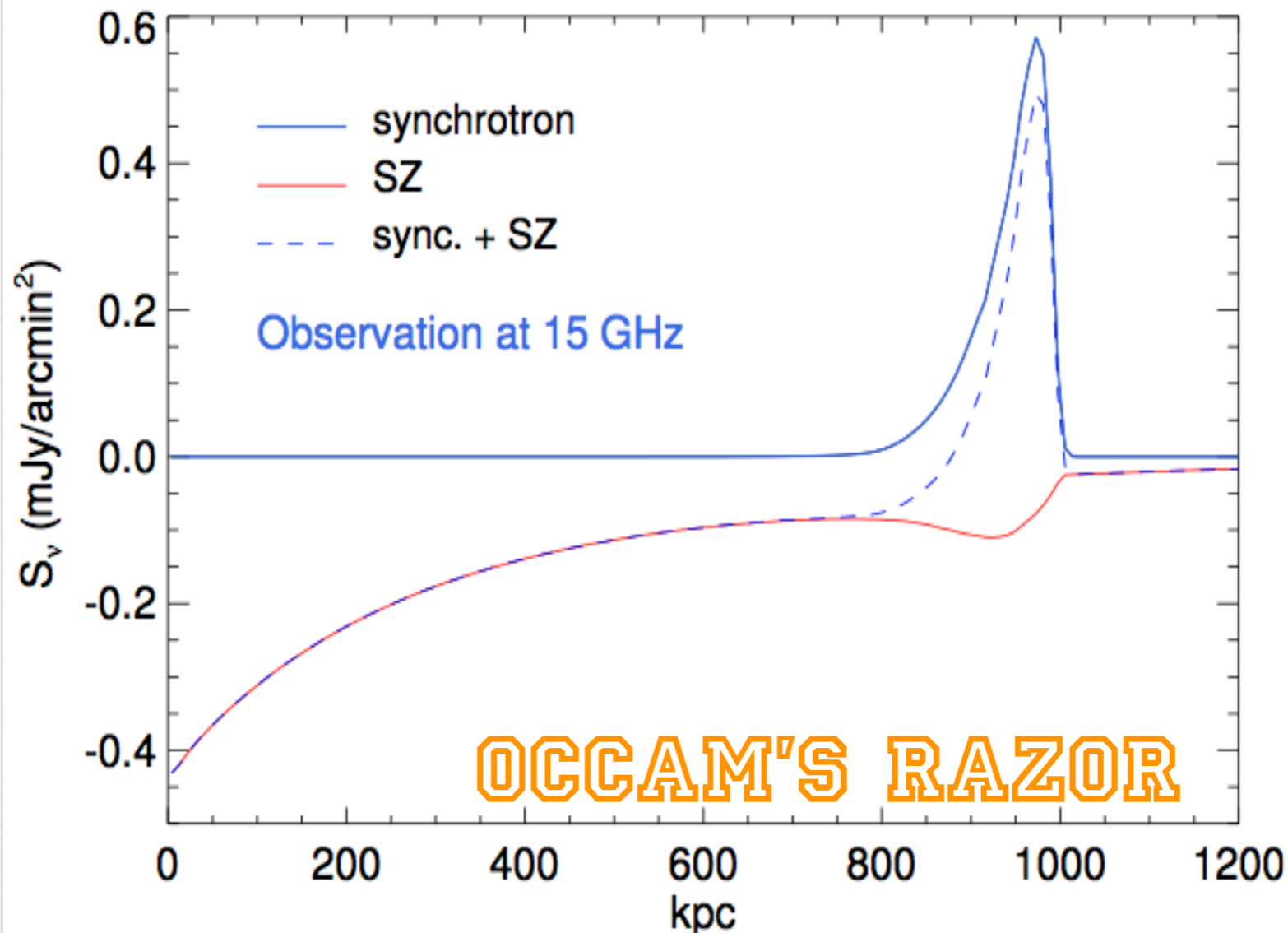
The widest frequency radio relic spectra: observations from 150 MHz to 30 GHz

Andra Stroe,^{1*†} Timothy Shimwell,¹ Clare Rumsey,² Reinout van Weeren,³
Maja Kierdorf,⁴ Julius Donnert,¹ Thomas W. Jones,⁵ Huub J. A. Röttgering,¹
Matthias Hoeft,⁶ Carmen Rodríguez-Gonzálvez,⁷ Jeremy I. Harwood⁸

016

Relic spectral steepening at >10 GHz

RE-ACCELERATION MODEL FOR RADIO RELICS WITH SPECTRAL CURVATURE



Basu et al. (2016), A&A, 591

150 MHz

Andra Stroe,^{1*}† Timothy Shimwell,¹ Clare Rumsey,² Reinout van Weeren,³
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016

Turbulent
Curved R
Sausage

Yutaka FUJITA

Magnetic Field
CIZA J2242.8+

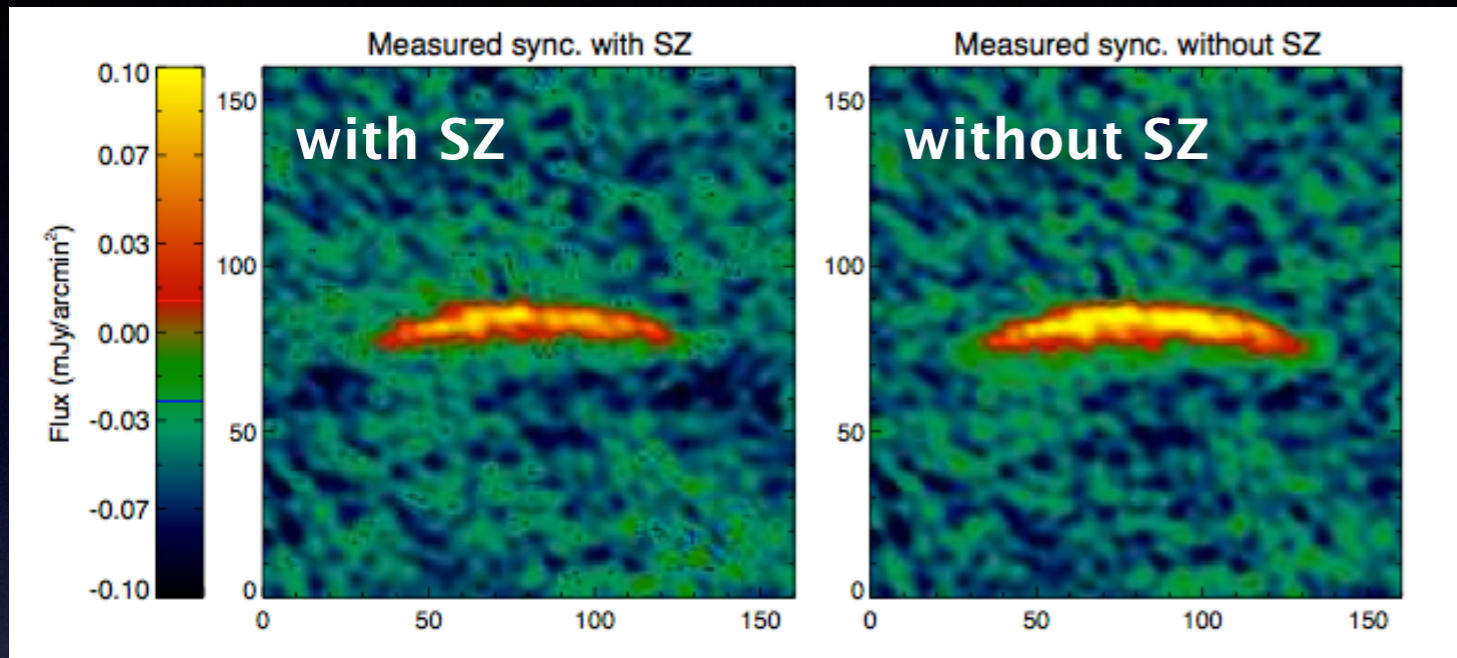
J. M. F. Donnert^{1,2,3}

¹ Leiden Observatory, Leiden Univ

² INAF-Istituto di Radioastronomia

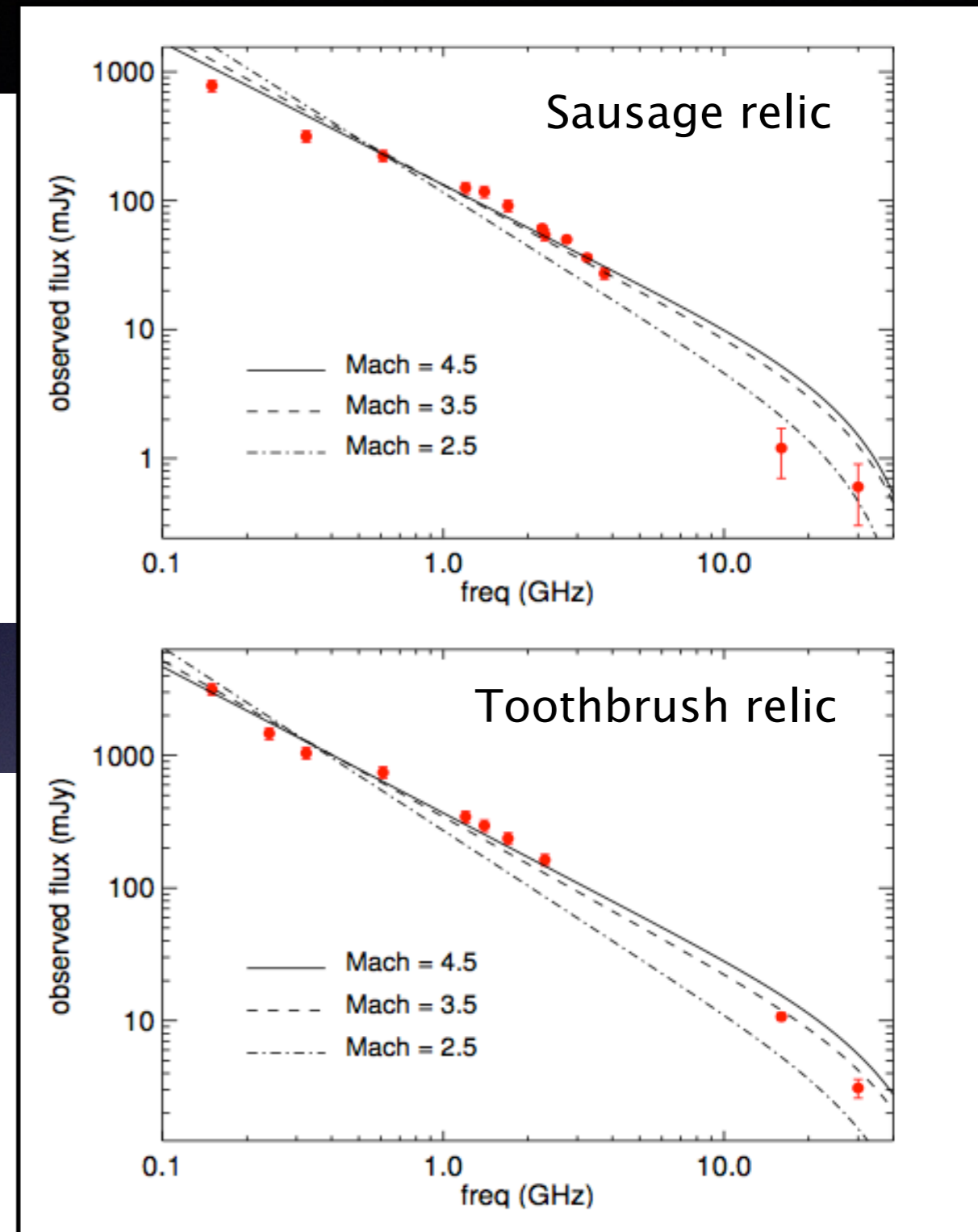
A non-negligible effect at >10 GHz

Simulated interferometric observation at 10 GHz



10%–50% flux loss at 10 GHz

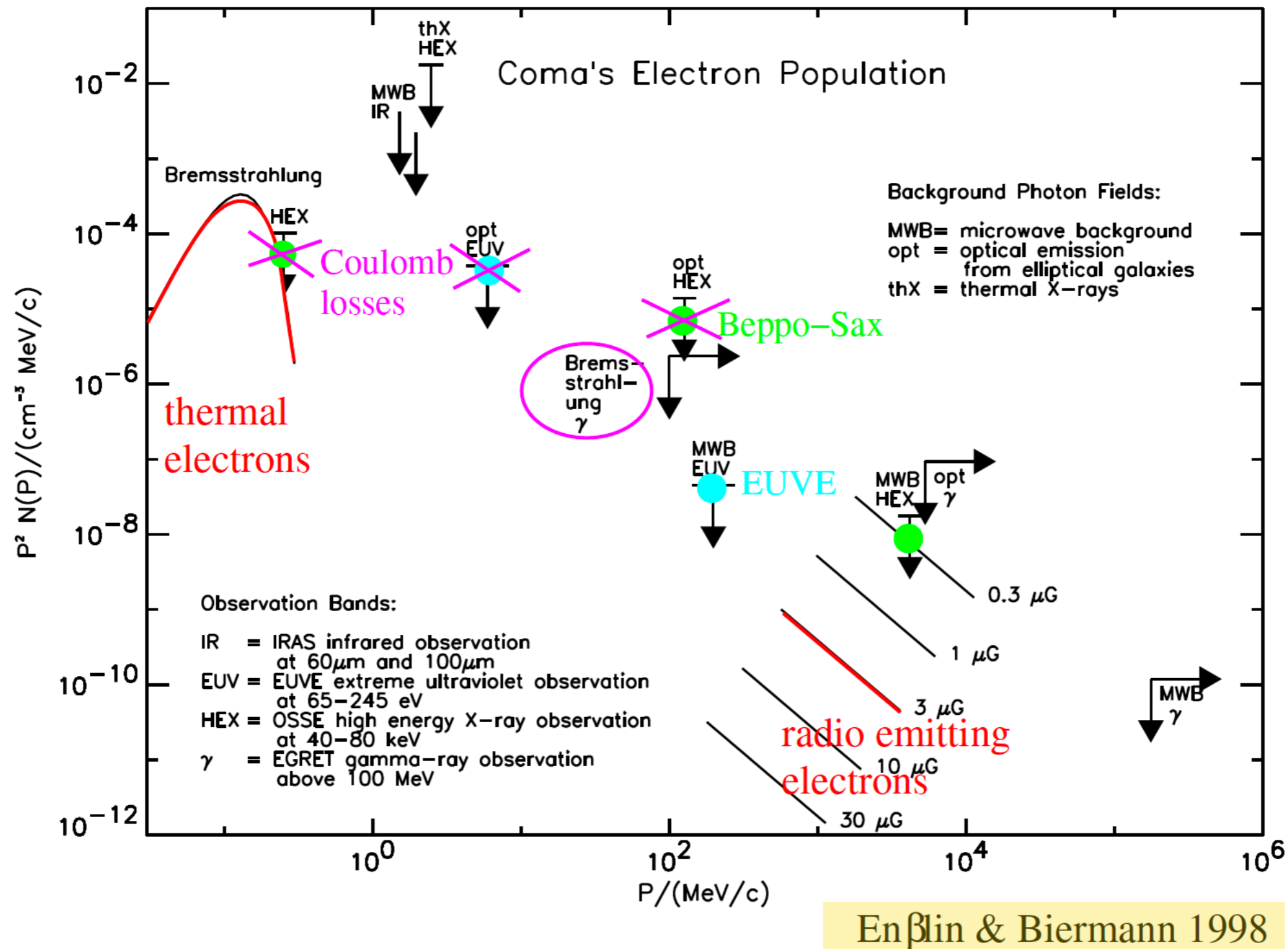
	3 GHz	5 GHz	10 GHz	15 GHz	20 GHz	30 GHz
Sausage relic ($\mathcal{M} = 2.5$)	<1%	<1%	4%	11%	24%	58%
($\mathcal{M} = 3.5$)	<1%	<1%	3%	10%	21%	49%
($\mathcal{M} = 4.5$)	<1%	<1%	4%	12%	24%	52%
Toothbrush relic ($\mathcal{M} = 3.5$)	<1%	<1%	3%	9%	18%	43%
($\mathcal{M} = 4.5$)	<1%	<1%	3%	10%	20%	46%
El Gordo relic ($\mathcal{M} = 2.5$)	<1%	3%	23%	53%	81%	>100%
A2256 relic ($\mathcal{M} = 2.0$)	1%	3%	28%	66%	96%	>100%



Basu et al. (2016), A&A, 591

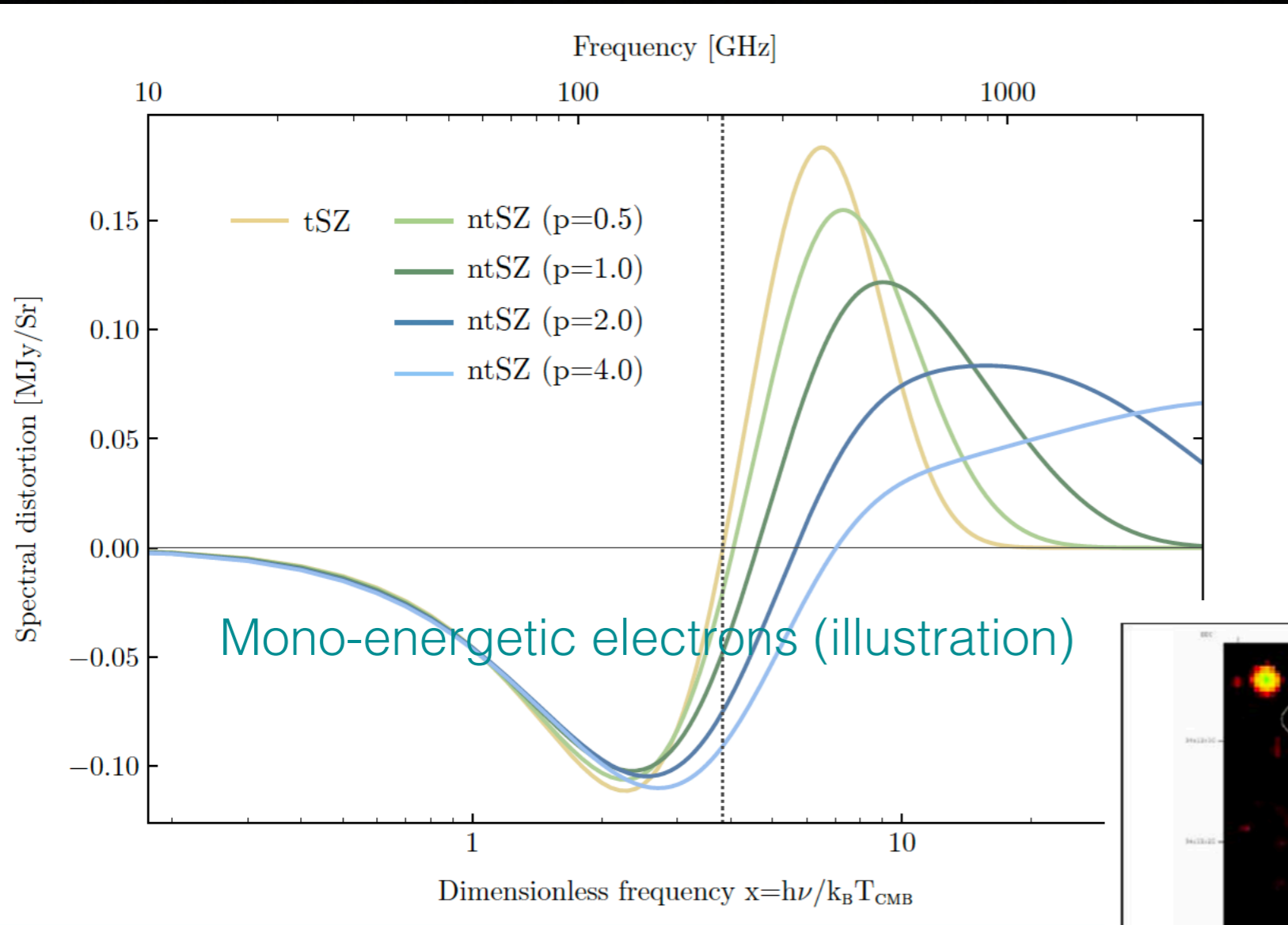
Part III: Nonthermal SZ

Nonthermal electrons in clusters



(reproduced in Enßlin 2004, with annotations)

Nonthermal SZ effect



CMB photons will also scatter off other sources of free electrons, e.g. power-law distribution with a high-energy tail.

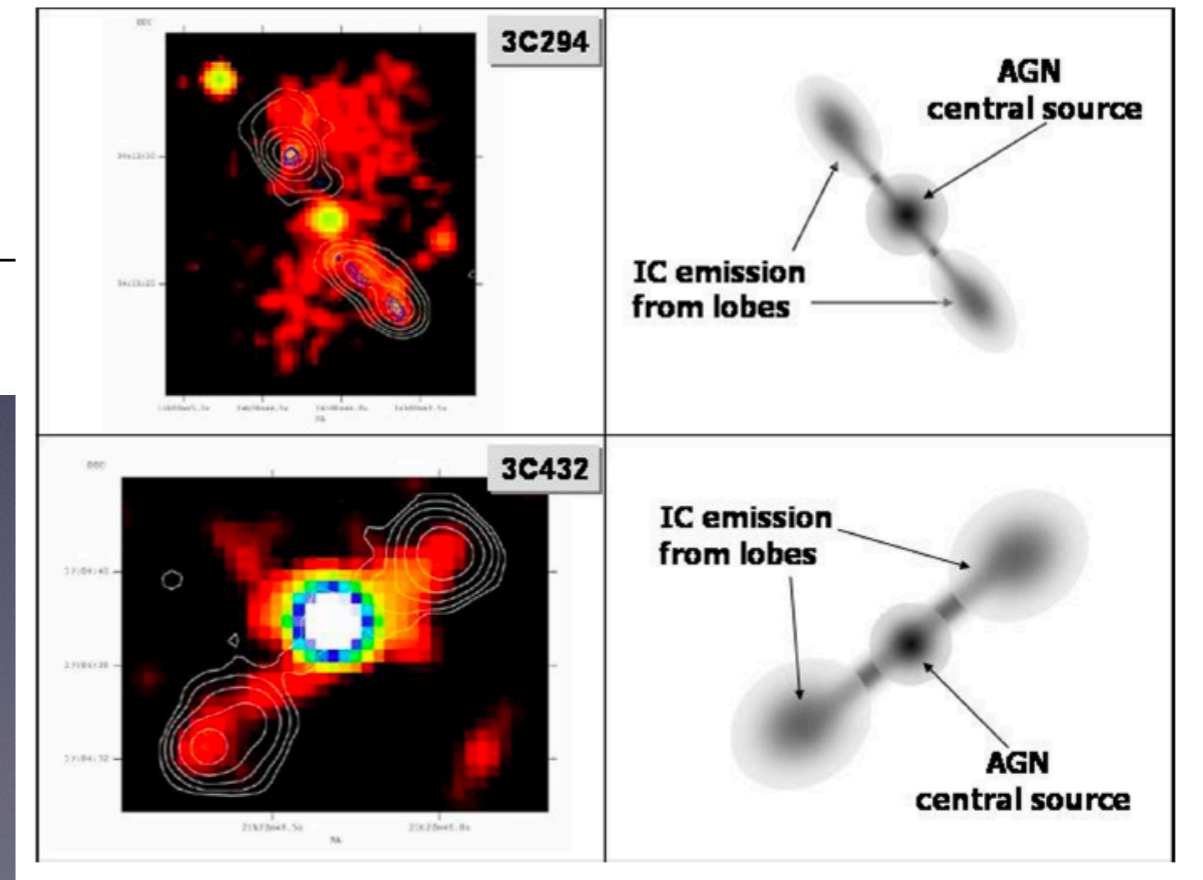
IC emission is routinely observed in hard X-ray band, from AGN lobes!

Fig. From Colafrancesco et al. (2013)

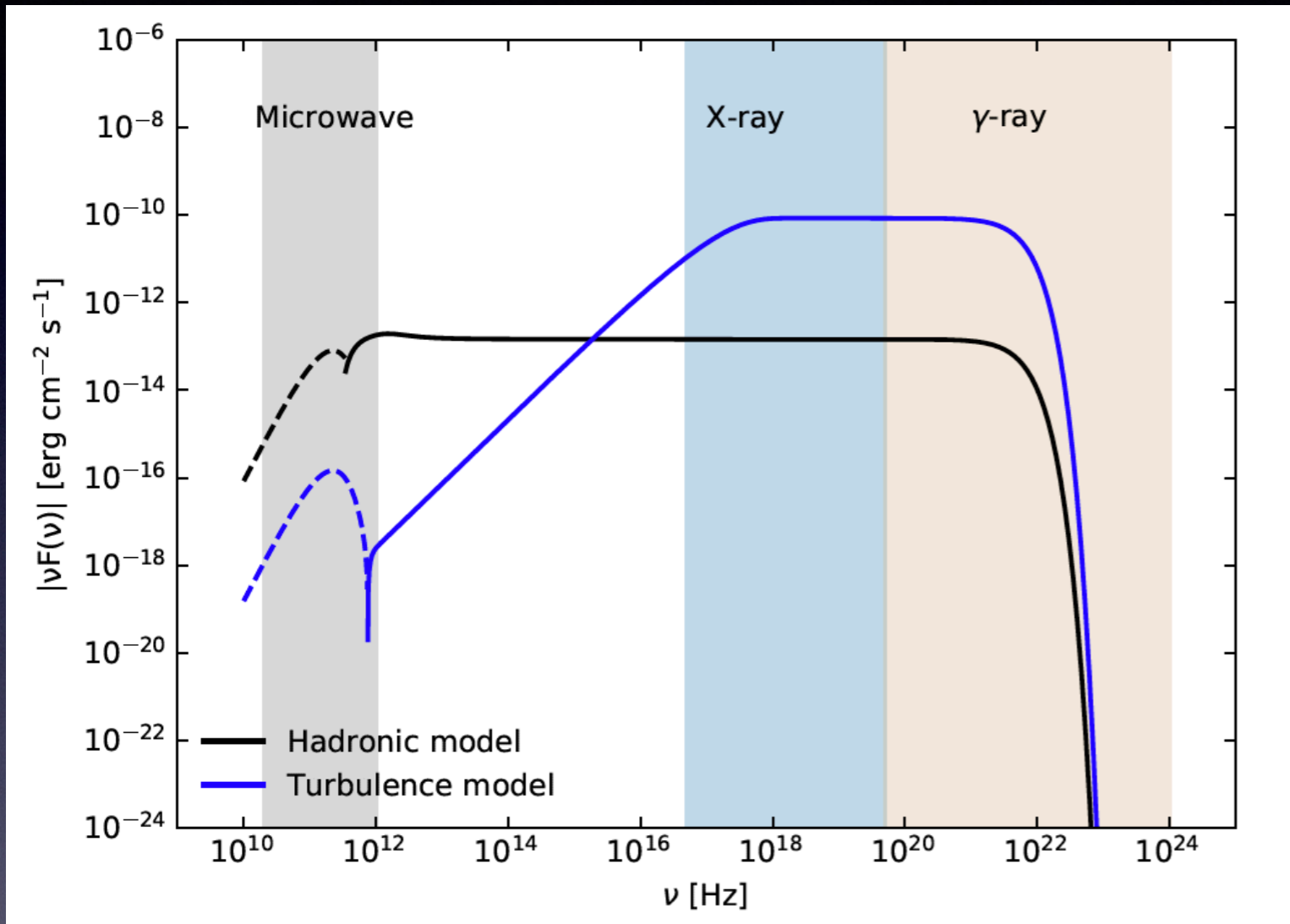
For a general isotropic momentum distribution $f(p)$:

$$\Delta I_\nu \approx I_0 x^3 \tau_e \int_0^\infty f(p) p^2 dp \int_{-s_m(p)}^{s_m(p)} P(s, p) [n_{\text{bb}}(x e^s) - n_{\text{bb}}(x)] ds.$$

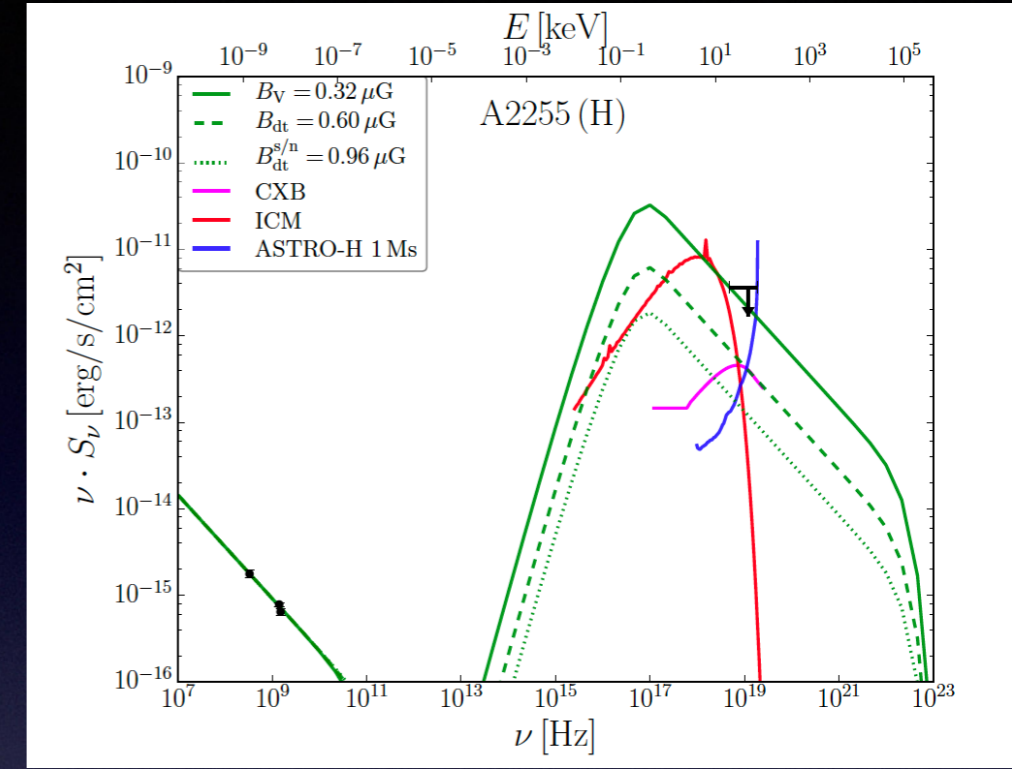
Figure above shows four cases of mono-energetic electrons (i.e. delta function distributions)



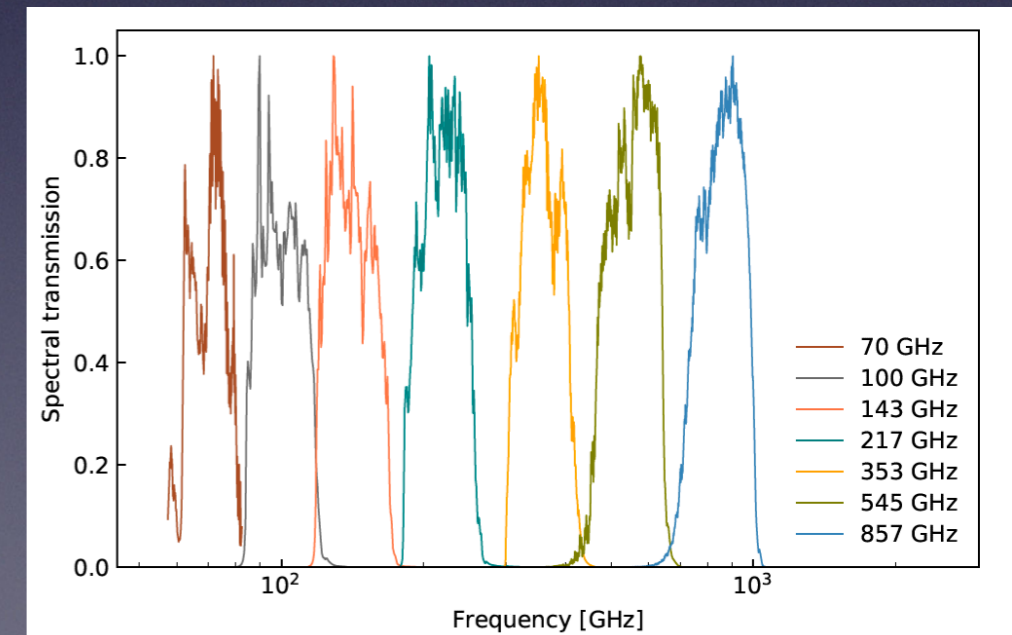
The spectrum of ntSZ effect



The same Inverse Compton signal is observable from millimeter-wave to gamma-ray energies!



X-ray IC predictions for RH clusters, from Bartels et al. (2018)

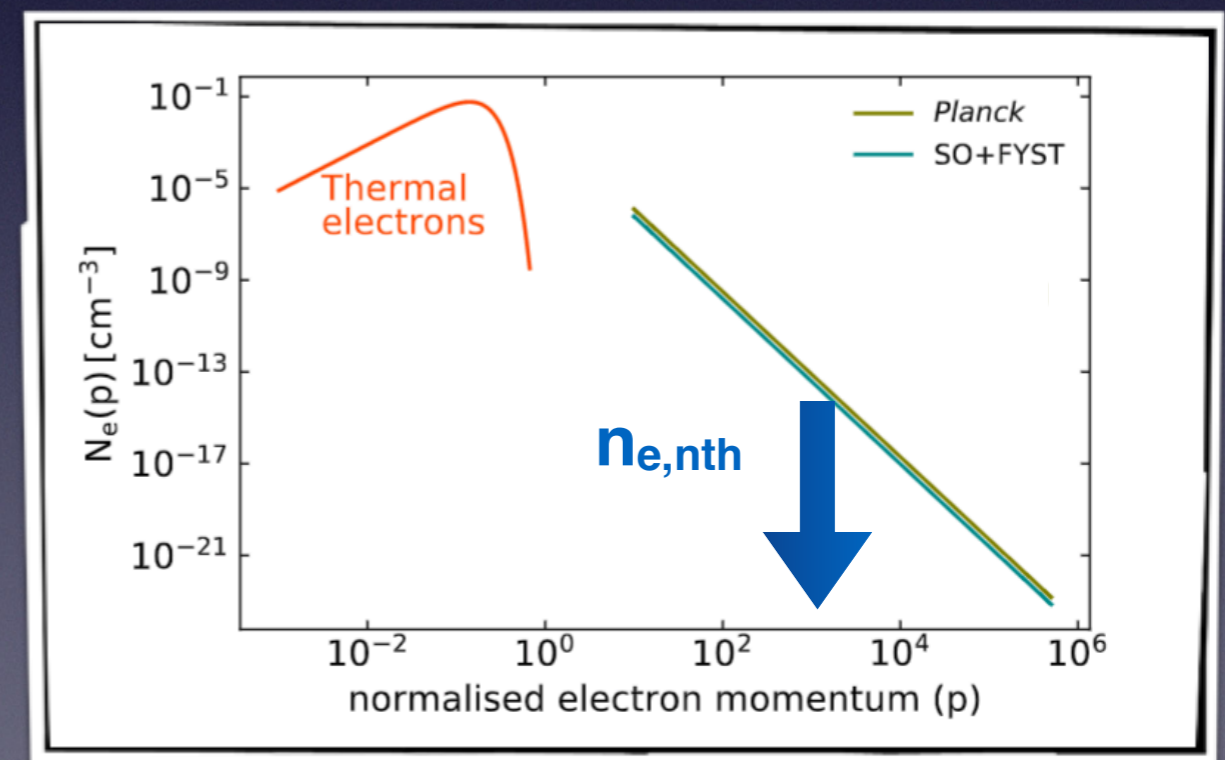
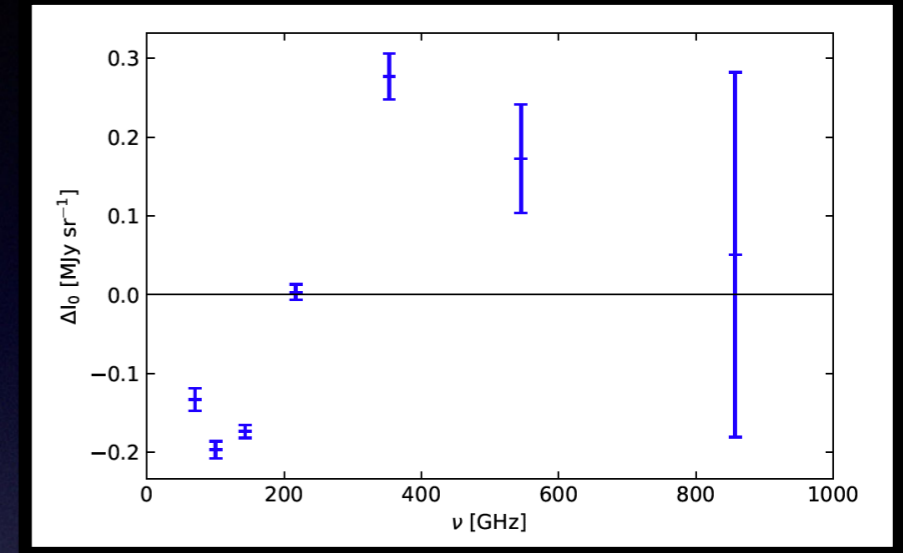
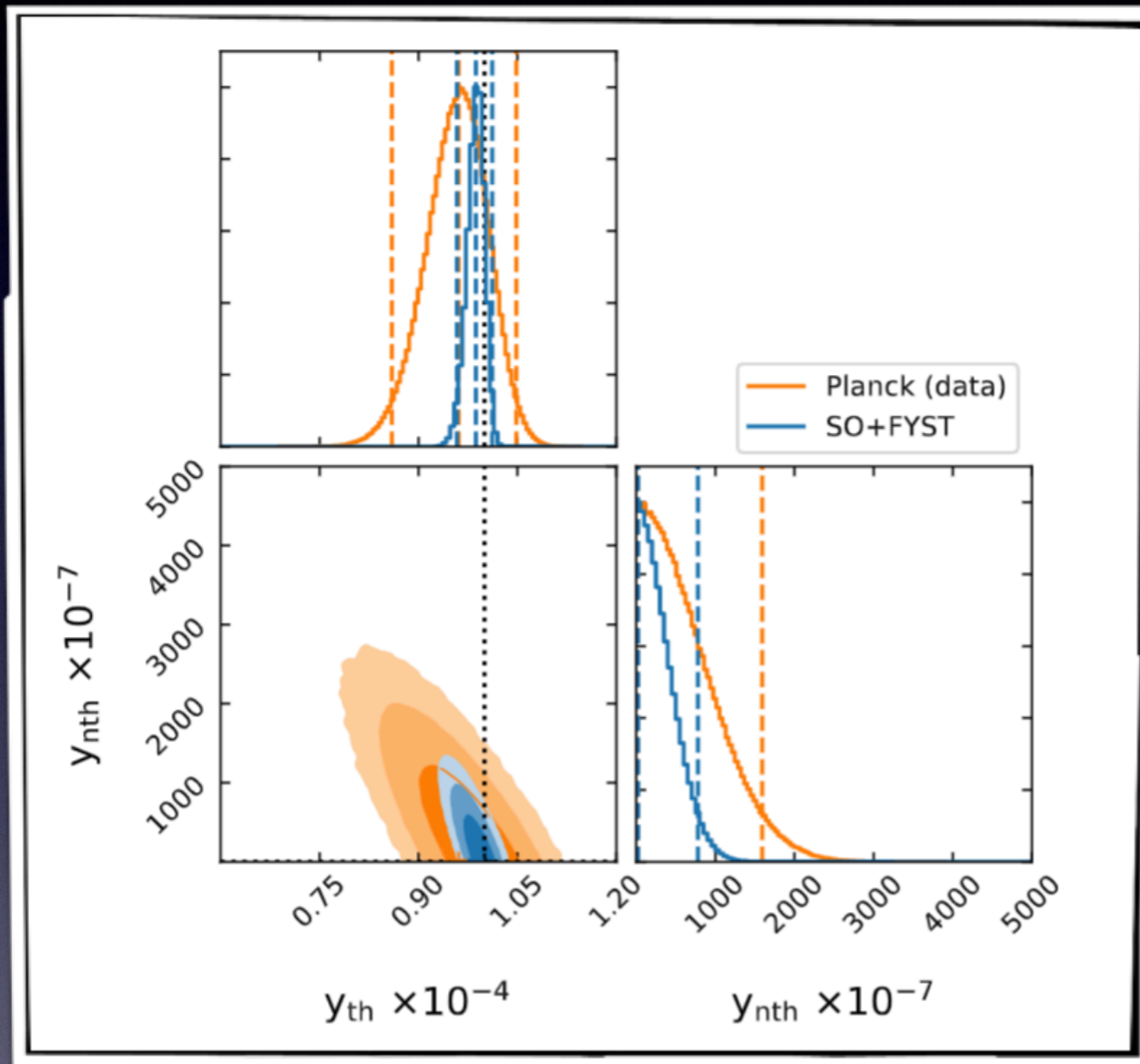


Planck frequency bands

Results from ~40 RH clusters

In clusters, ntSZ is $\approx 1\%$ of the tSZ signal 😞

Stacked spectrum of of 40 clusters from *Planck* data, after matched-filtering

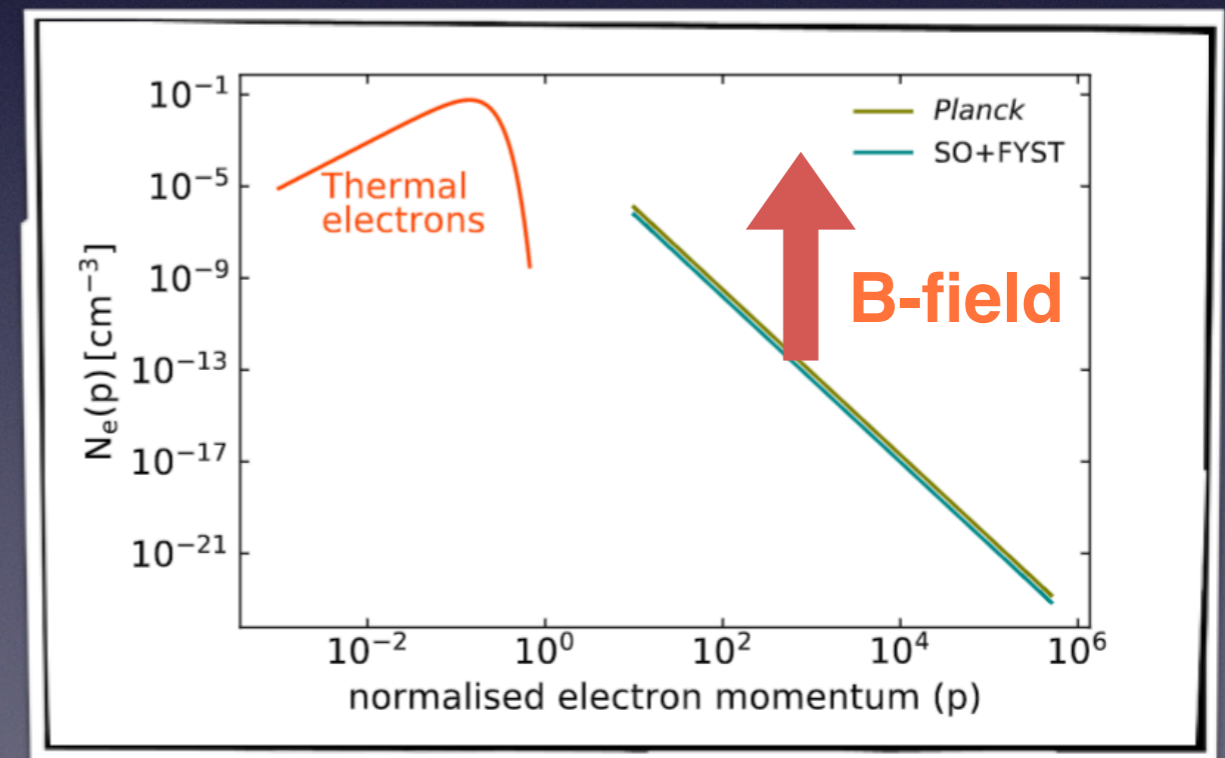
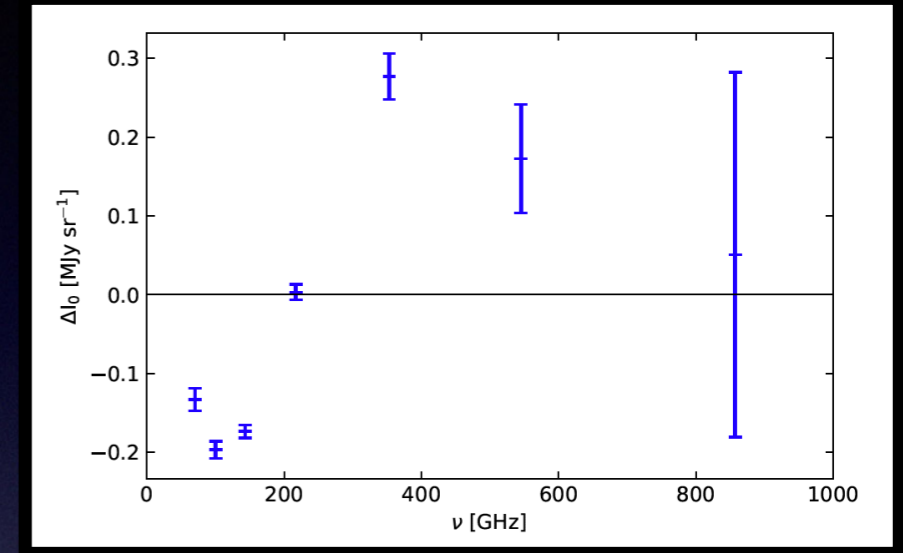
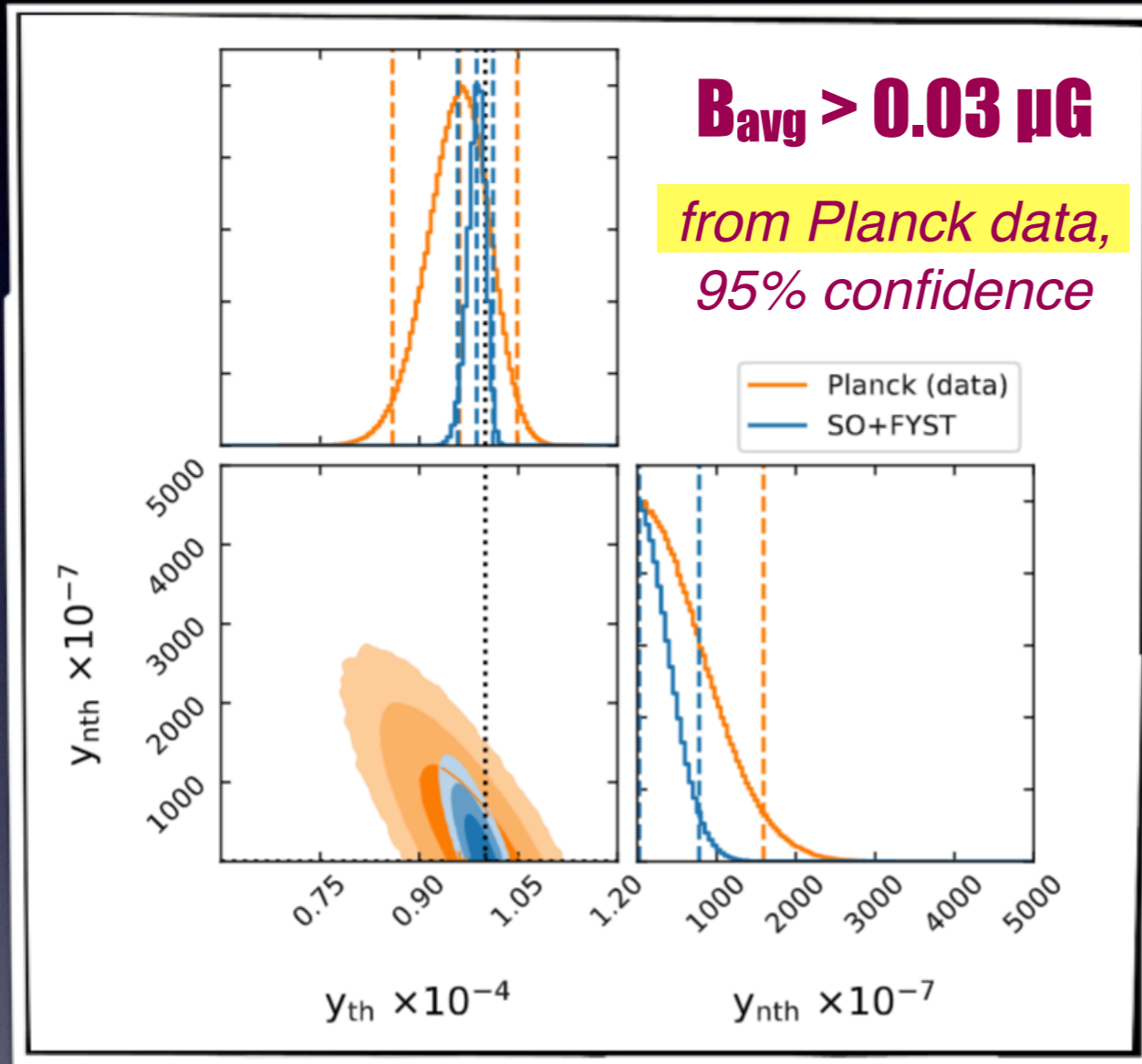


Muralidhara & Basu (in prep.)

Results from ~40 RH clusters

Muralidhara & Basu (in prep.)

Stacked spectrum of of 40 clusters from *Planck* data, after matched-filtering



With CCAT-prime and SO, order-of-magnitude better constraints on ntSZ-based B-field limits can be expected in the next ~5 years

Take home points

Many questions unanswered for cluster diffuse radio emissions. Need to look beyond radio data.

SZ selection for radio halos provide unbiased statistics. SZ observation of radio relics measure the underlying shocks.

Nonthermal SZ measurement from next-generation CMB experiments is a potential new frontier.

